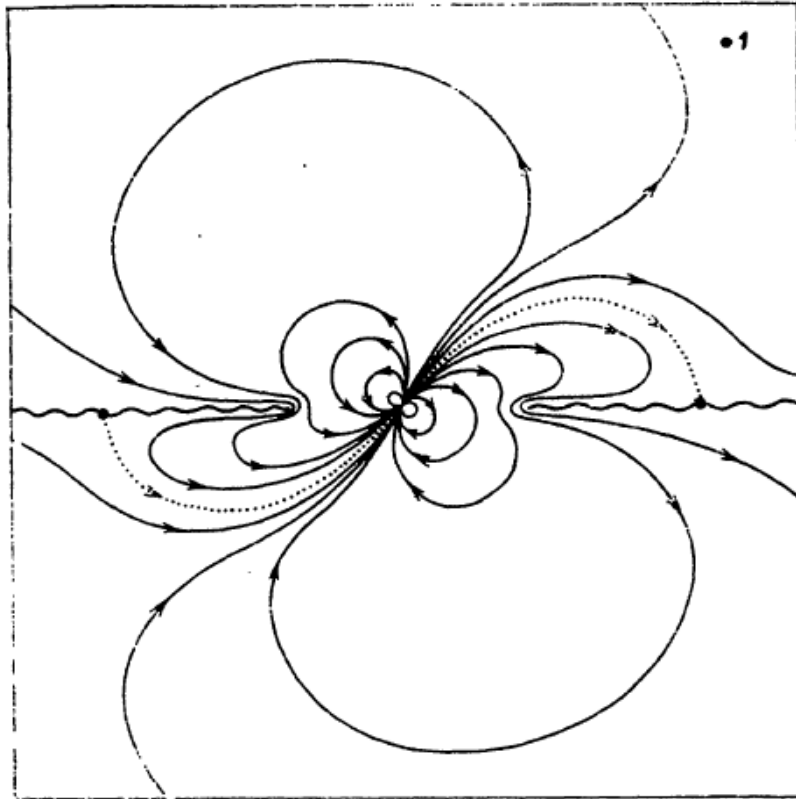
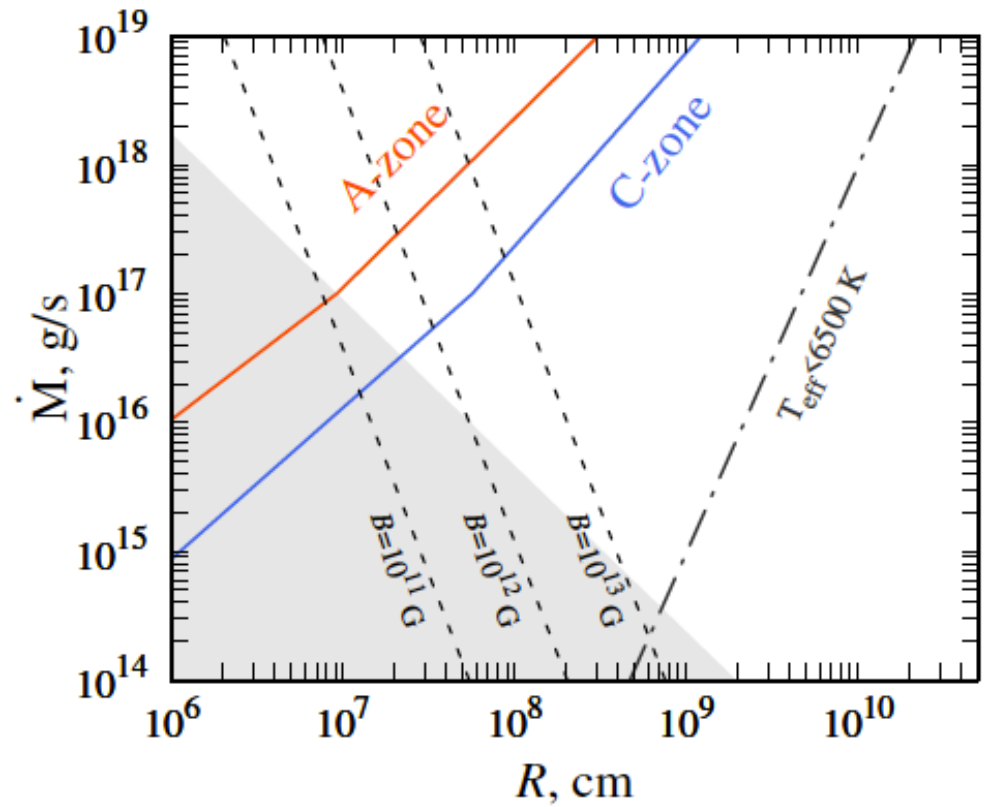


Accreting magnetized neutron stars: X-ray pulsars

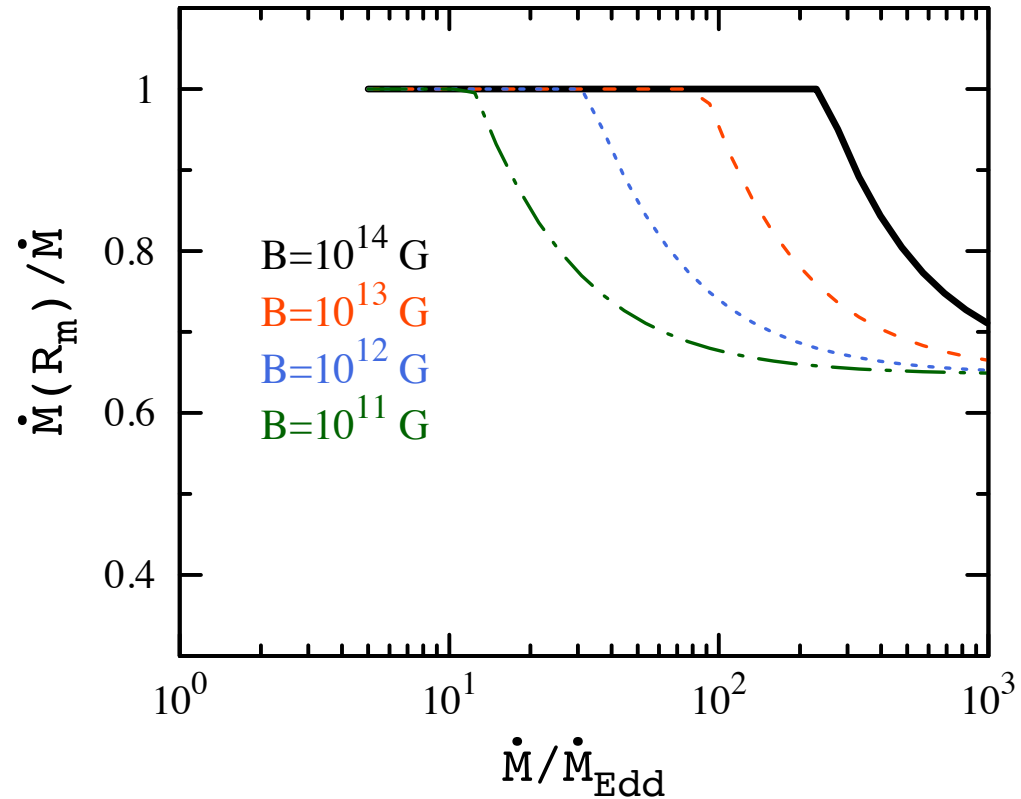
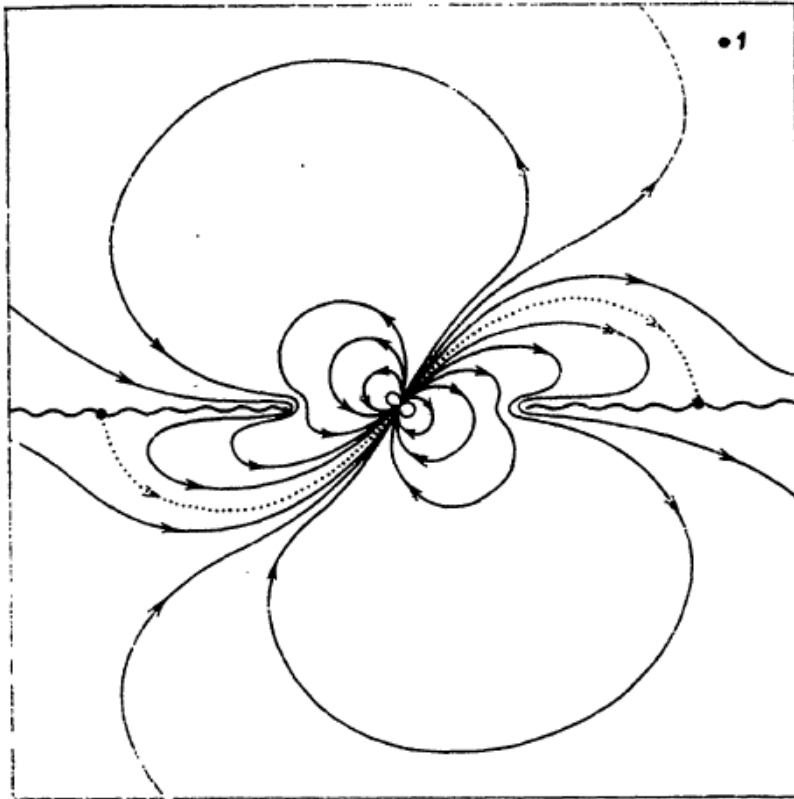


