

Extreme Accretion onto Strongly Magnetized Neutron Stars

Alexander Mushtukov

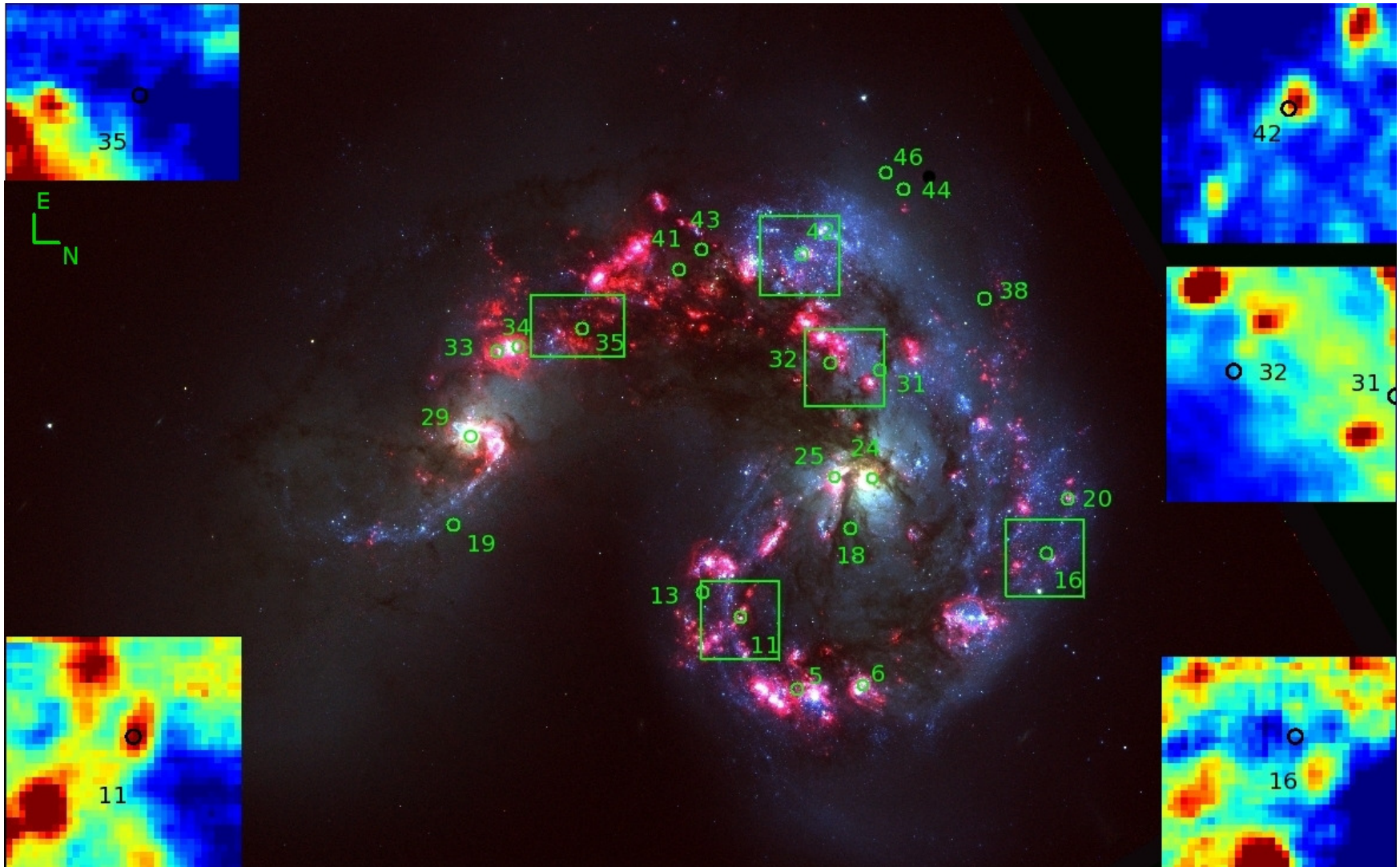
12 April 2021



**Universiteit
Leiden**

Ultraluminous X-ray Sources (ULXs)

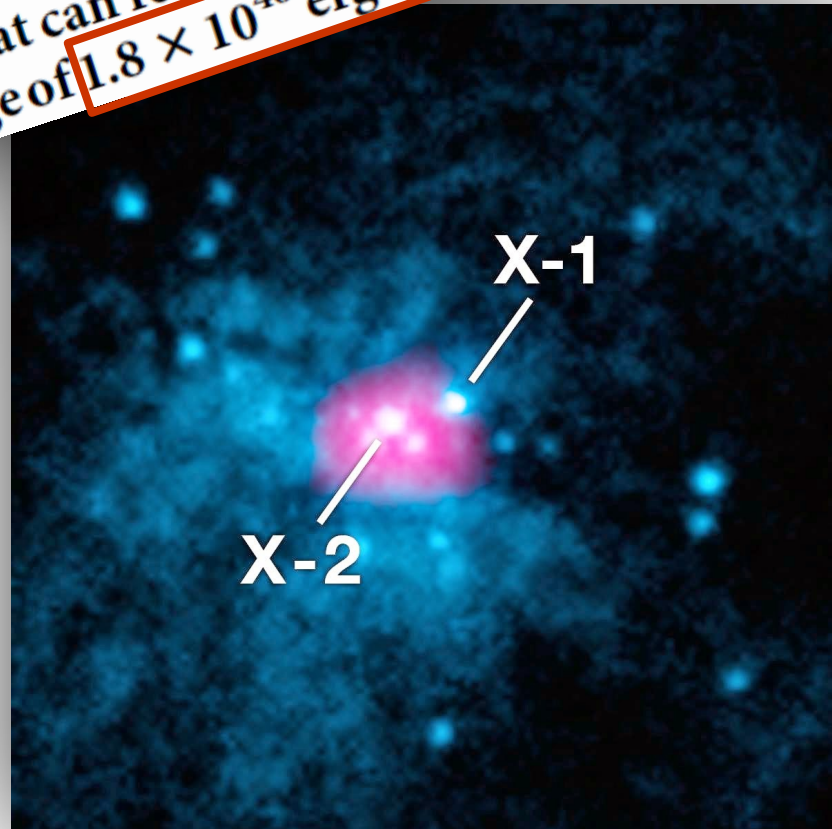
- Off-center bright X-ray sources in nearby galaxies
- Discovered with Einstein X-ray observatory 30 years ago
- X-ray luminosity: $L_x = 10^{39} - 10^{41} \text{ erg s}^{-1}$



Pulsations from ULX in M82

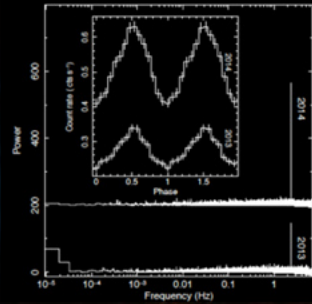
and the modulation arises from its binary orbit. The pulsed flux alone corresponds to an X-ray luminosity in the 3–30 kiloelectronvolt range of 4.9×10^{39} ergs per second. The pulsating source is spatially coincident with a variable source⁴ that can reach an X-ray luminosity in the 0.3–10 kiloelectronvolt range of 1.8×10^{40} ergs per second¹. This

**Pulse period
~1.37 sec**



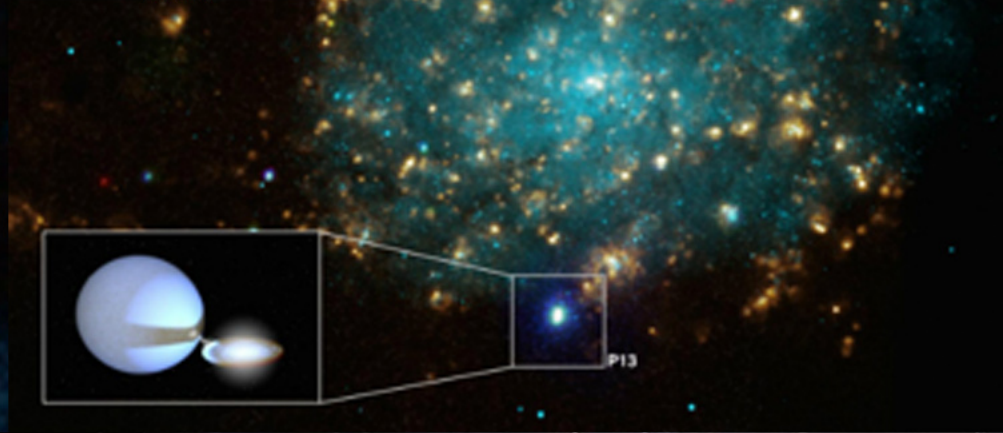
M82

$\sim 10^{40}$ erg s⁻¹



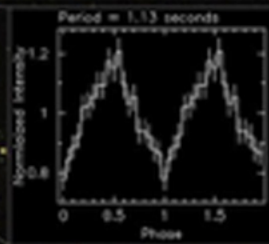
NGC 7793

$\sim 5 \cdot 10^{39}$ erg s⁻¹



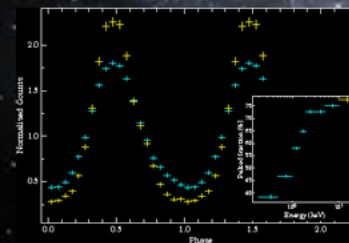
NGC 300

$\sim 5 \cdot 10^{39}$ erg s⁻¹



$\sim 2 \cdot 10^{41}$ erg s⁻¹

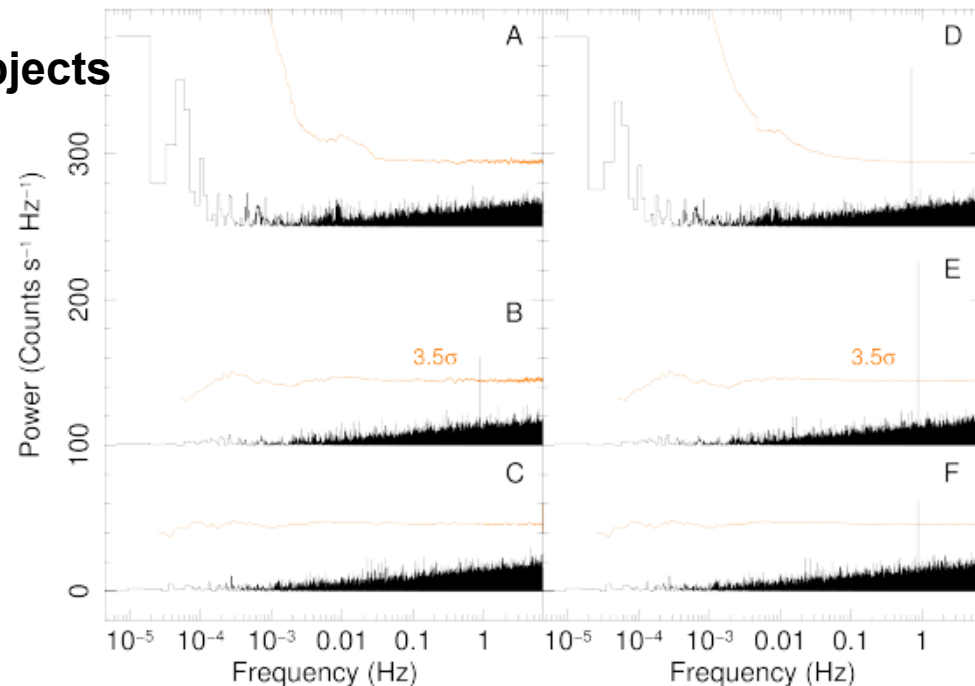
NGC 5907



ULX-pulsars in a nutshell

name	M82 ULX2	NGC 7793 P13	NGC5907 ULX1	NGC300 ULX1
L_X (max) [erg s ⁻¹]	1.8×10^{40}	5×10^{39}	10^{41}	4.7×10^{39}
P_s [s]	1.37	0.42	1.13	31.5
$\dot{\nu}$ [s ⁻²]	10^{-10}	4×10^{-11}	$4. \times 10^{-9}$	5.6×10^{-10}
P_{orb} [d]	2.52	64	5.3	
M_2 [M_\odot]	≥ 5.2	18 – 23		