Accretion discs in XRP

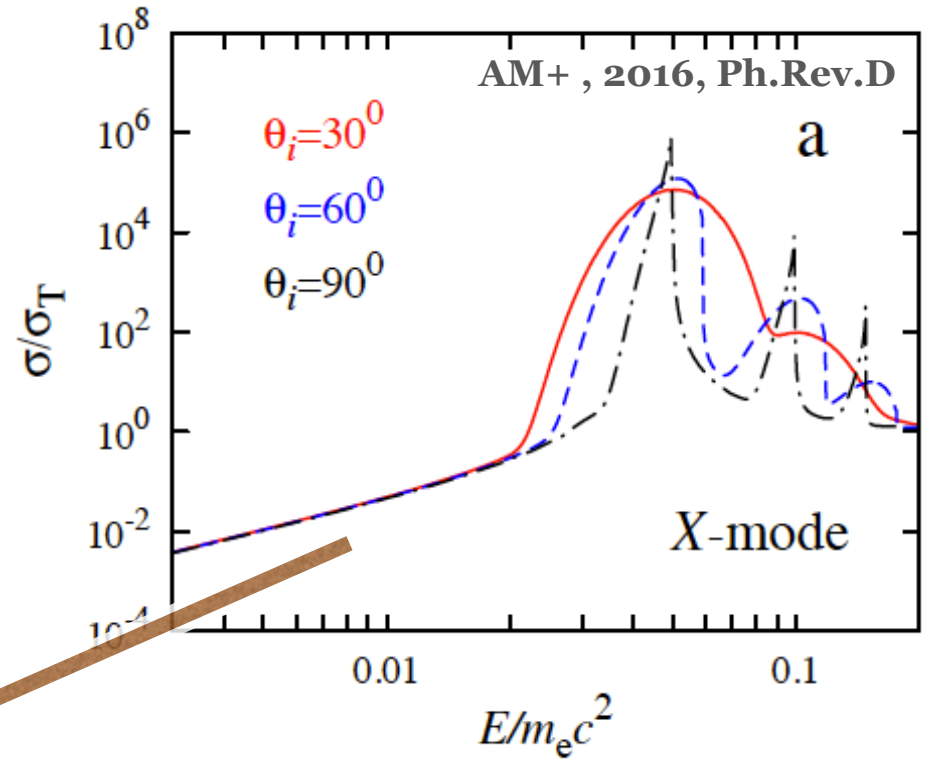
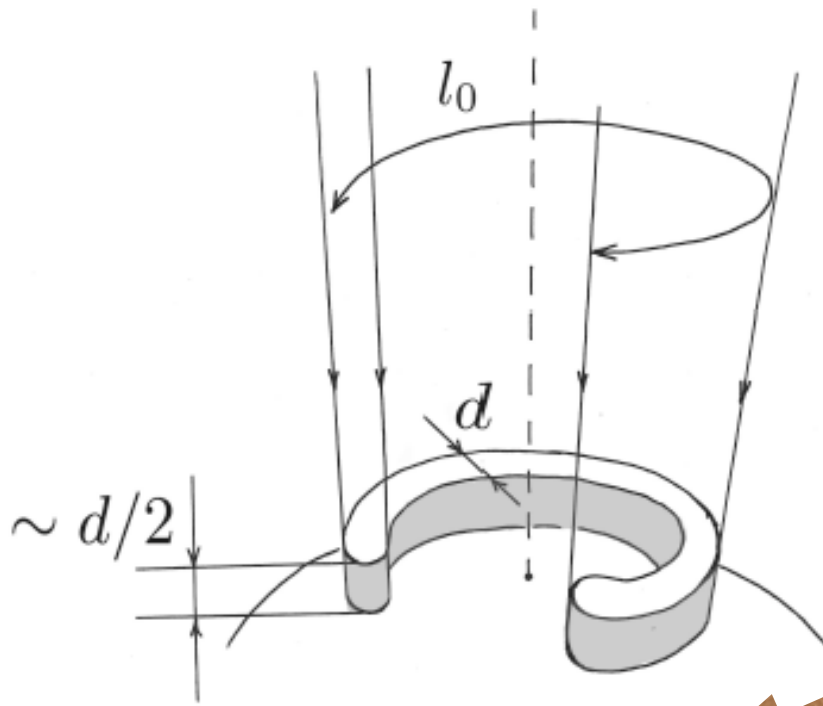


Accreting magnetized neutron stars: X-ray pulsars

Mass outflow and mass accretion rates at the magnetosphere



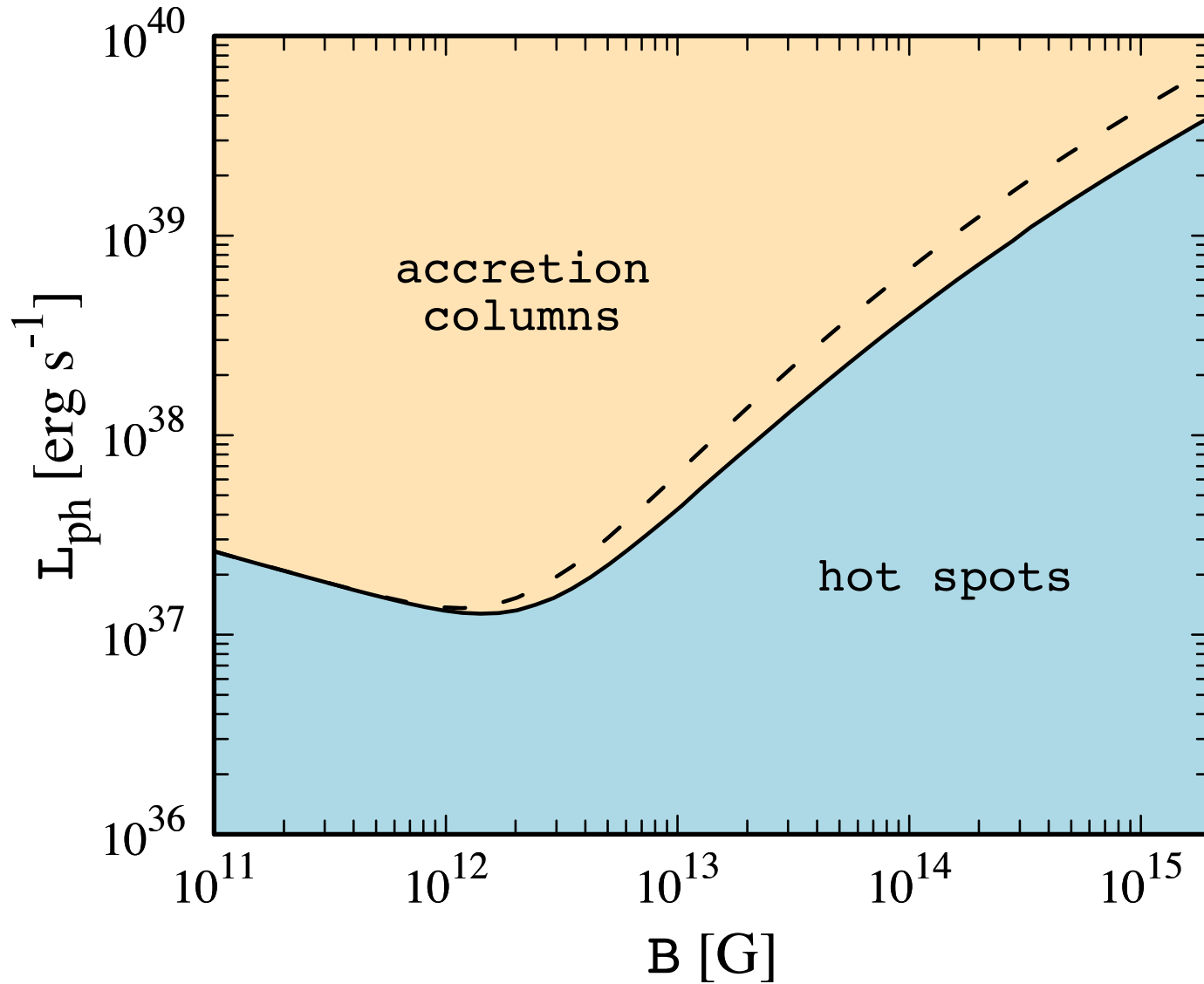
Critical luminosity: matter is stopped by radiation



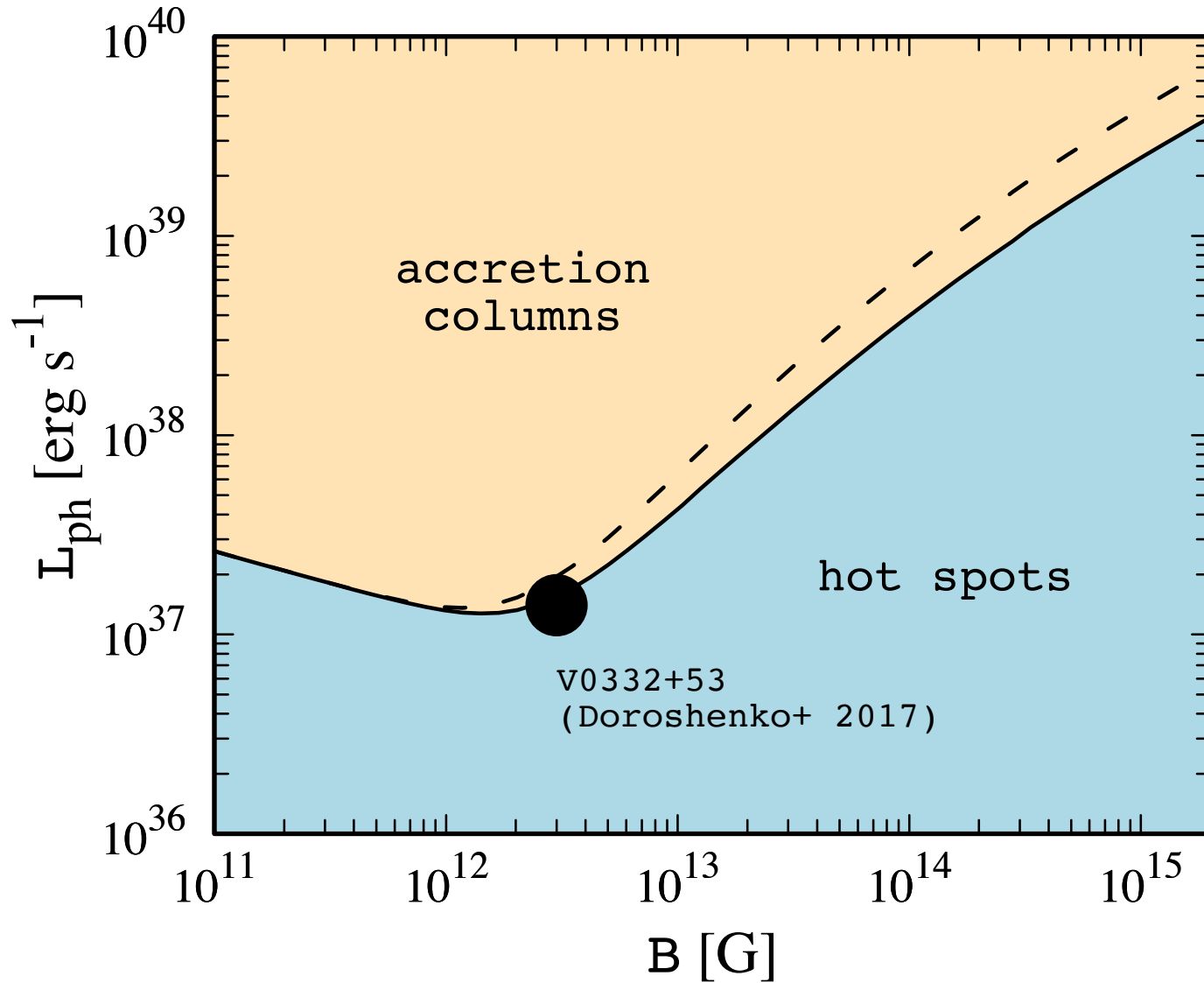
$$L^* = 4 \times 10^{36} \left(\frac{\sigma_T}{\sigma_{eff}} \right) \left(\frac{l_0}{2 \times 10^5 [cm]} \right) \frac{1}{R_6} \frac{M}{M_\odot} \quad [\text{erg s}^{-1}]$$

$$L_{\text{Edd}} = \frac{4\pi GM m_p c}{\sigma_T} \approx 1.3 \times 10^{38} \frac{M}{M_\odot} \text{ erg s}^{-1}$$

Critical luminosity



Critical luminosity



Above the critical luminosity: accretion column

Dipole magnetic field.

*

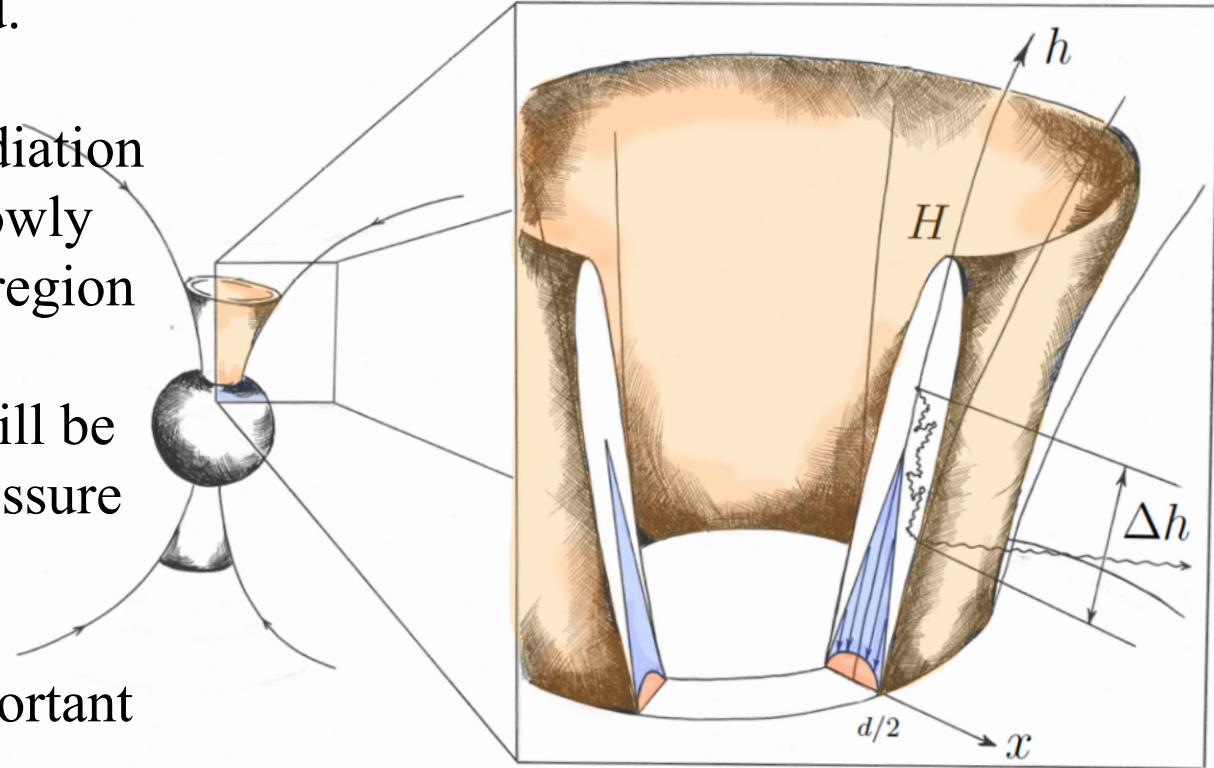
Accretion flow stops at radiation dominated shock and slowly settles in inside a sinking region

*

The gravitational force will be offset by the radiation pressure gradient only

*

The gas pressure is unimportant

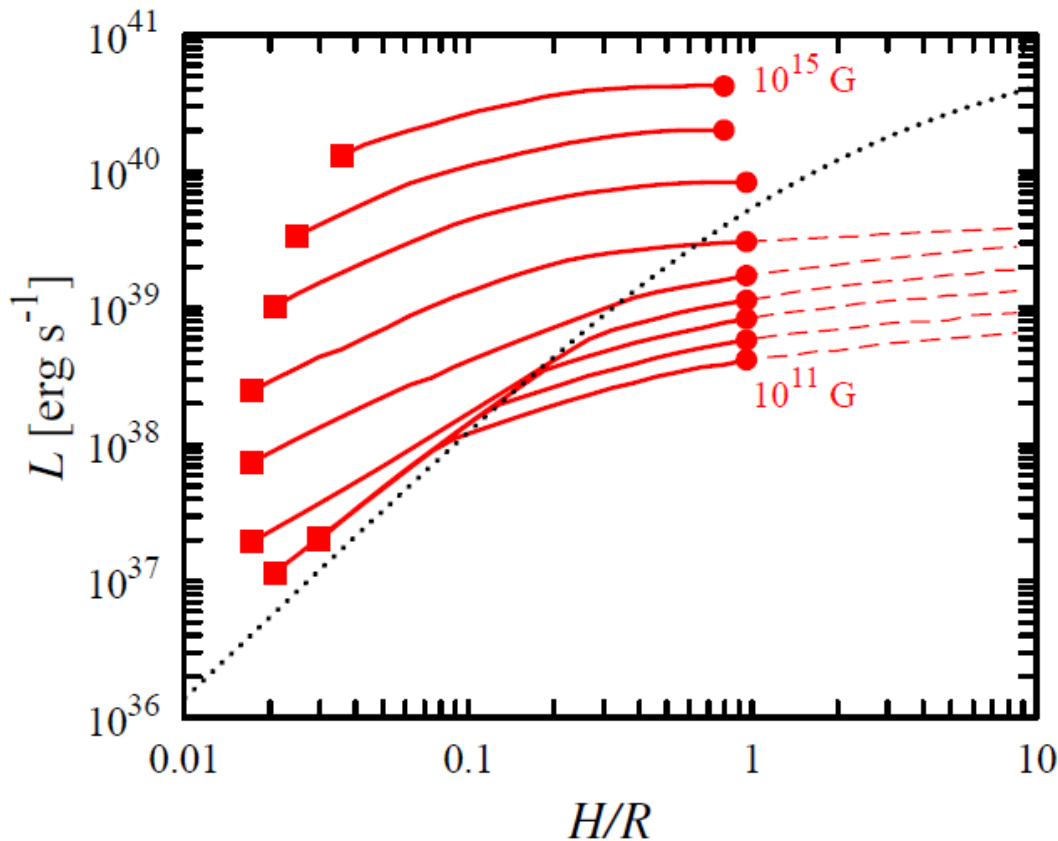


Analytical estimation:

$$L(H = R) \approx 1.8 \times 10^{39} \left(\frac{l_0/d_0}{50} \right) \left(\frac{\kappa_T}{\kappa_{\perp}} \right) \frac{M}{M_{\odot}} \text{erg s}^{-1}$$

Accretion column: height and luminosity

Luminosity of accretion column as a function of column height



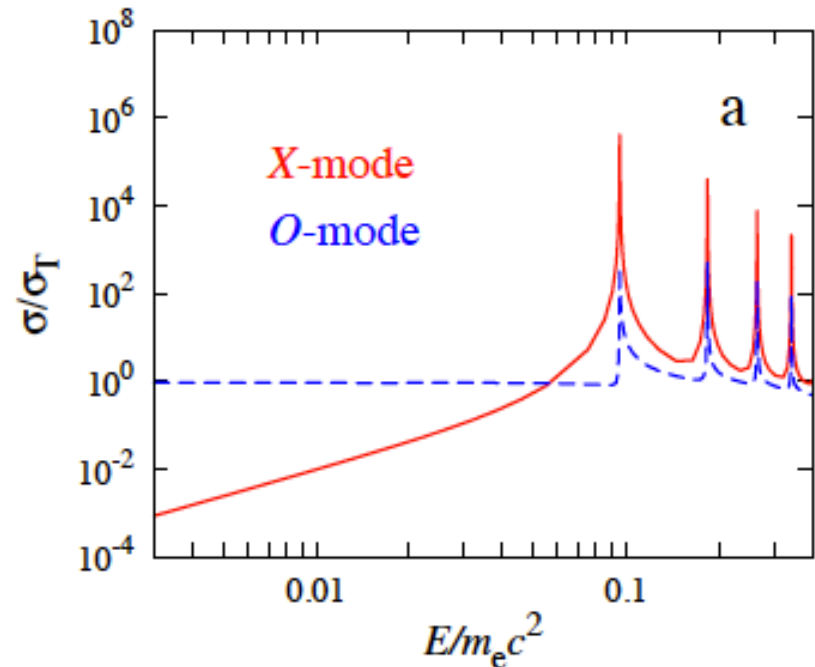
Scattering cross-section:

The higher the mass accretion rate

=> the higher the temperature

=> the higher photon energy

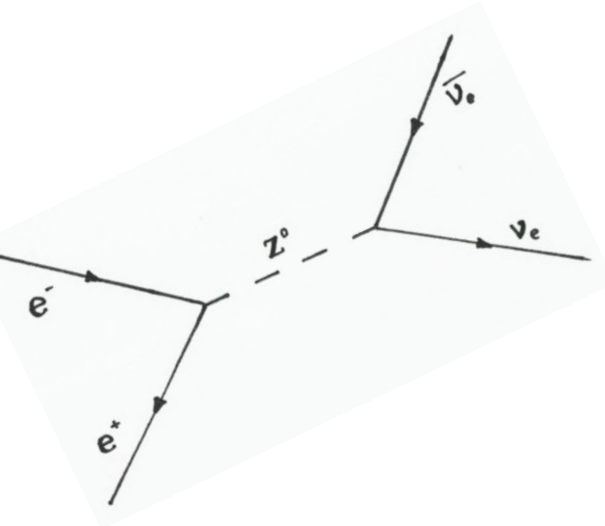
=> the higher the cross-sections



Neutrino pulsars

Typical time of photon escape:

$$t_{\text{diff}} = \frac{\tau d}{2c} \approx 5 \times 10^{-4} \frac{\dot{m}_{10} d_4^2 \kappa_e}{\beta} \text{ s}$$