- Large pulsed fraction, 20-30%, in all objects
- Smooth pulse profiles
- Huge variations in luminosity
- Multi-colour blackbody spectrum
- No (or very weak?) cyclotron lines



Bachetti+ 2014, Nature, 514

Israel+, 2017, Science, 355

Israel+, 2017, MNRAS, 466

Fürst+, 2016, ApJ, 831

Carpano+, 2018, MNRAS, 476

Neutron Stars

Product of supernova
explosions



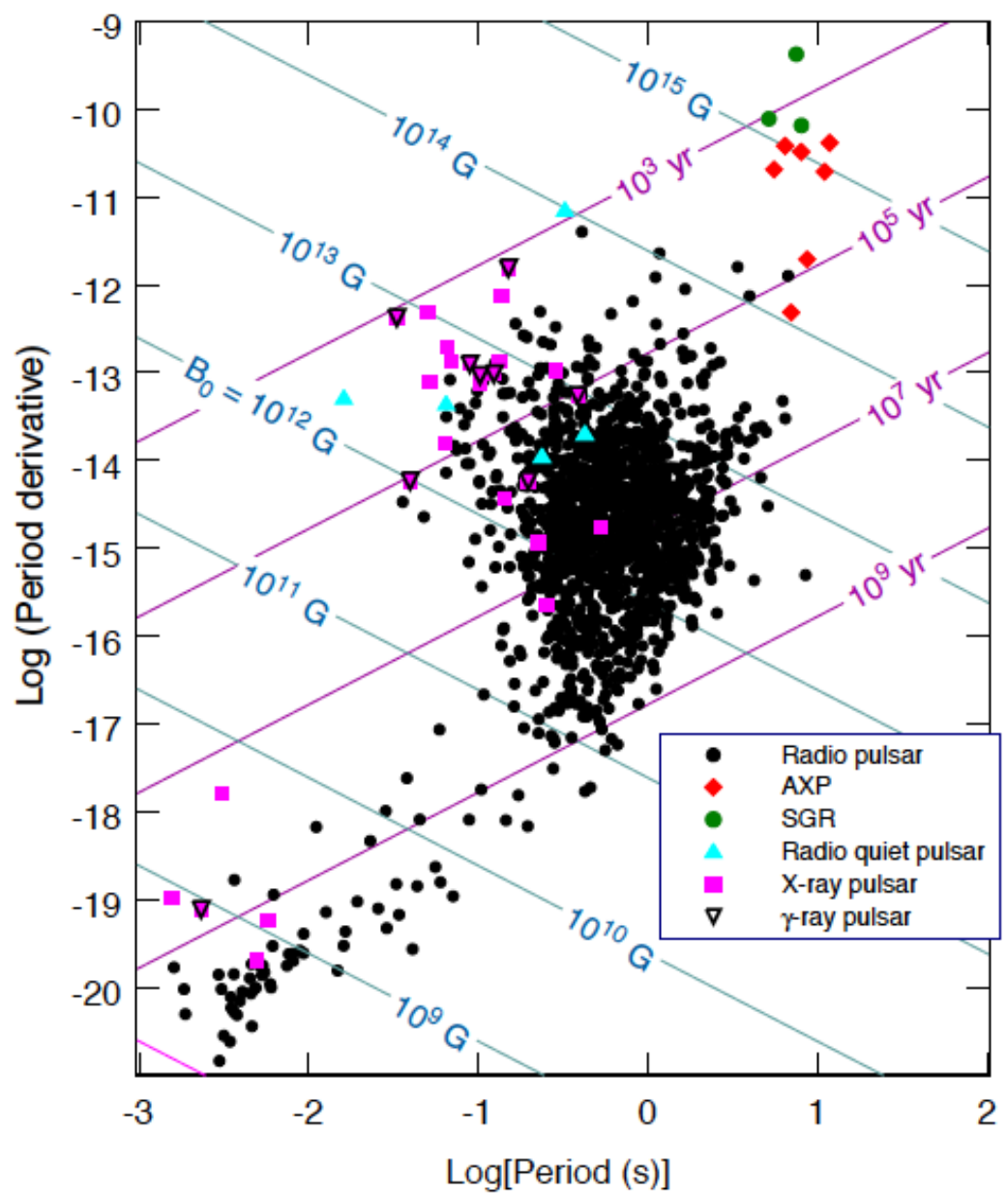
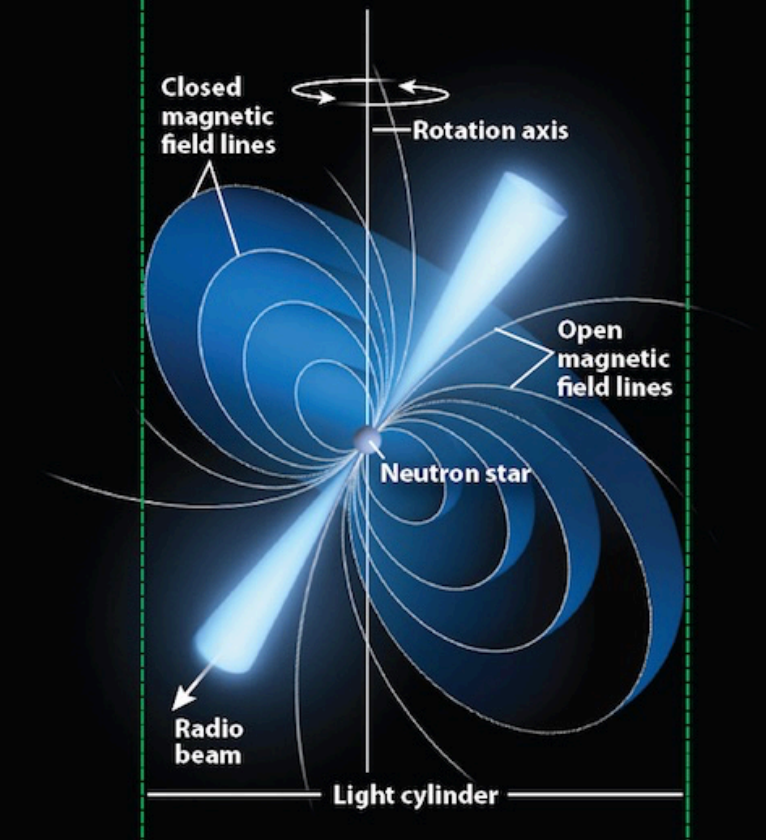
Mass: $1.5 M_{\text{sun}}$
Radius: 10 km

Mass density

$>10^{14} \text{ g cm}^{-3}$

**The Highest Density
in the Universe**

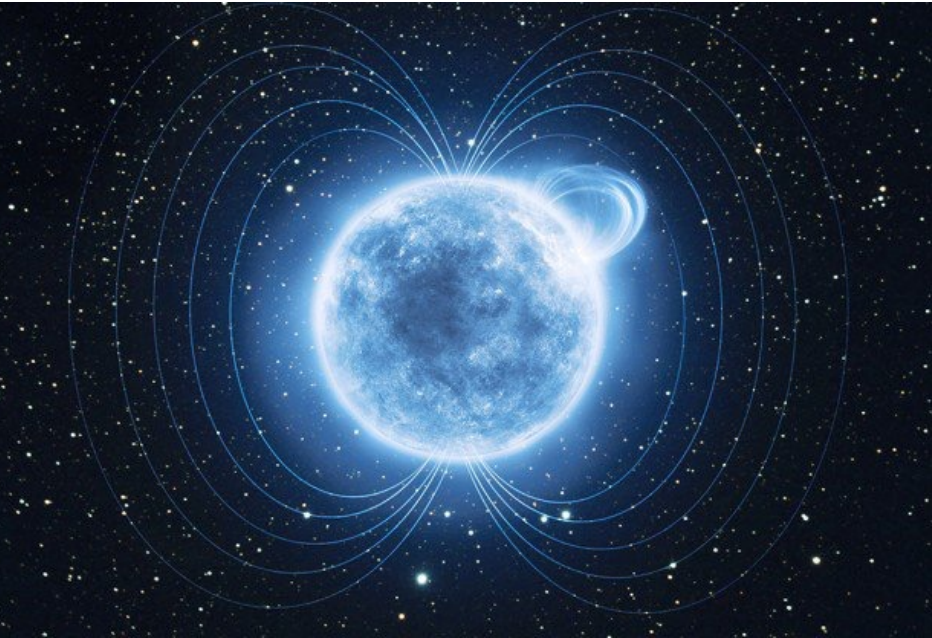
Magnetic fields



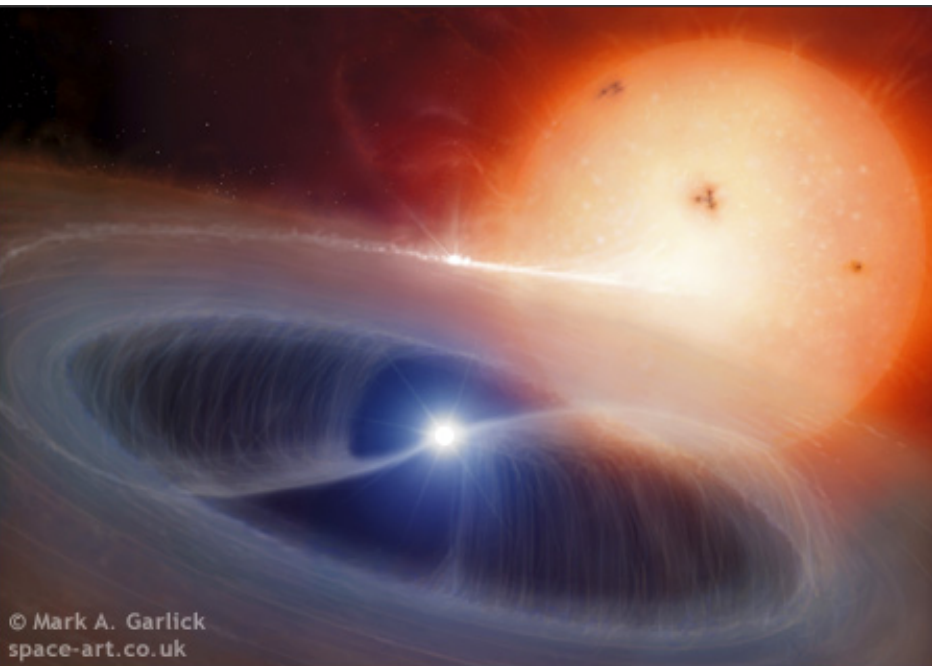
$$\dot{P}_{15} \equiv \dot{P} / (10^{-15} \text{ s s}^{-1})$$

$$B_s = \left(\frac{3Ic^3 P \dot{P}}{2\pi^2 R^6} \right)^{1/2} \simeq 2 \times 10^{12} \text{ G} (P \dot{P}_{15})^{1/2}$$

What if neutron star has a companion?