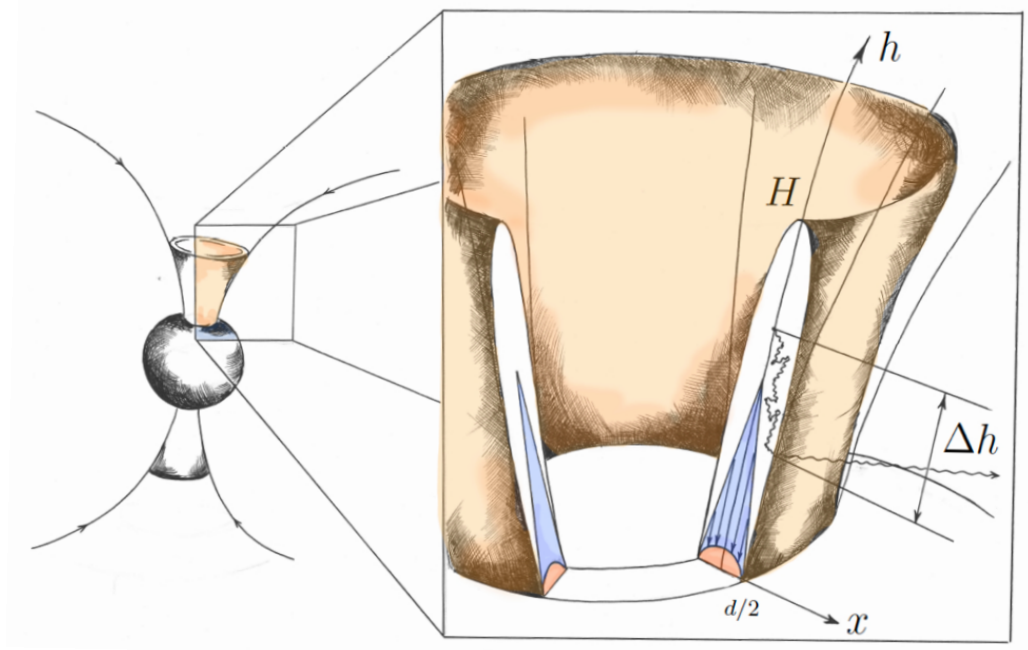


$$\frac{\partial}{\partial h} \left[\left(-\frac{\rho GM}{R+h} + \frac{\rho v^2}{2} + \epsilon_{\text{tot}} + P_{\text{tot}} + 2n_+ m_e c^2 \right) v \right] = Q^-$$



The total accretion luminosity:

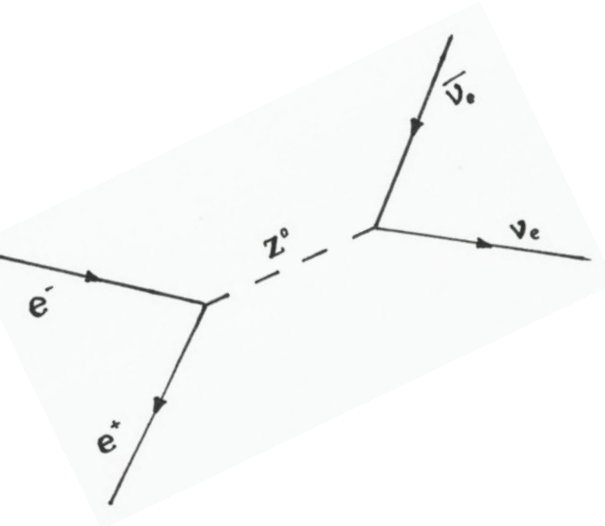
$$L_{\text{tot}} = \frac{GM\dot{M}}{R} = L_{\text{ph}} + L_{\nu}$$



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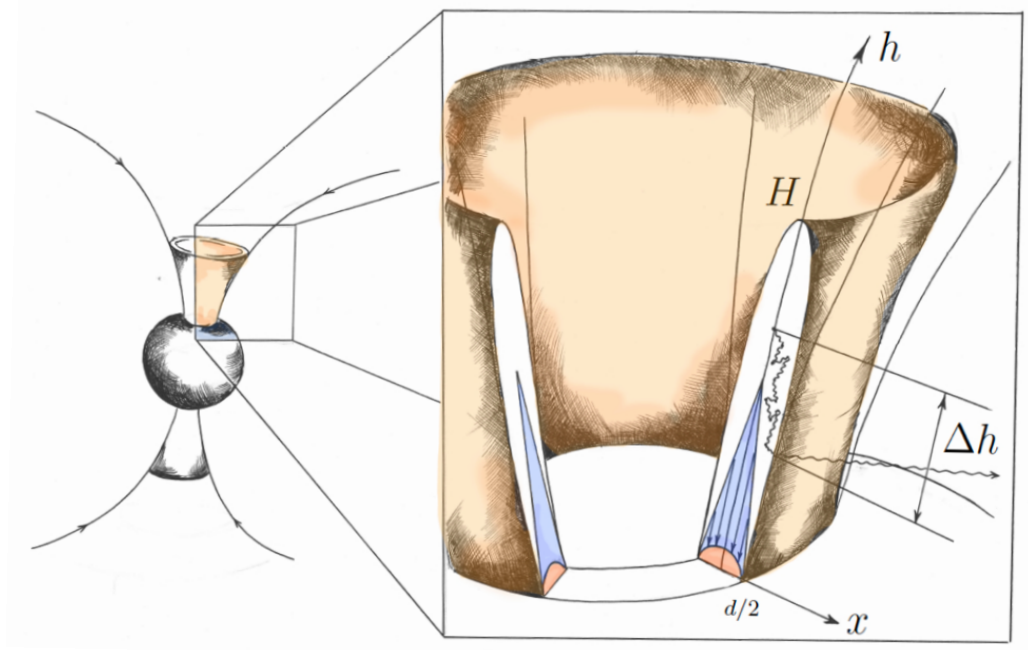


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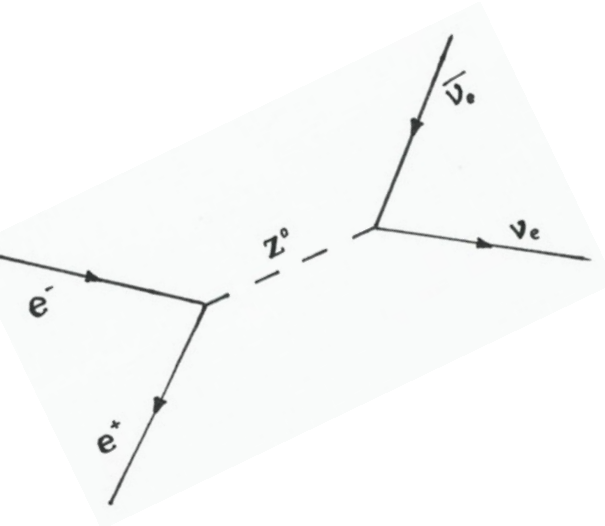
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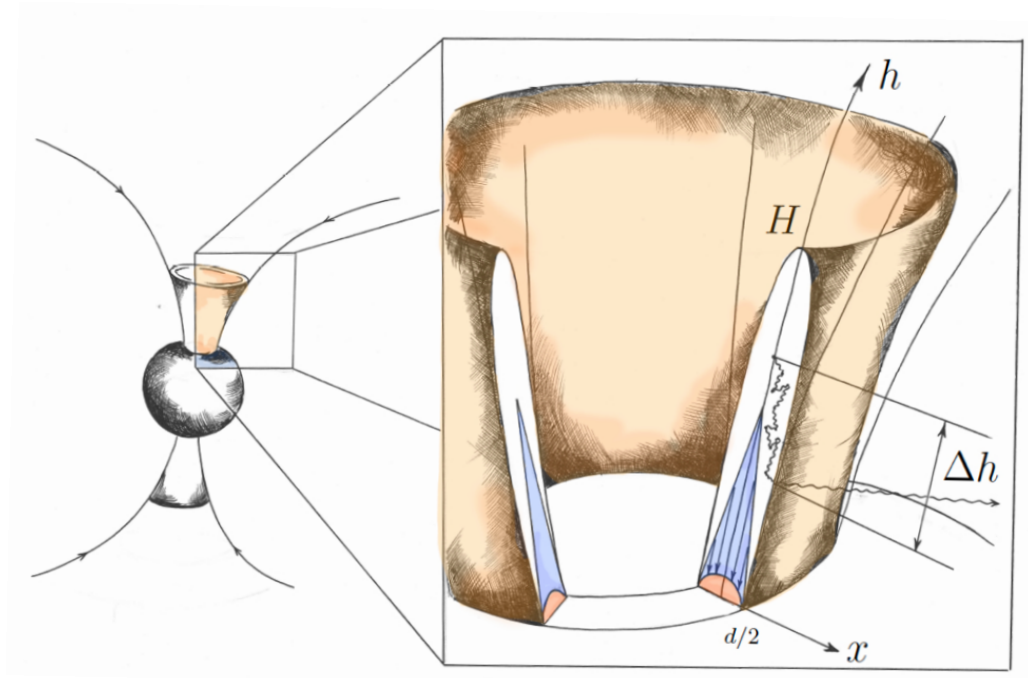


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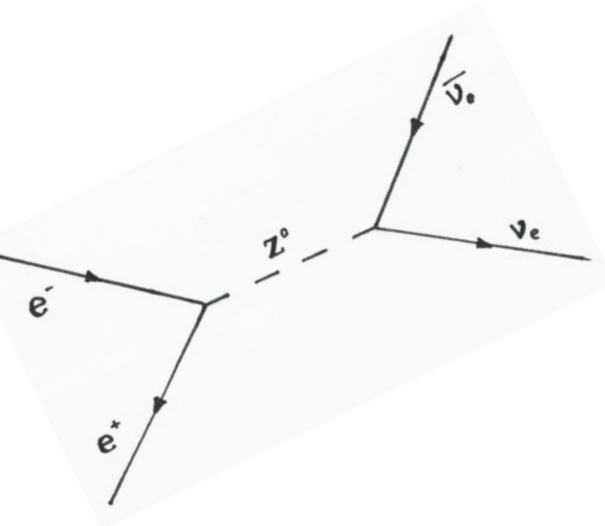
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Neutrino pulsars

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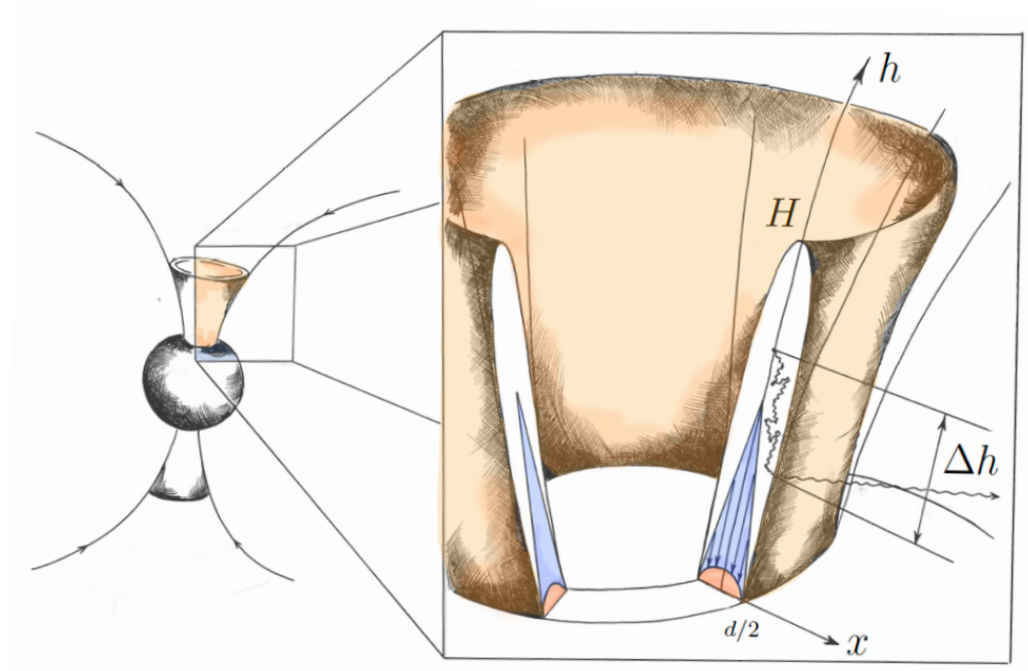


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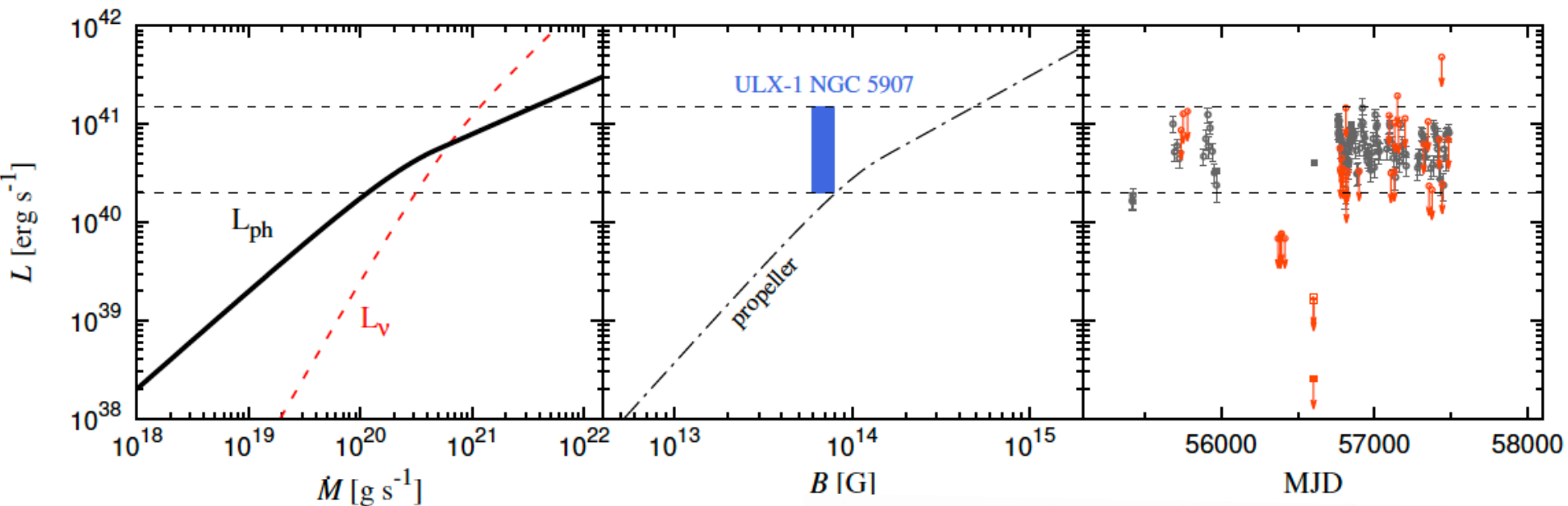


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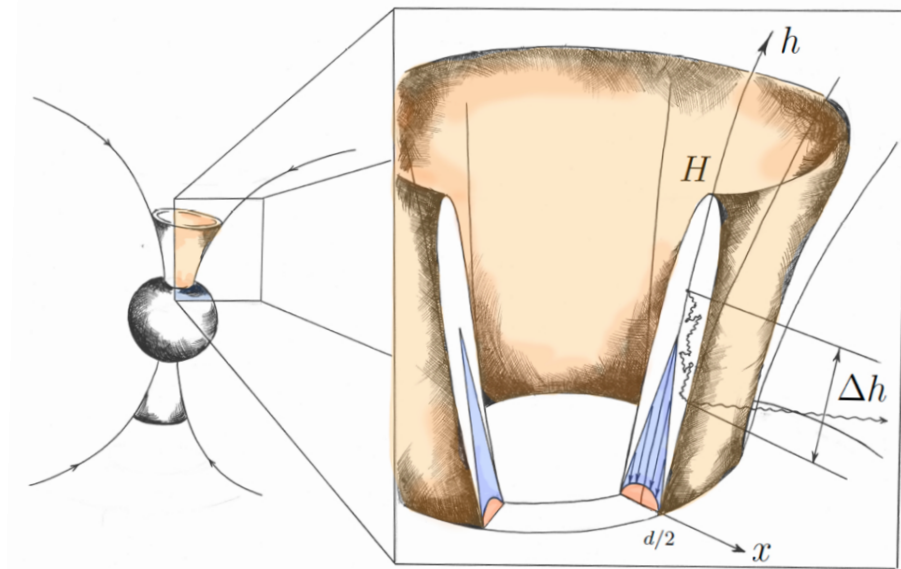


Neutrino pulsars



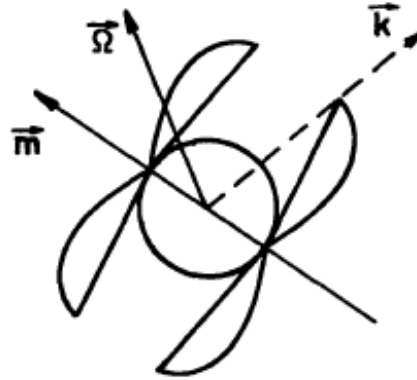
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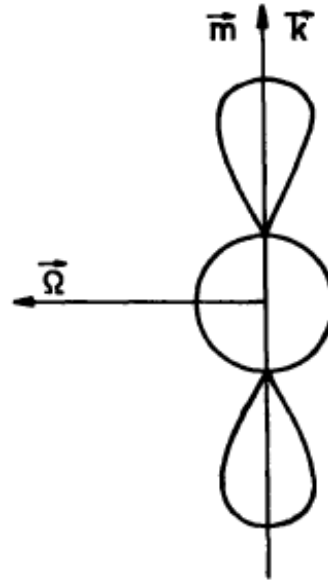