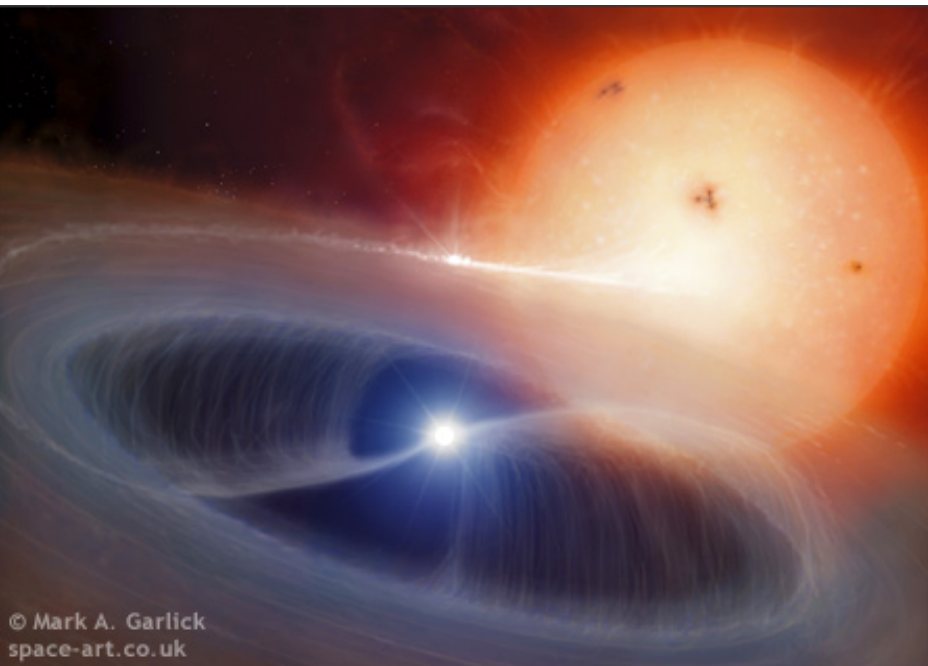
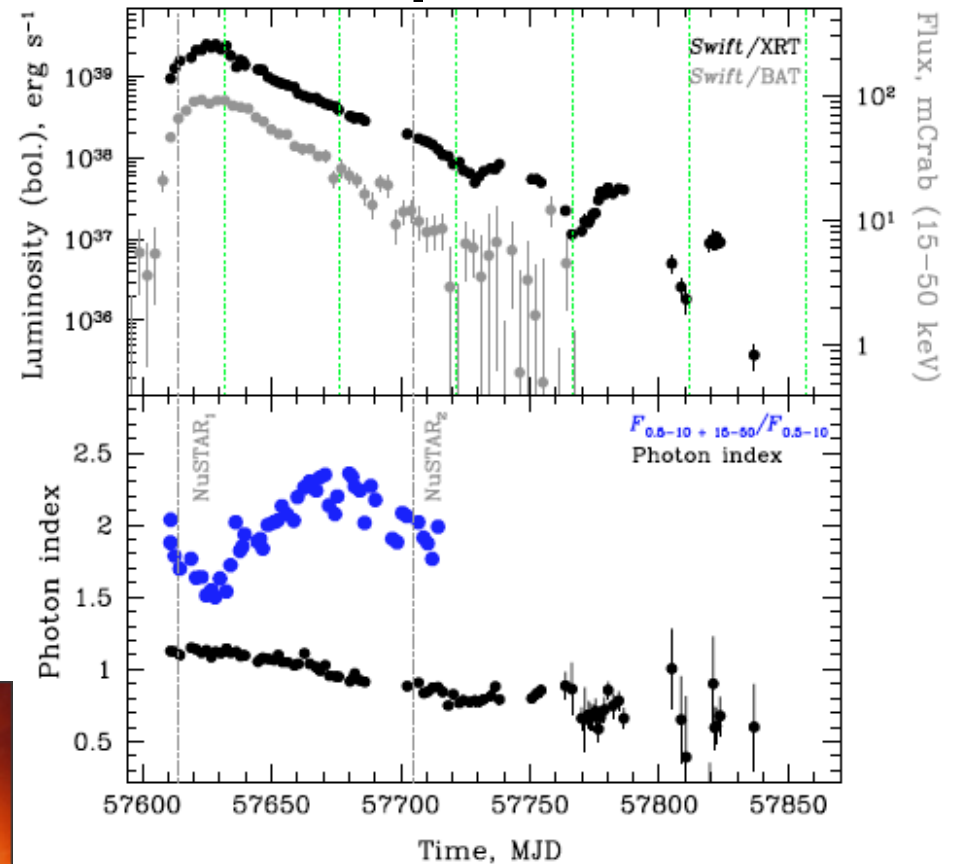


Some neutron stars are isolated and dim

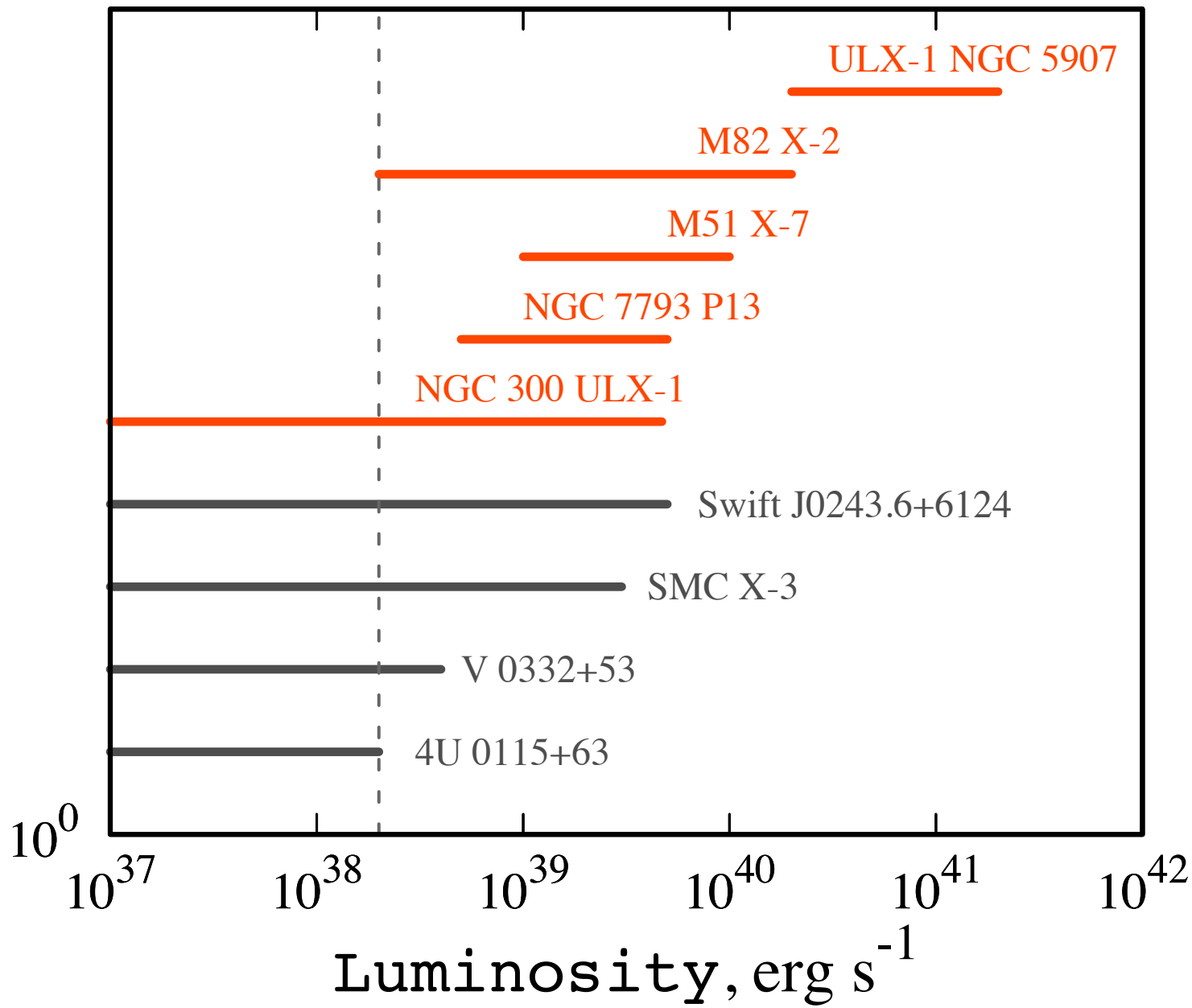


If neutron star has a close companion, it absorbs material and can be extremely bright.

What if neutron star has a companion?

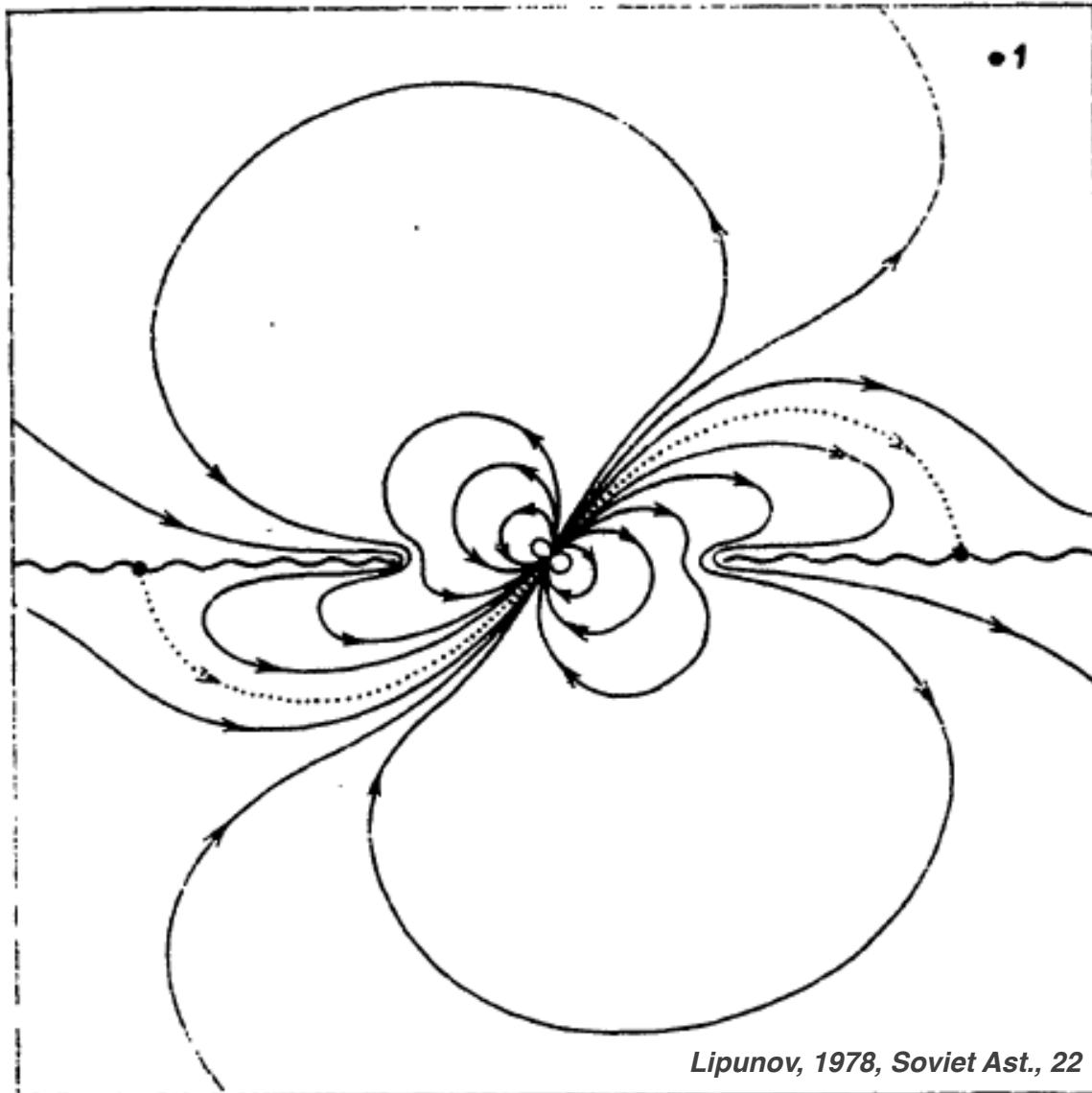


**If neutron star has a close companion,
it absorbs material
and
can be extremely bright.**



X-ray pulsar

Accretion Disc and its Interaction with B-field



Lipunov, 1978, Soviet Ast., 22

The inner disc radius:

$$r_A = \left(\frac{\mu^4}{2GM\dot{M}^2} \right)^{1/7}$$

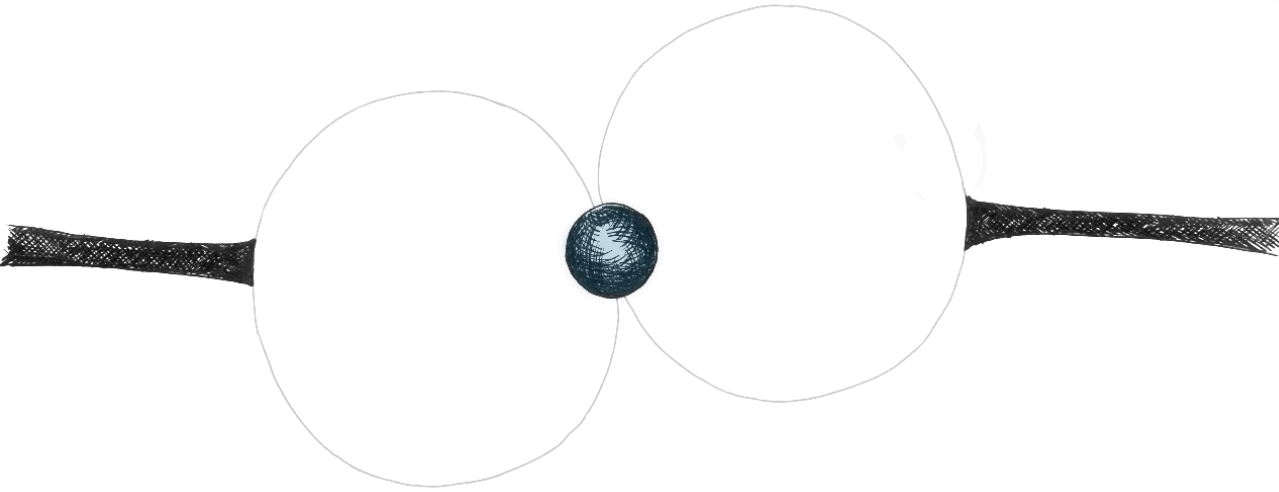
$$r_m = \xi r_A$$

Co-rotational radius:

$$r_{co} = \left(\frac{GM}{\Omega^2} \right)^{1/3}$$

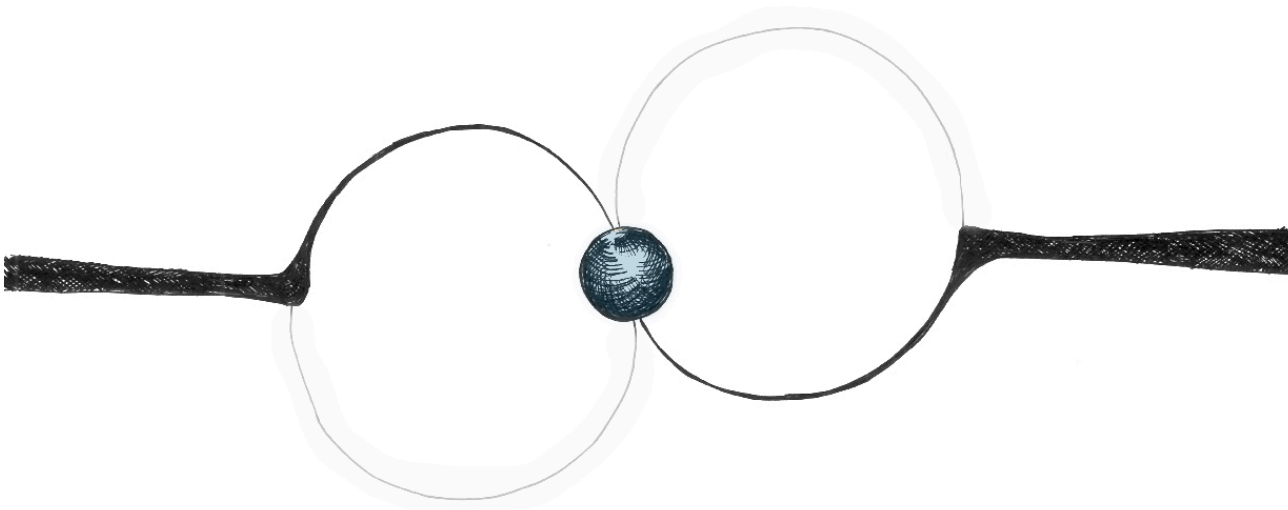
Keplerian and stellar-rotation frequencies are equal

“Propeller” state



$r_m > r_{co}$
accretion is prohibited due
to centrifugal barrier

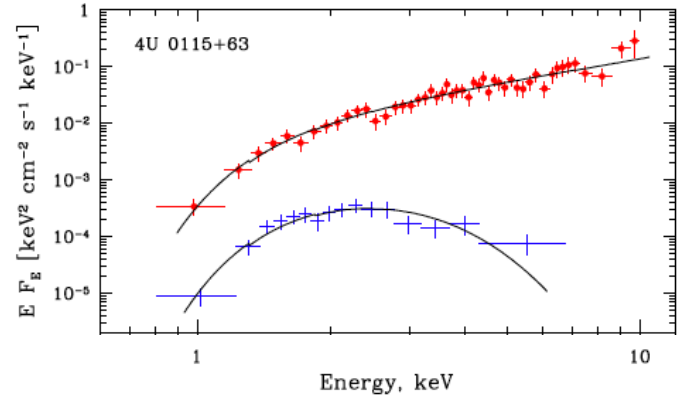
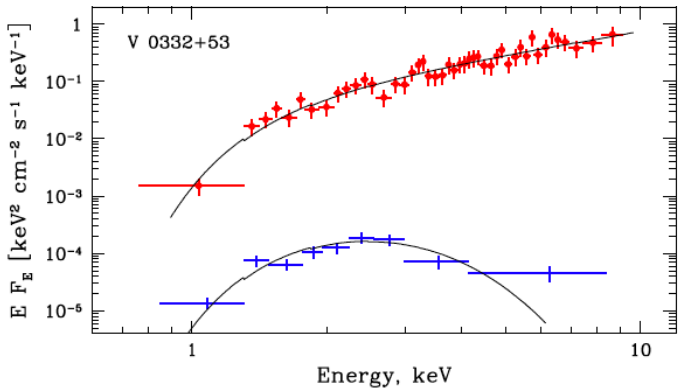
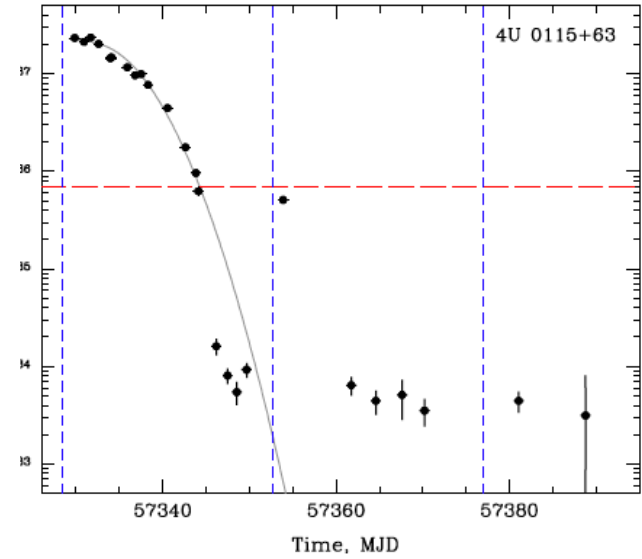
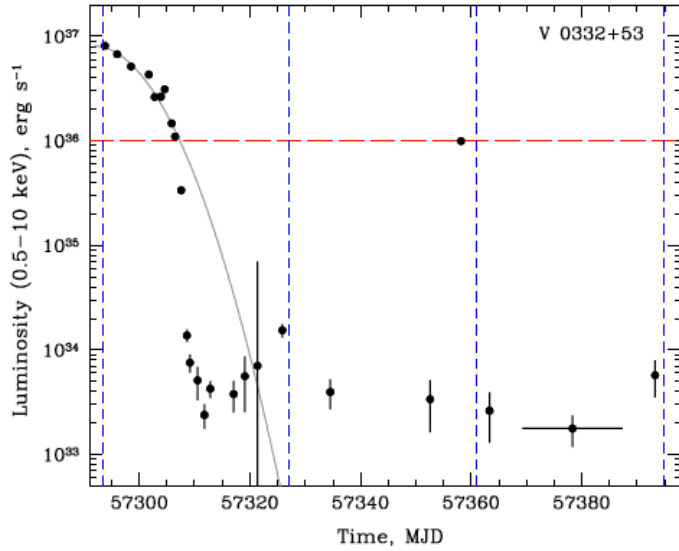
Accretion state



$r_m < r_{co}$
accretion is possible

“Propeller” effect

Detection

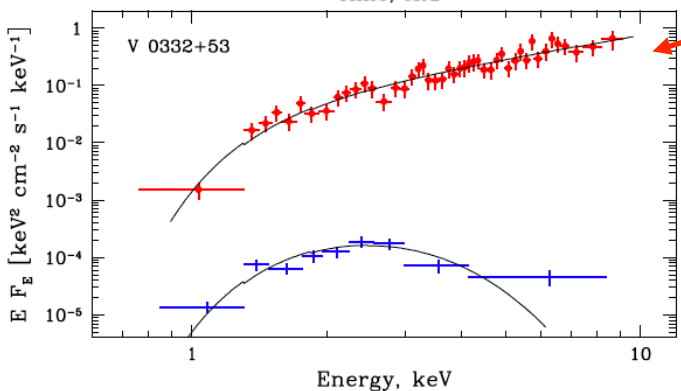
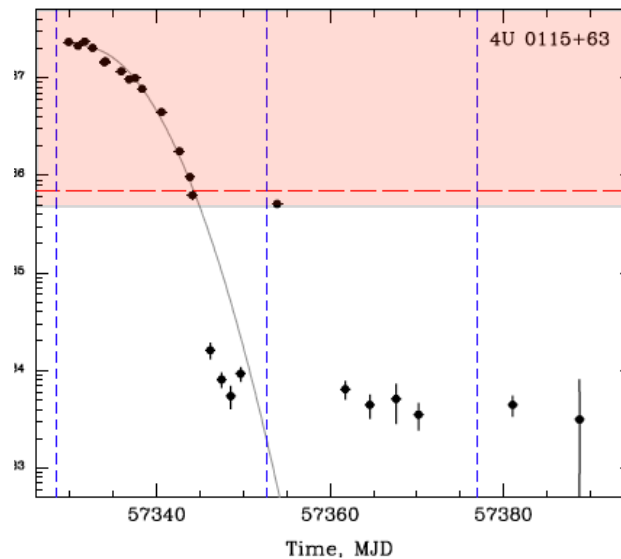
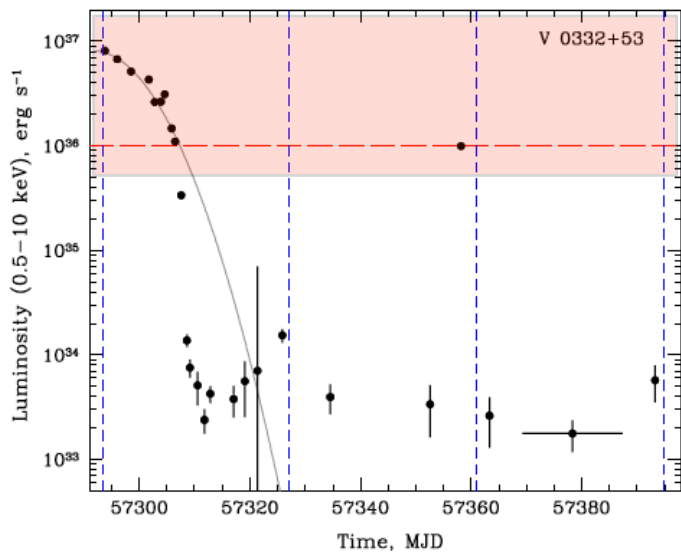