Accretion column: radiation beaming

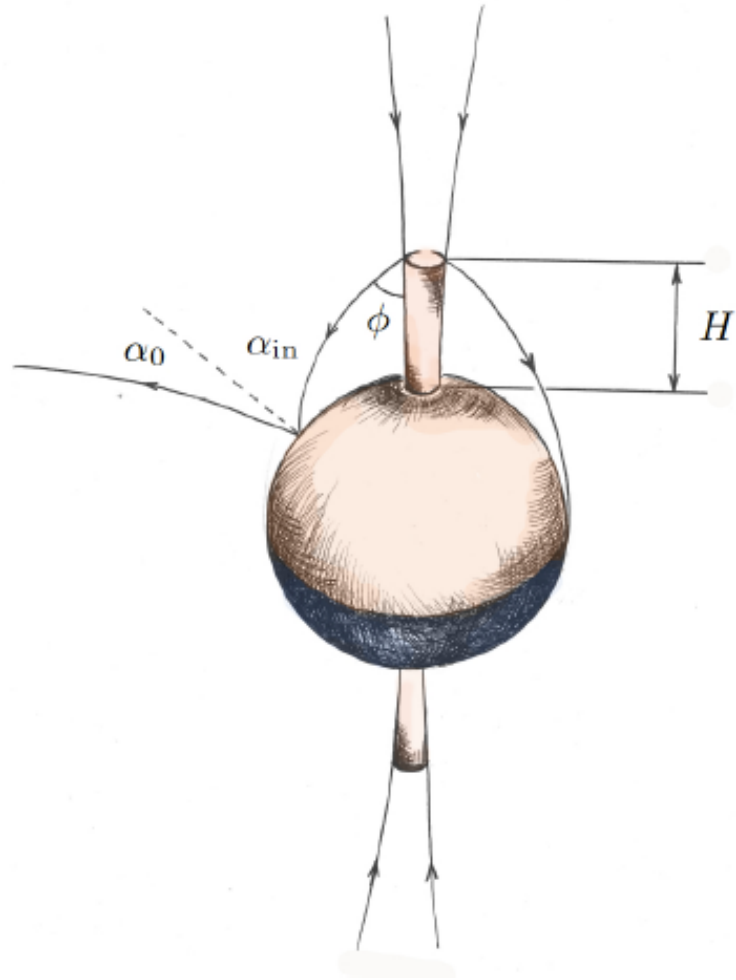
Accretion column case



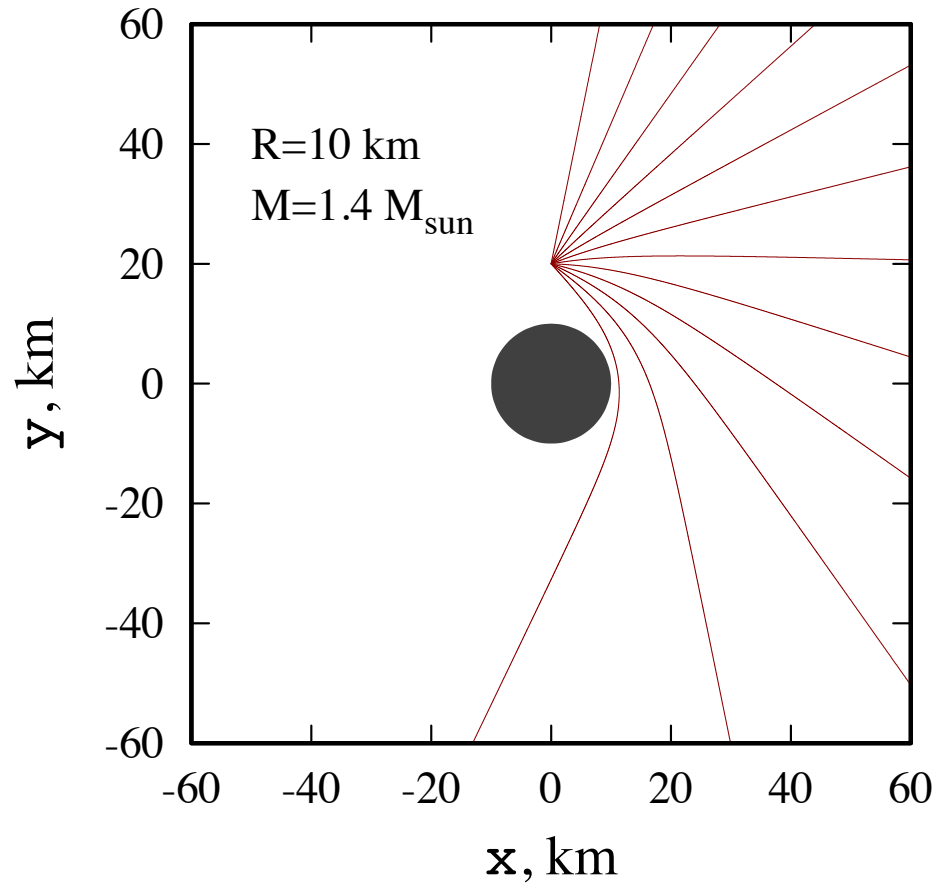
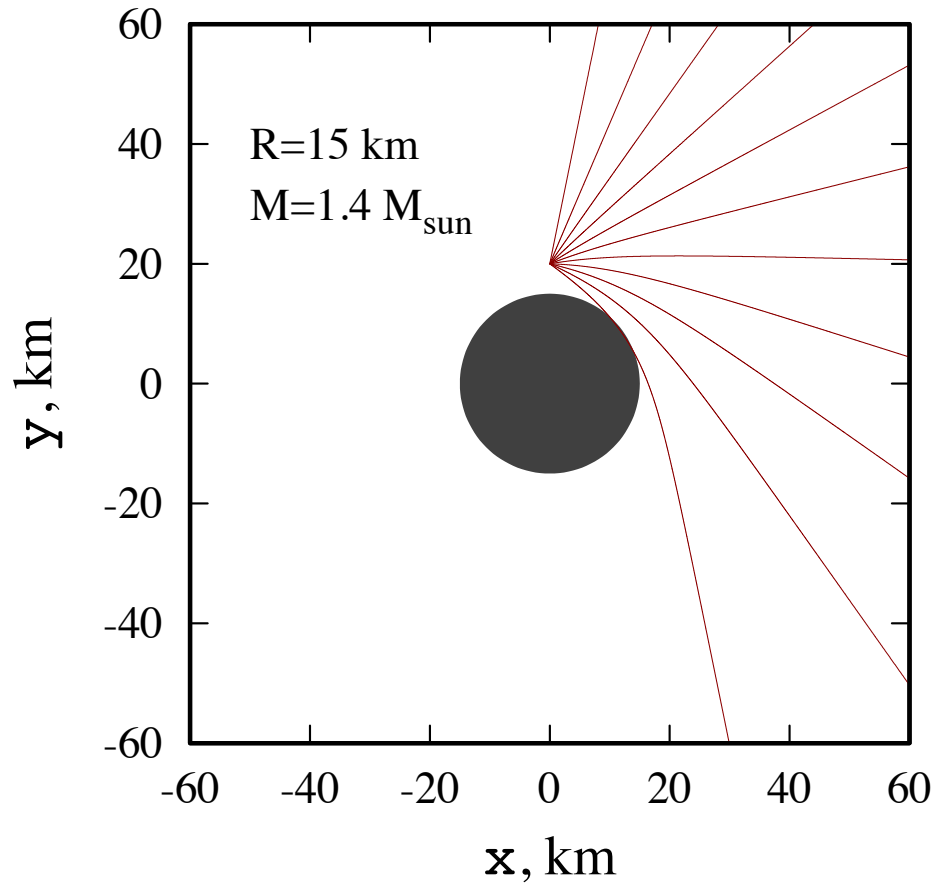
Hot spots case



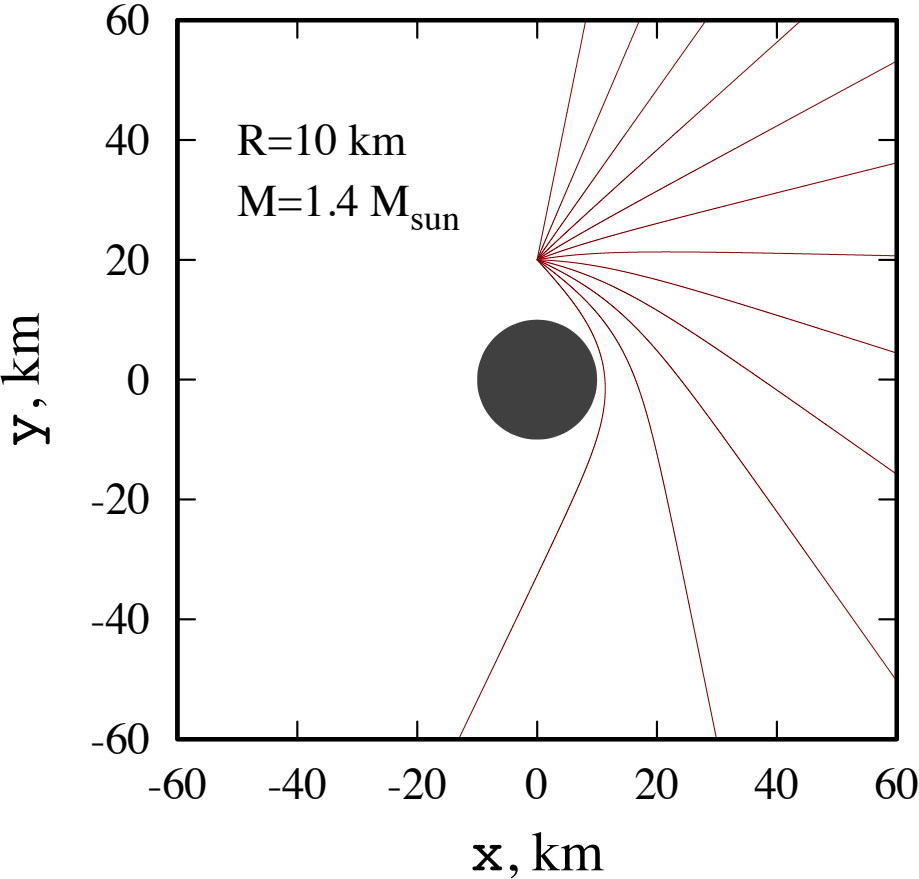
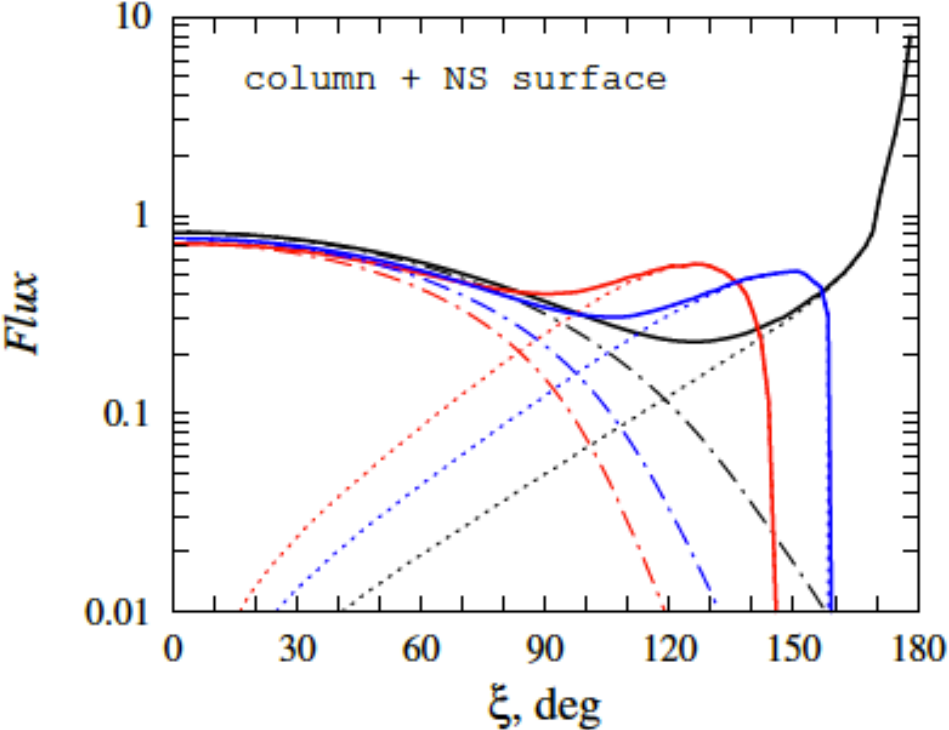
Accretion column: radiation beaming