Propeller luminosity:

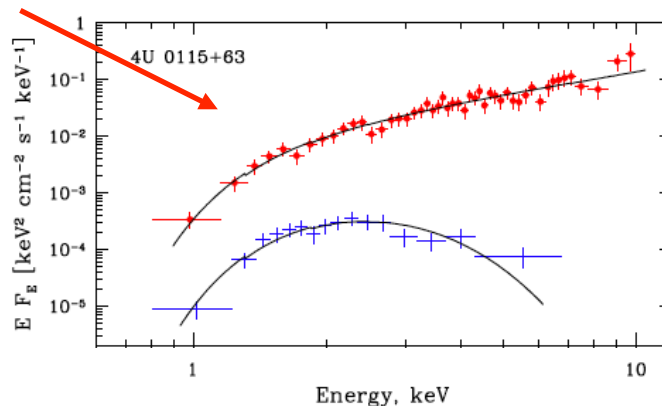
$$L_{\text{prop}} \approx 3.5 \times 10^{36} B_{12}^2 P^{-7/3} M_{1.4}^{-2/3} R_6^5 \text{ erg s}^{-1}$$

“Propeller” effect

Detection



Absorbed power-law

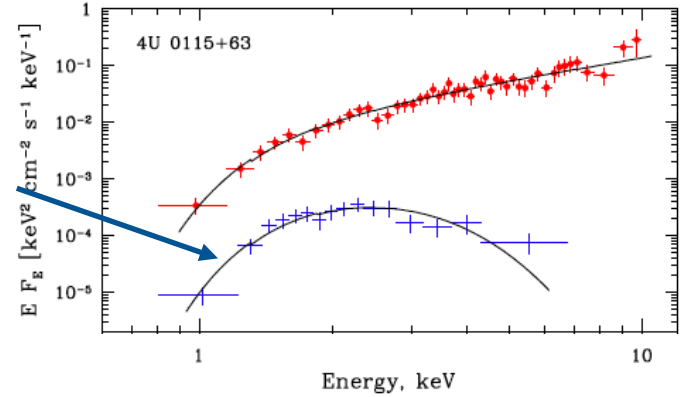
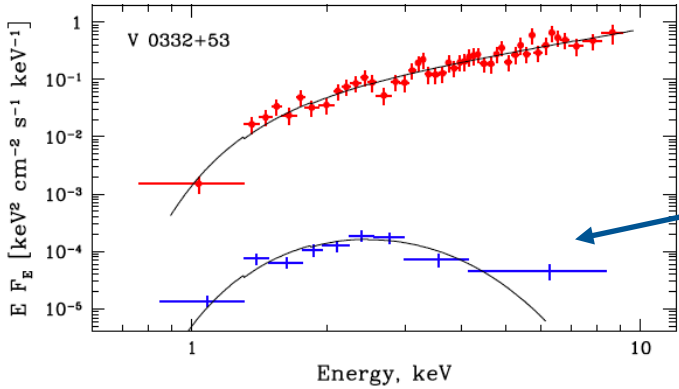
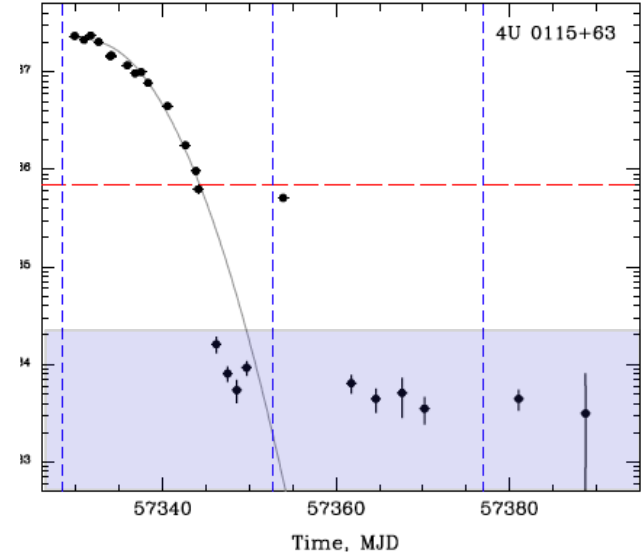
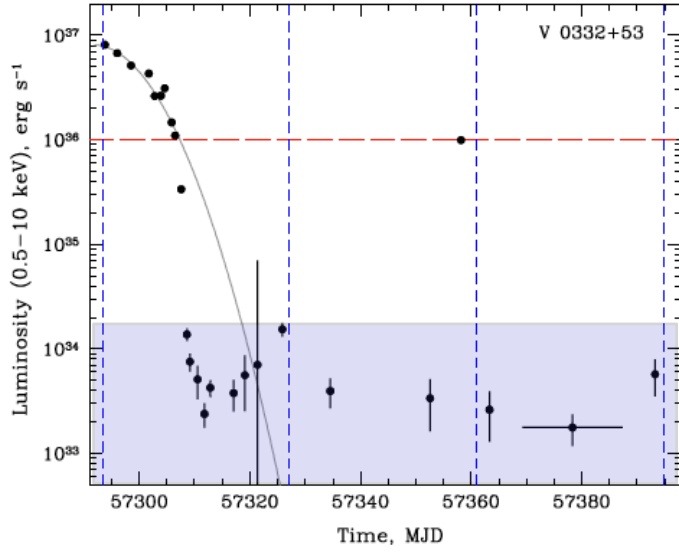


Propeller luminosity:

$$L_{\text{prop}} \approx 3.5 \times 10^{36} B_{12}^2 P^{-7/3} M_{1.4}^{-2/3} R_6^5 \text{ erg s}^{-1}$$

“Propeller” effect

Detection

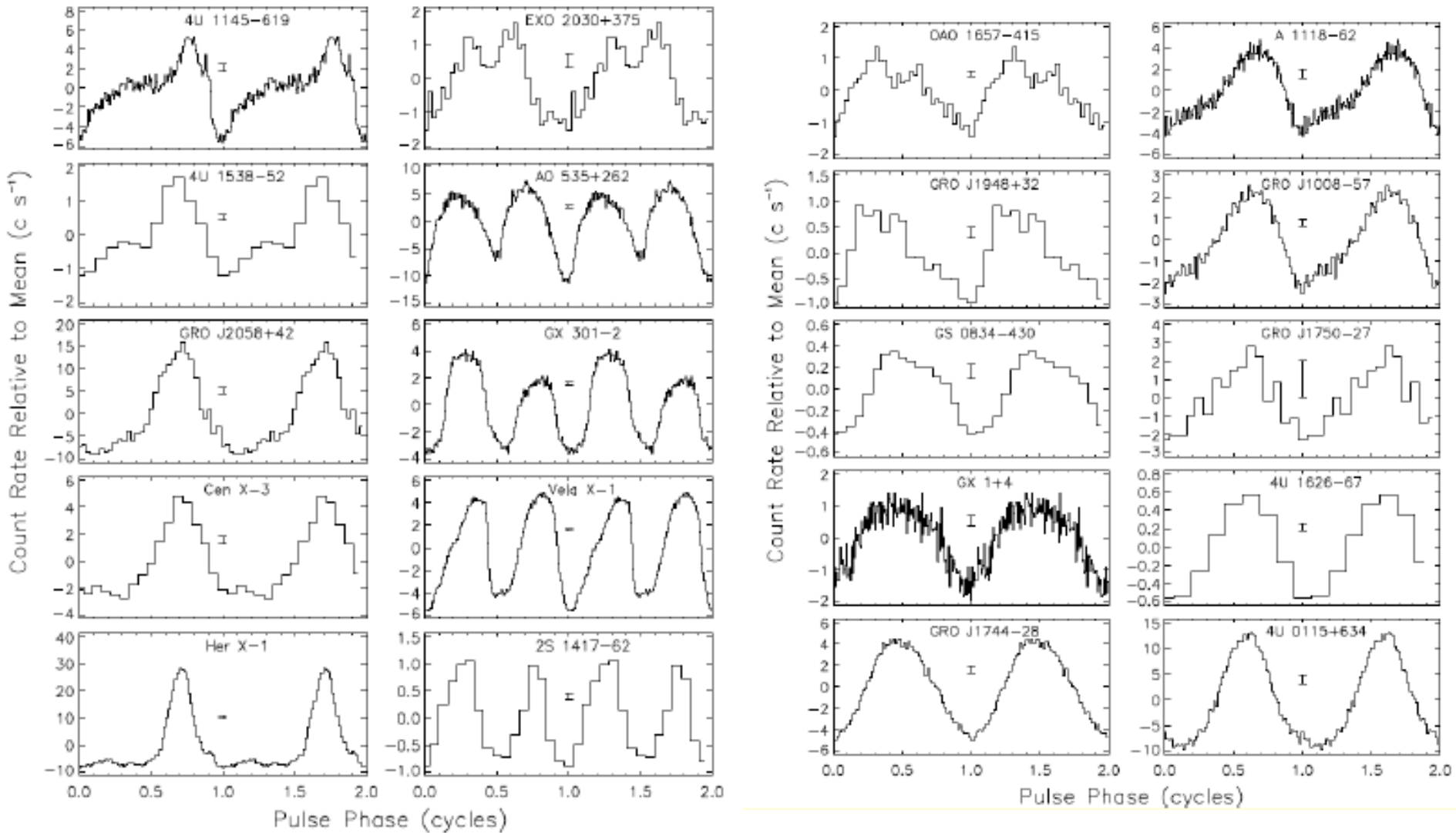


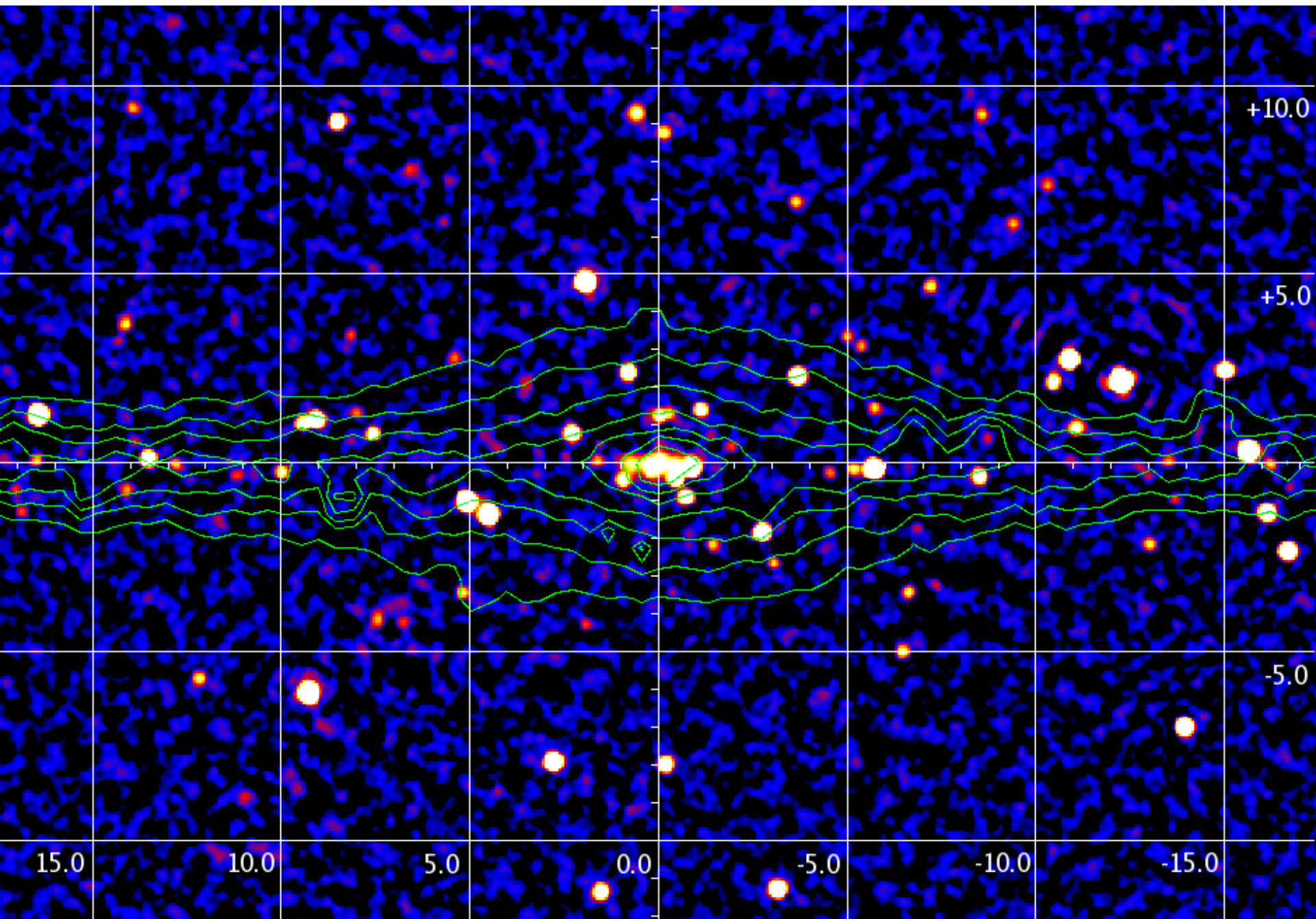
Black body
with
 $T=0.5$ keV

Propeller luminosity:

$$L_{\text{prop}} \approx 3.5 \times 10^{36} B_{12}^2 P^{-7/3} M_{1.4}^{-2/3} R_6^5 \text{ erg s}^{-1}$$

X-ray pulsars: pulse profiles





Magnetic field strengths



Magnetic field strengths



Magnetic field strengths

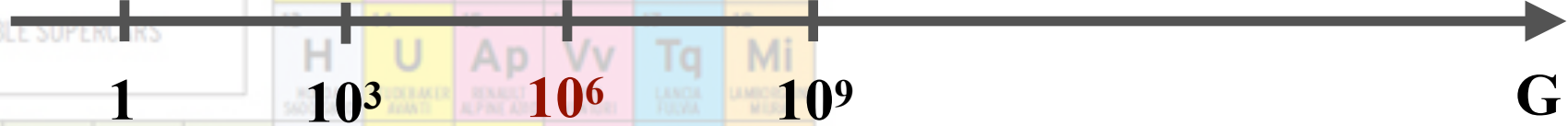


Magnetic field strengths

- POSEUR-OIDS
- FRANCO-OIDS
- HALOGEN
- NOBLE SUPERSTARS

Earth

Stars
Strongest field
in lab



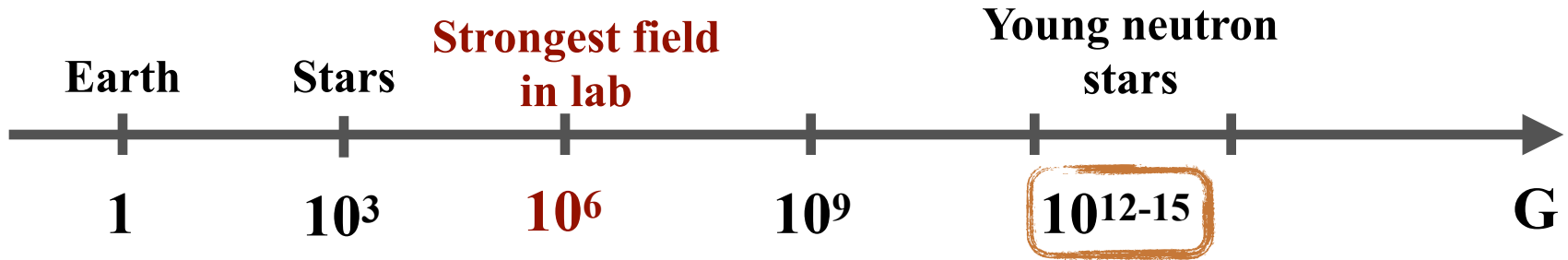
27 Wo BENTLEY SPEED SIX	28 Mg MG 3-SERIES	29 Bm BMW 328	30 Ss JAGUAR XK000	31 T TOYOTA 2000GT	32 Ca CADILLAC ALLANTE	33 Br BUICK REATTA	34 Re RENAULT SPORT	35 Hf LANCIA STRATOS	36 Ff FERRARI
45 Lm LAMBORGHINI MICHELANGELO	46 Bf BRISTOL FIGHTER	47 Sl MERCEDES SLS	48 Mc MCLAREN MP4-02C	49 Da DATSUN 2000	50 Z NISSAN Z	51 Tc CHRYSLER PT CRUZER	52 Hr PLYMOUTH Prowler	53 La LANCIA RALLY 037	54 Cf MCLAREN F1
77 F PONTIAC FIREBIRD	78 W WOLSKEL CONSULET	79 Lz DODGE Viper	80 Le LOTUS ELISE	81 Rx MAZDA RX-7	82 Mr TOYOTA MR2	83 Mx MAZDA MAZ3	84 XI CADILLAC ELR	85 Q AUDI QUATTRO	86 Ez FERRARI ENZO
109 Dt DE TOMASO PANTEGRA	110 Ma MORGAN AERO 8	111 Sk SPYKER C8	112 Ev LOTUS EVORA	113 Su TOYOTA SUPRA	114 Ax ACURA NSX	115 Ho HONDA SC2000	116 Lx LEXUS LFA	117 Rs JORD GORDON	118 Ru BUGATTI VEYRON

the end of school chemistry

63 Fx FERRARI F119	64 Bb FERRARI SC288	65 Pi FERRARI 308	66 Lj LAMBORGHINI JAURA	67 Au FERRARI TESTAROSSA	68 Cv FERRARI F355	69 X FERRARI F50	70 Mo FERRARI 360	71 It FERRARI 458
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95 Cc Koenigsegg CC	96 Pz PAGANI ZONDA	97 G Gumpert Apollo	98 Ct CARRARO TI	99 Ua ULTIMA GT-R	100 Bs Mazda's G-Force	101 Fu Suzuki's Ultimate Aero	102 Wm Wolker W7000	103 N Saleen S7
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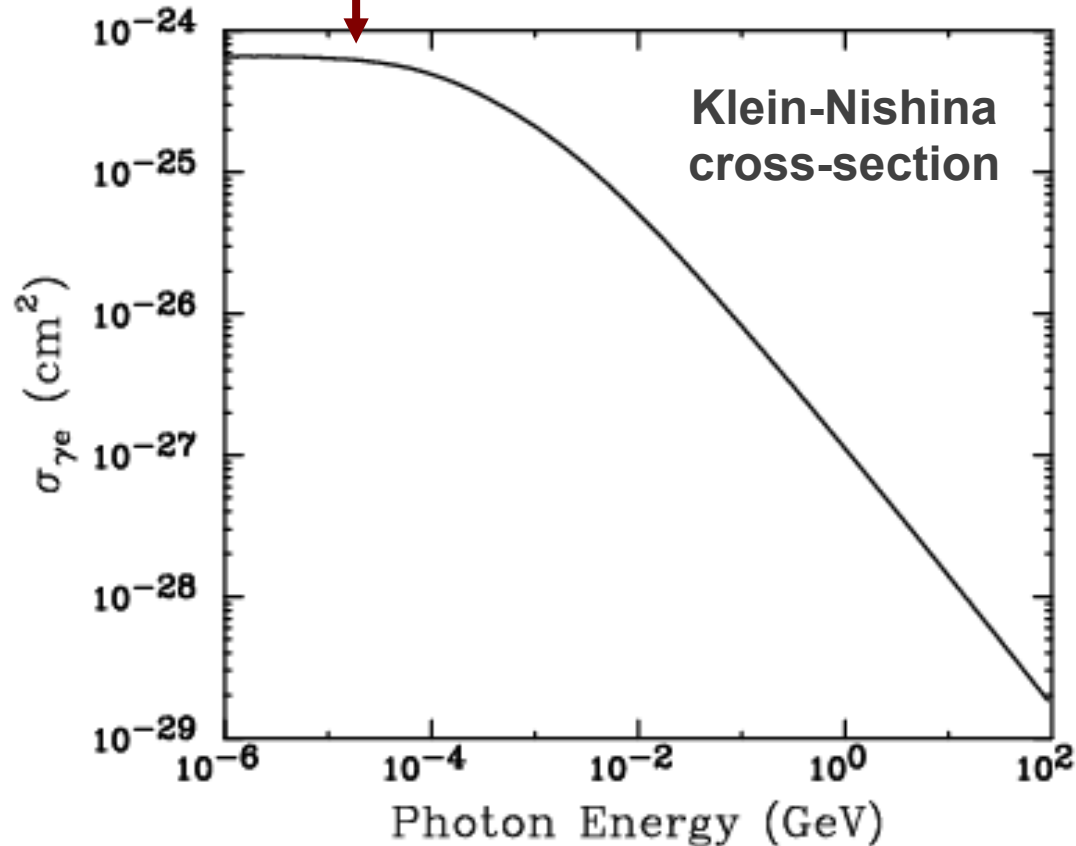
Magnetic field strengths



Extreme physics
Deviations from
Quantum
Electrodynamics?