Accretion column: radiation beaming

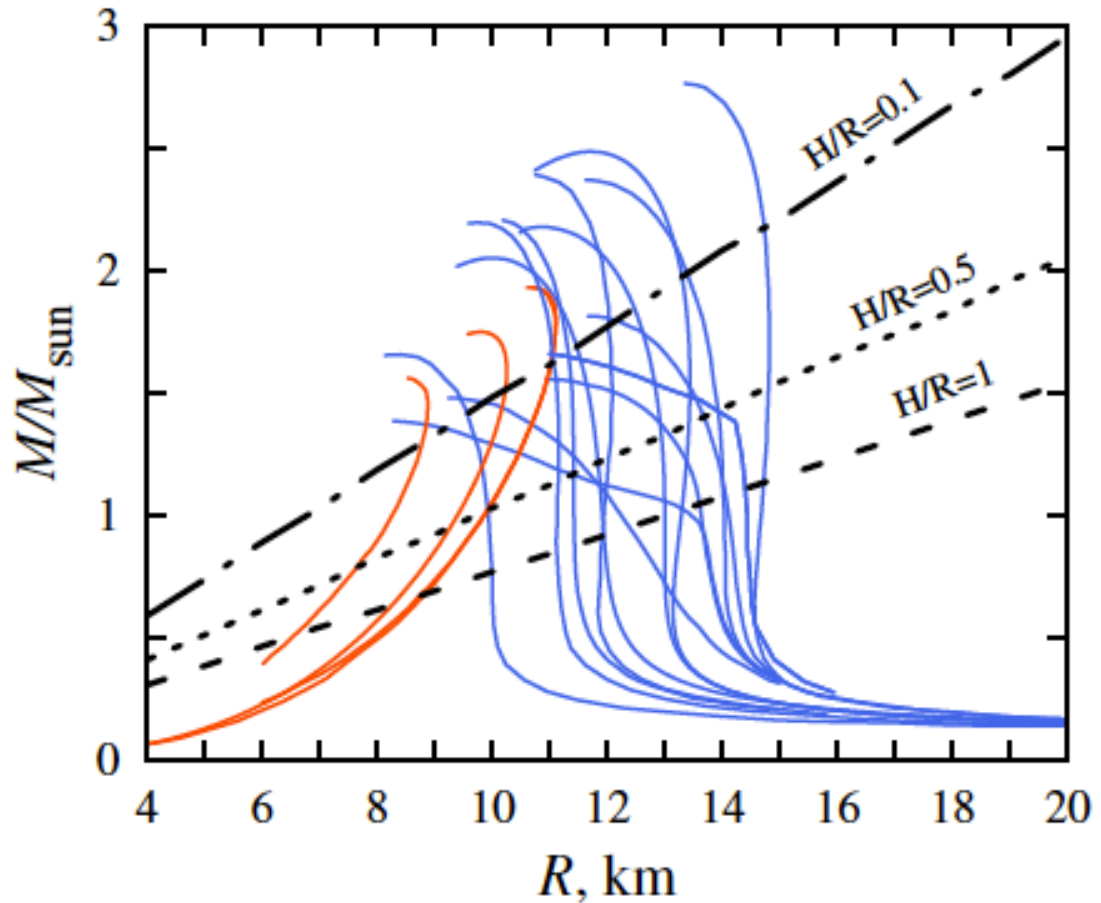
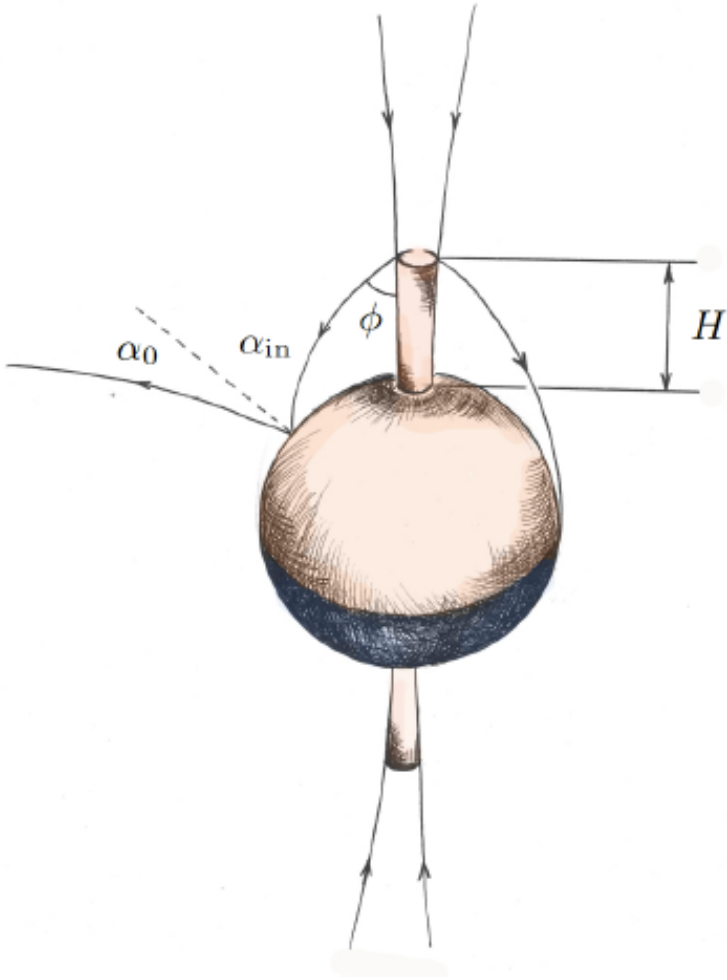