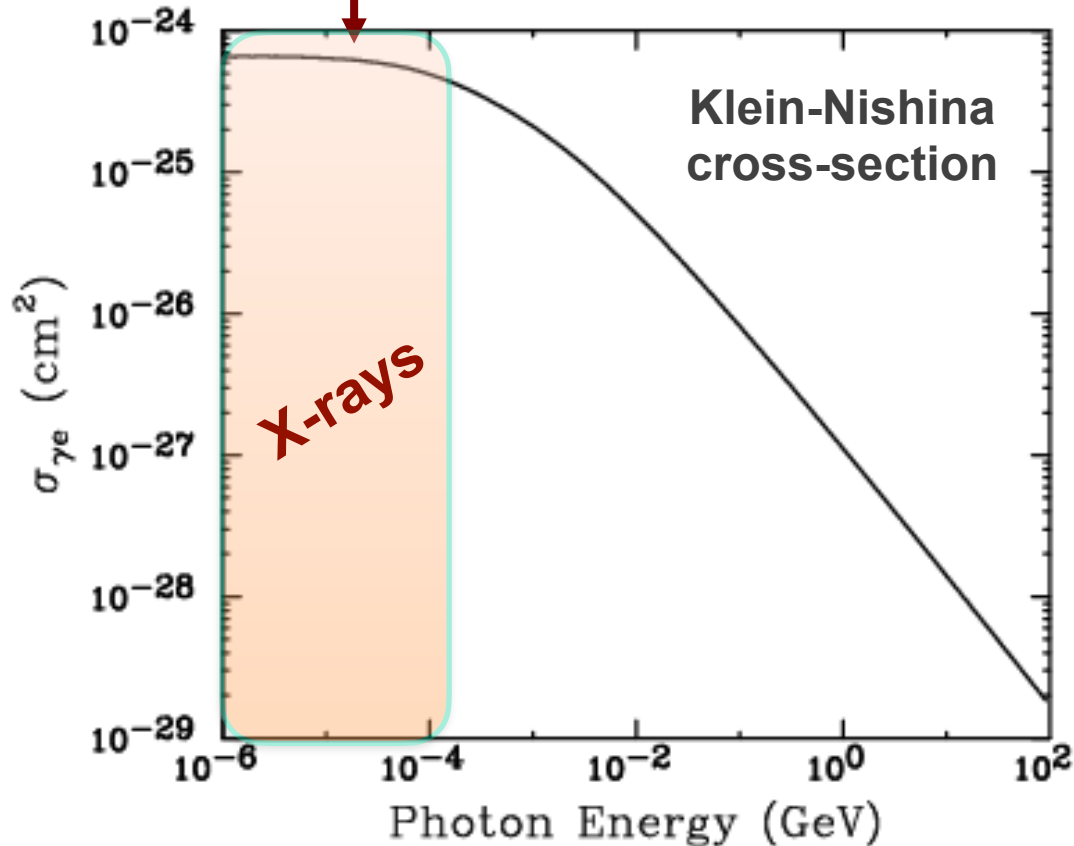
Compton scattering: non-magnetic case

$$\sigma_T = \frac{8\pi}{3} \left(\frac{e^2}{m_e c^2} \right)^2 \approx 6.65 \times 10^{-25} \text{ cm}^2$$



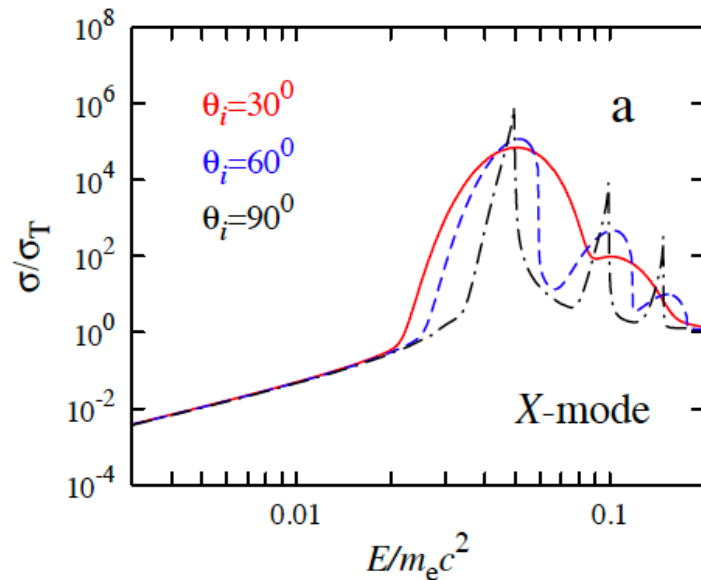
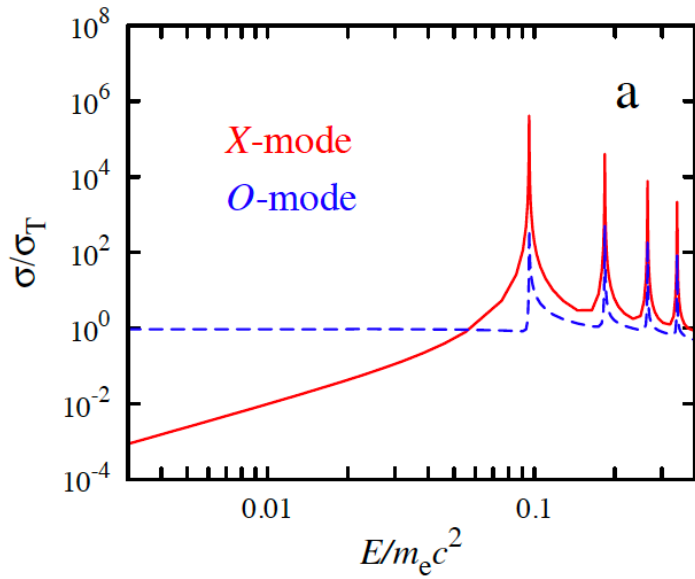
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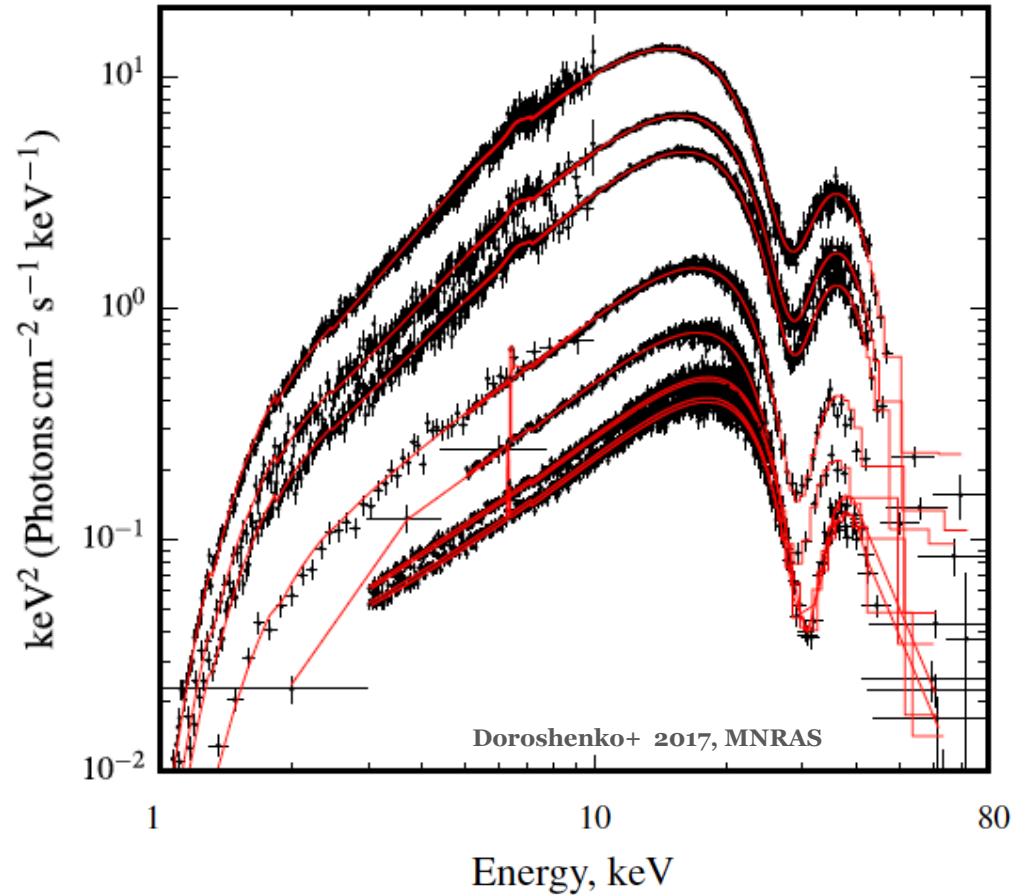


X-ray pulsar

Typical spectra



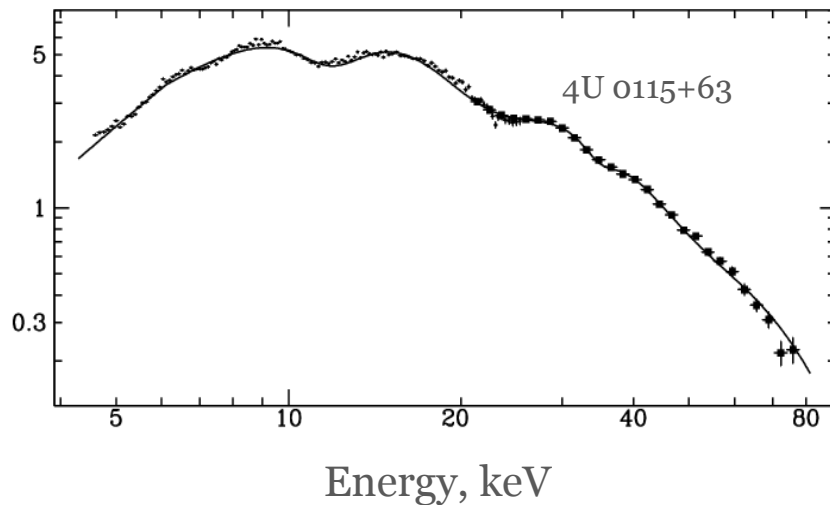
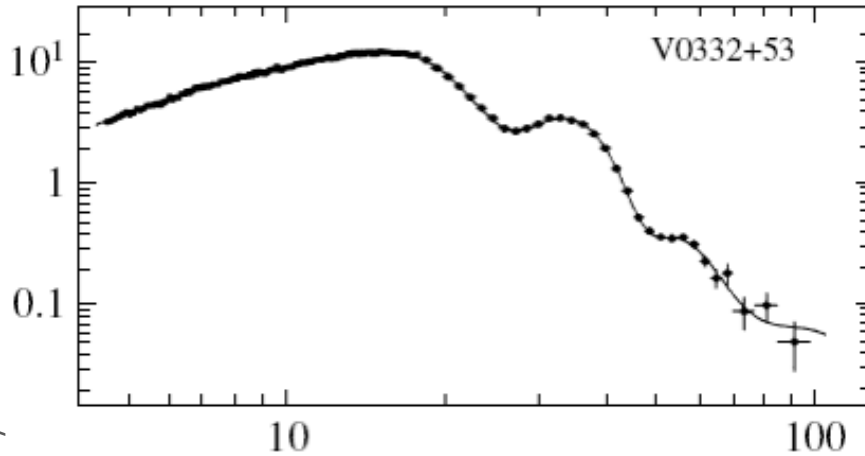
Spectra of V 0332+53 at different luminosity states



$$E_{\text{cyc}} = 11.6 B_{12} \text{ keV}$$

X-ray pulsar

Typical spectra

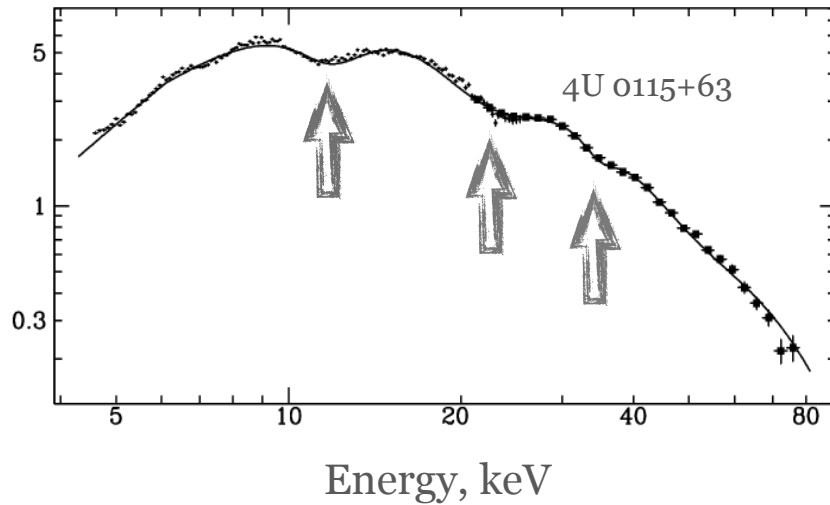
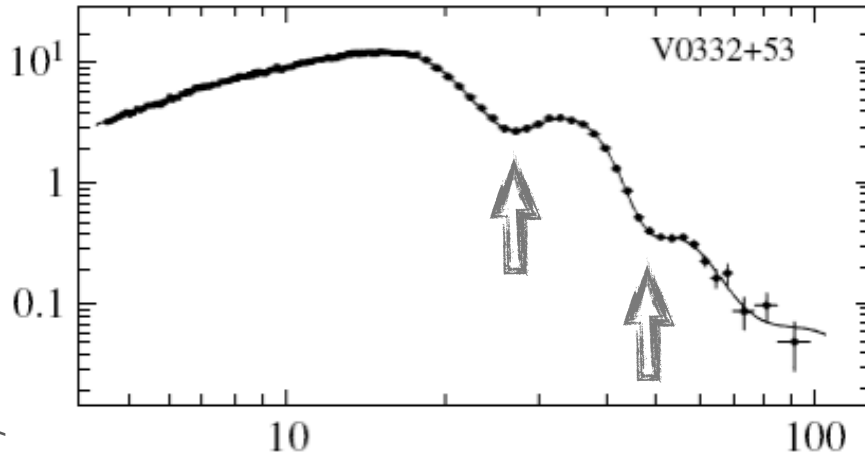