Accretion column: radiation beaming

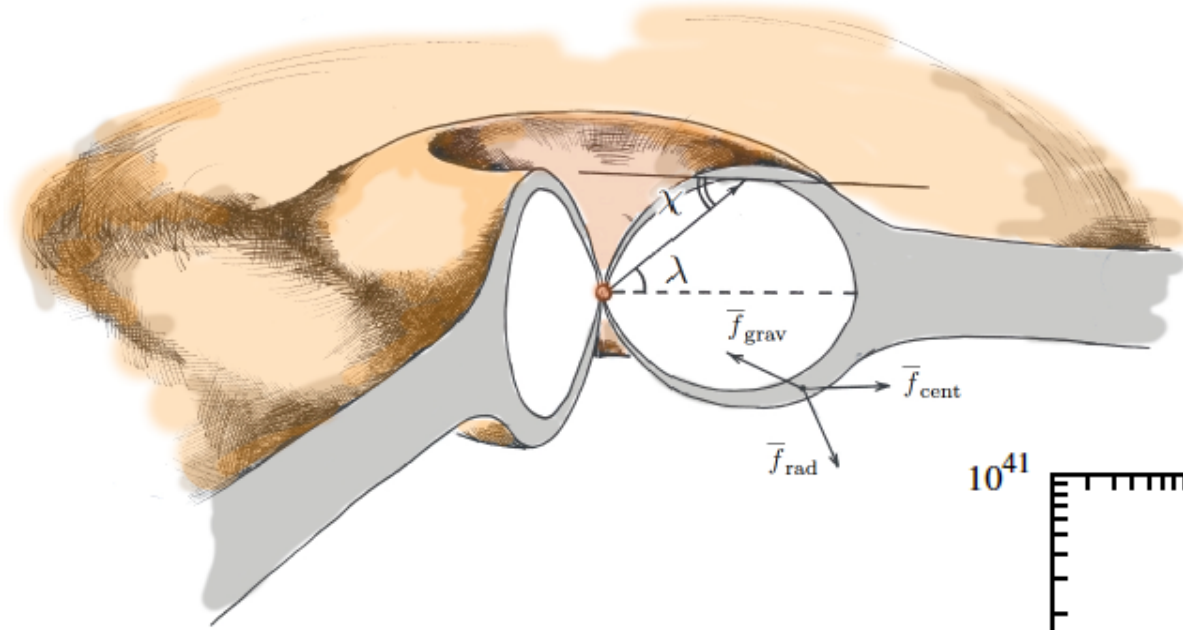


Accretion column: radiation beaming

Accretion column eclipses
and
neutron star mass-radius relation



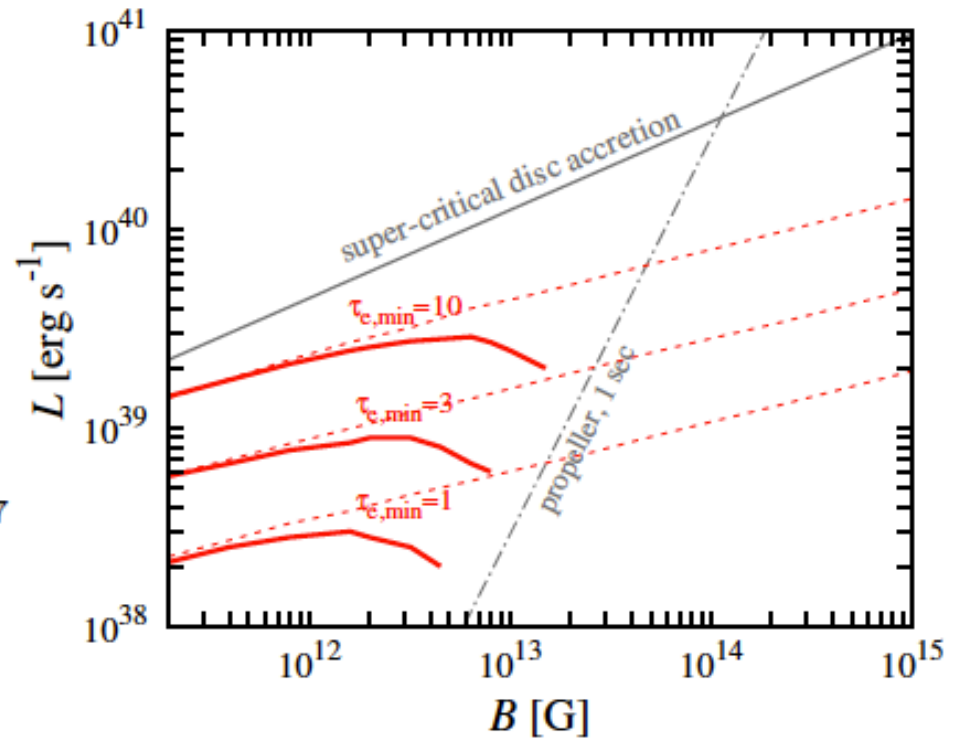
Accretion envelope



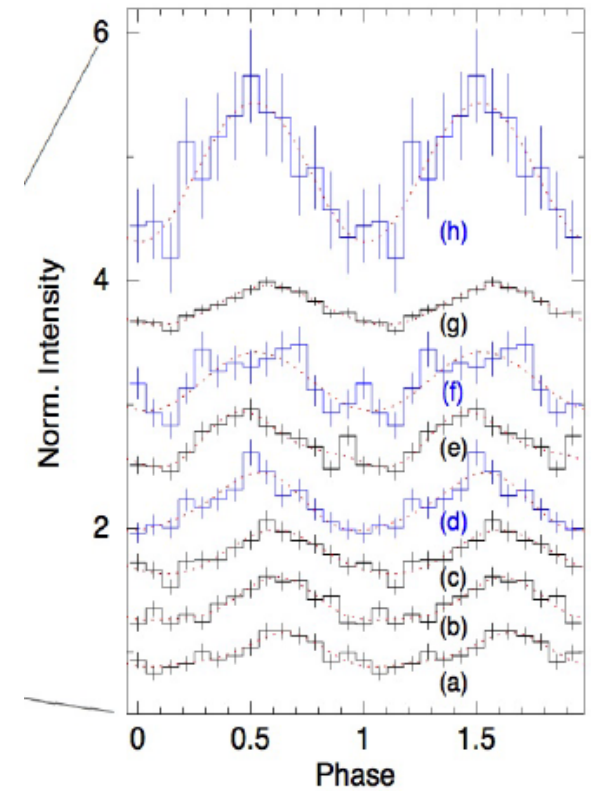
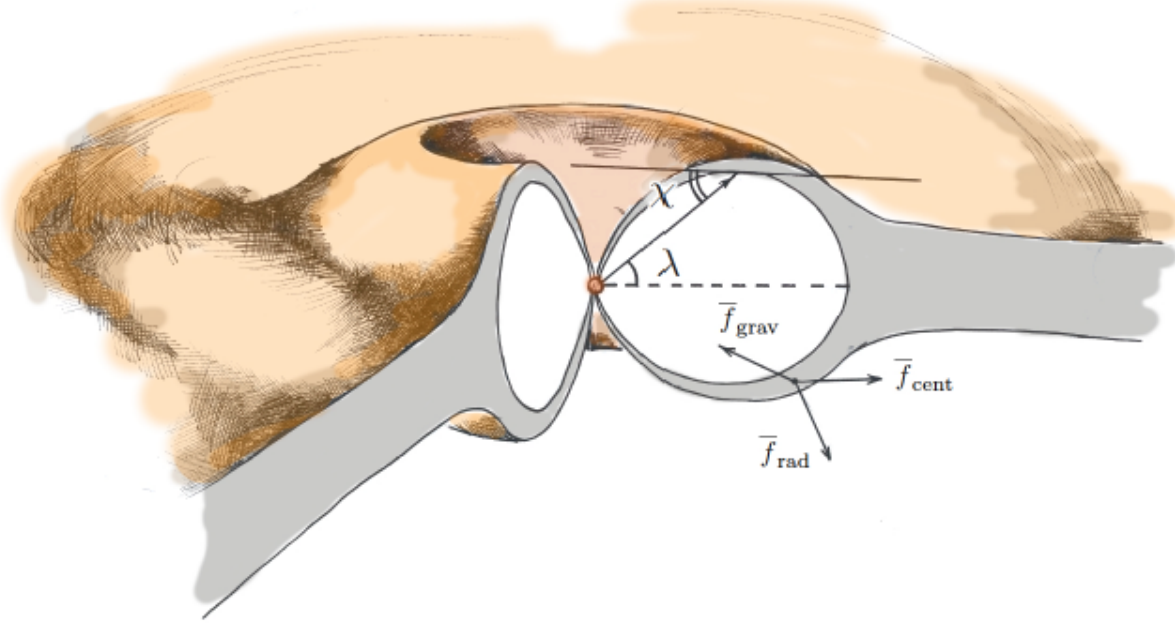
Stability

$$P_{\text{mag}}(\lambda = 0) \propto L^{12/7} B^{-10/7}$$

$$\frac{P_{\text{rad}}}{P_{\text{mag}}} \lesssim 7 \times 10^{-2} \tau L_{39}^{-1/7} B_{12}^{2/7} m^{2/7} R_6^{-2/7}$$



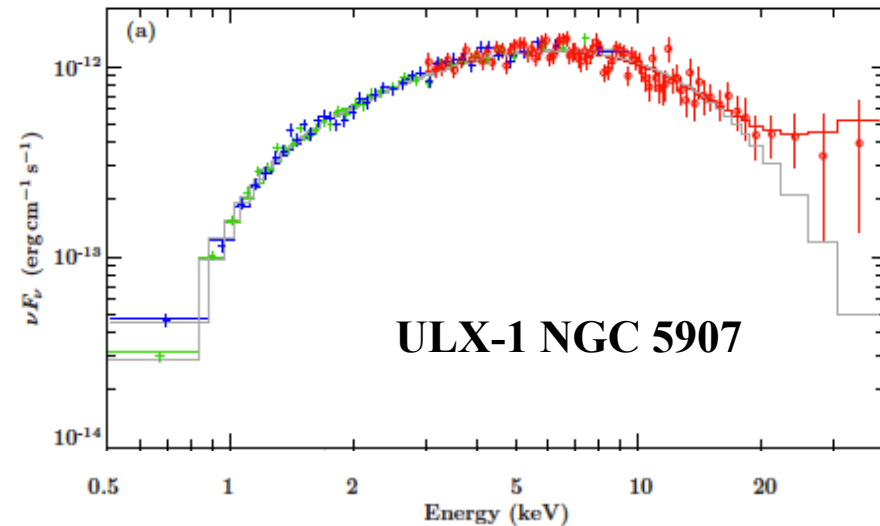
Accretion envelope



Israel+ Science, 2017

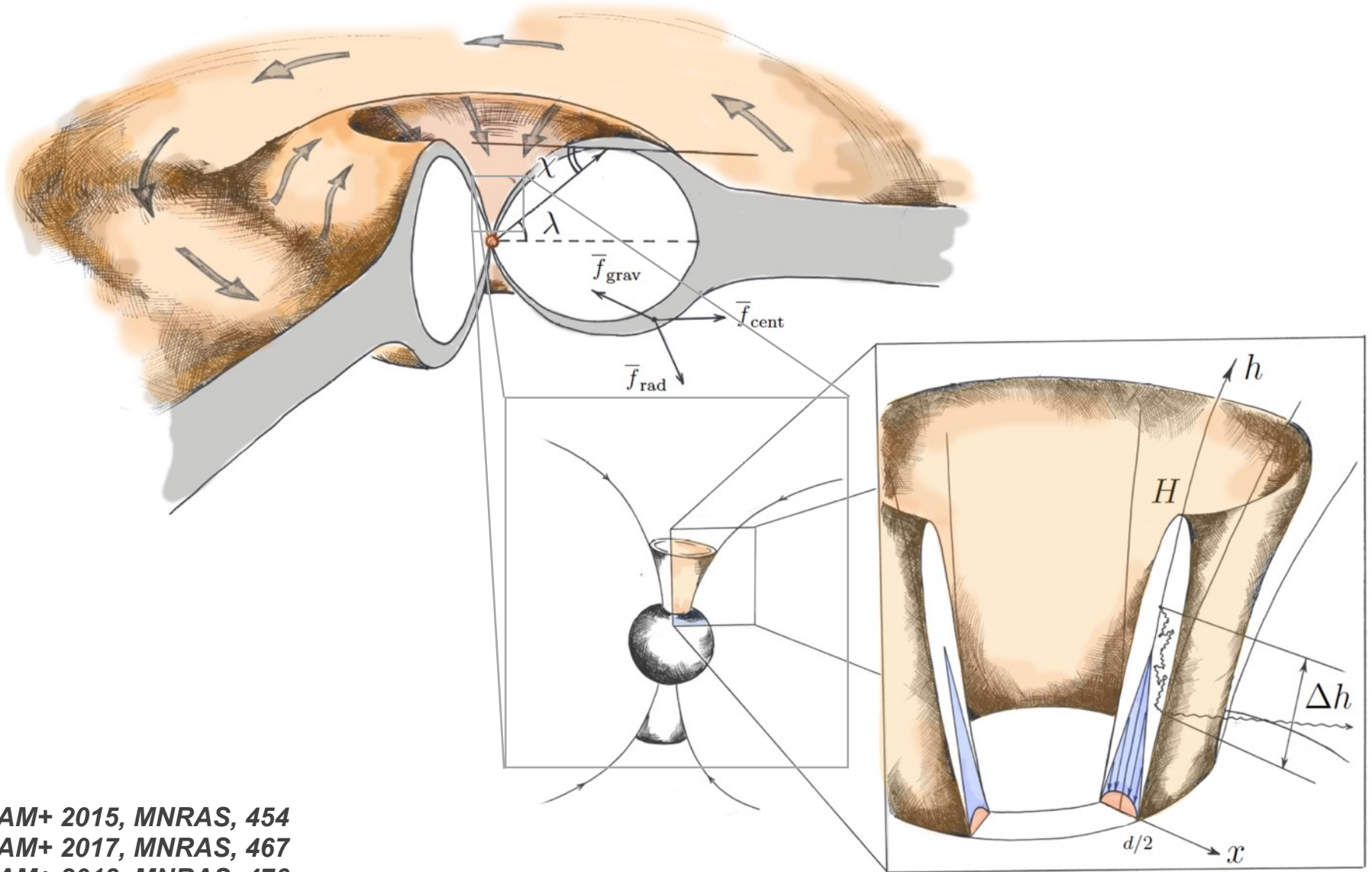
Some consequences:

- we hardly see the central NS directly
- multicolour black-body spectrum
- smooth pulse profiles
- super-orbital variability because of precession of magnetic dipole



Fuerst+ 2016

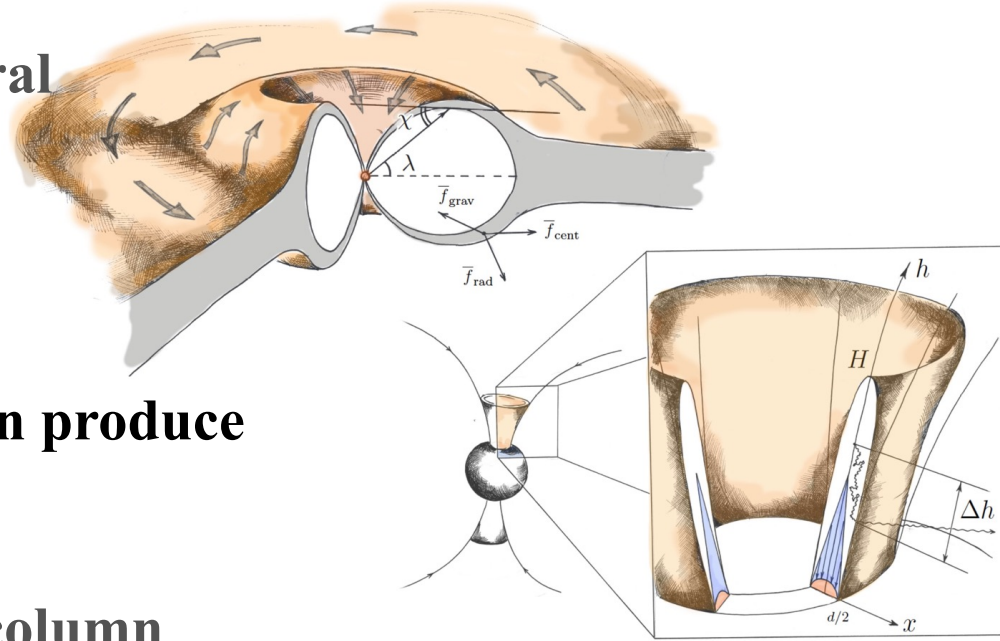
Geometry of ULXPs



AM+ 2015, MNRAS, 454
AM+ 2017, MNRAS, 467
AM+ 2018, MNRAS, 476

Conclusions

(1) Accretion columns are the central engines in ULXs; their luminosity is strongly affected by geometry (!)



(2) Advective accretion columns can produce strong neutrino emission

(3) The X-ray flux from accretion column is strongly lensed by NS gravitational field

(4) Bright ULXPs are surrounded by optically thick envelopes
- observational manifestation
- principal possibility of accretion

(5) We have to track bright transients to verify the theory of ULXPs