$$E_{\text{cyc}} = 11.6 B_{12} \text{ keV}$$

Source name	Cyclotron energy, keV
4U 0115+63 (–)	11.5, 20.1, 33.6, 49.5, 53
V 0332+53 (–)	28, 53, 74
4U 0352+309 (X Per)	29
RX J0440.9+4431	32
RX J0520.5-6932	31.5
A 0535+262	50, 110
MXB 0656–072	36
Vela X-1 (+)	27, 54
GRO J1008-57	88 [?] , 75.5
1A 1118–61	55
Cen X-3	28
GX 301–2	37, 48
GX 304–1 (+)	50.8
4U 1538–52	20, 47
Swift J1626.6–5156	10
4U 1626–67	37
Her X-1 (+)	42
OA0 1657–415	36
GRO J1744–28	4.7
IGR J18179–1621	21
GS 1843+00	20
4U 1907+09	19, 40
4U 1909+07	44 [?]
XTE J1946+274	36
KS 1947+300	12.5
EXO 2030+375	11 [?] , 36 [?] , 63 [?]
Cep X-4	30

X-ray pulsar

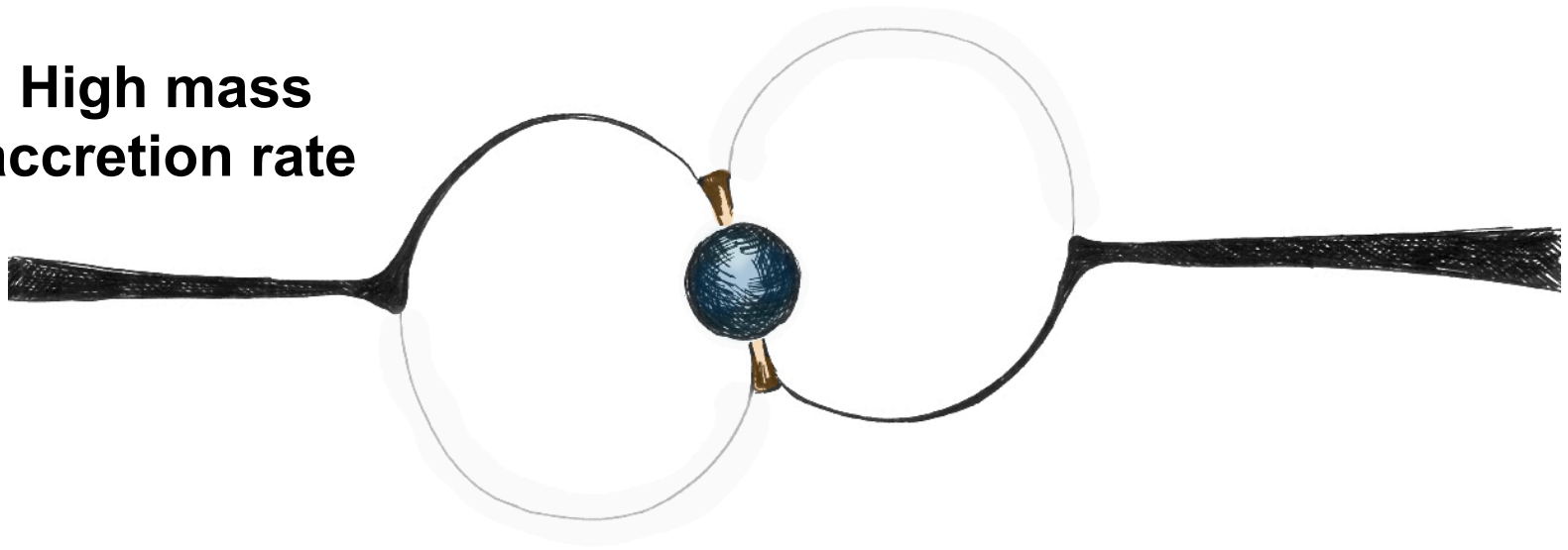
Typical spectra



$$E_{\text{cyc}} = 11.6 B_{12} \text{ keV}$$

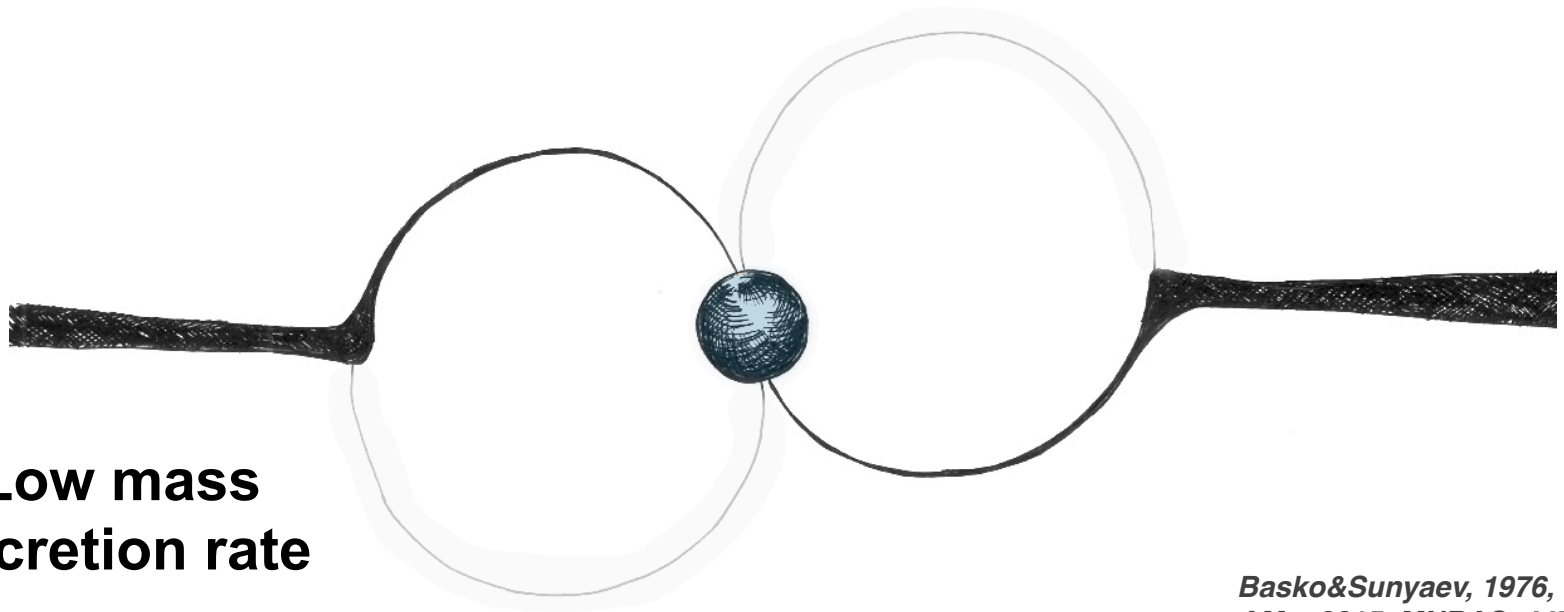
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**High mass
accretion rate**

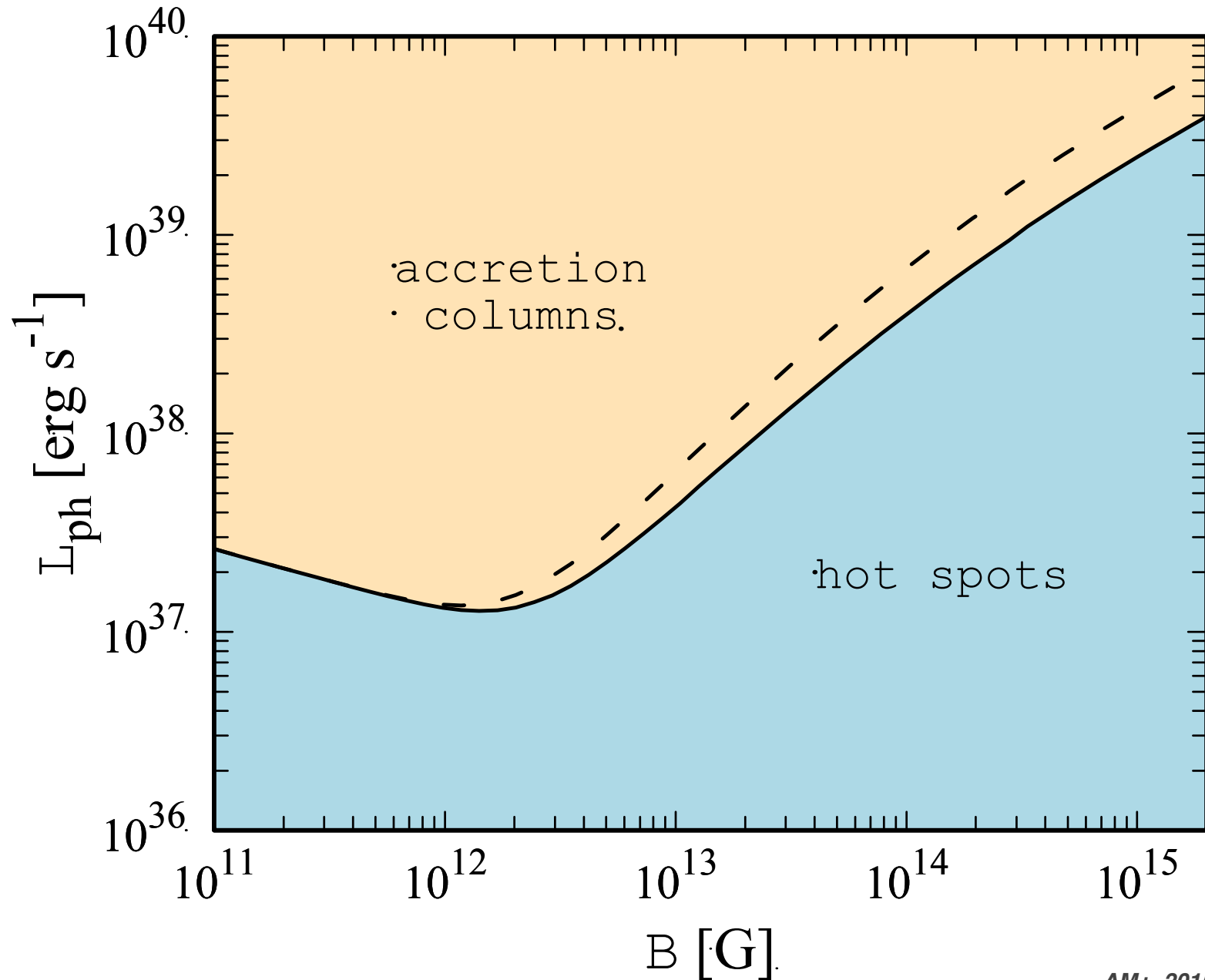


Critical luminosity

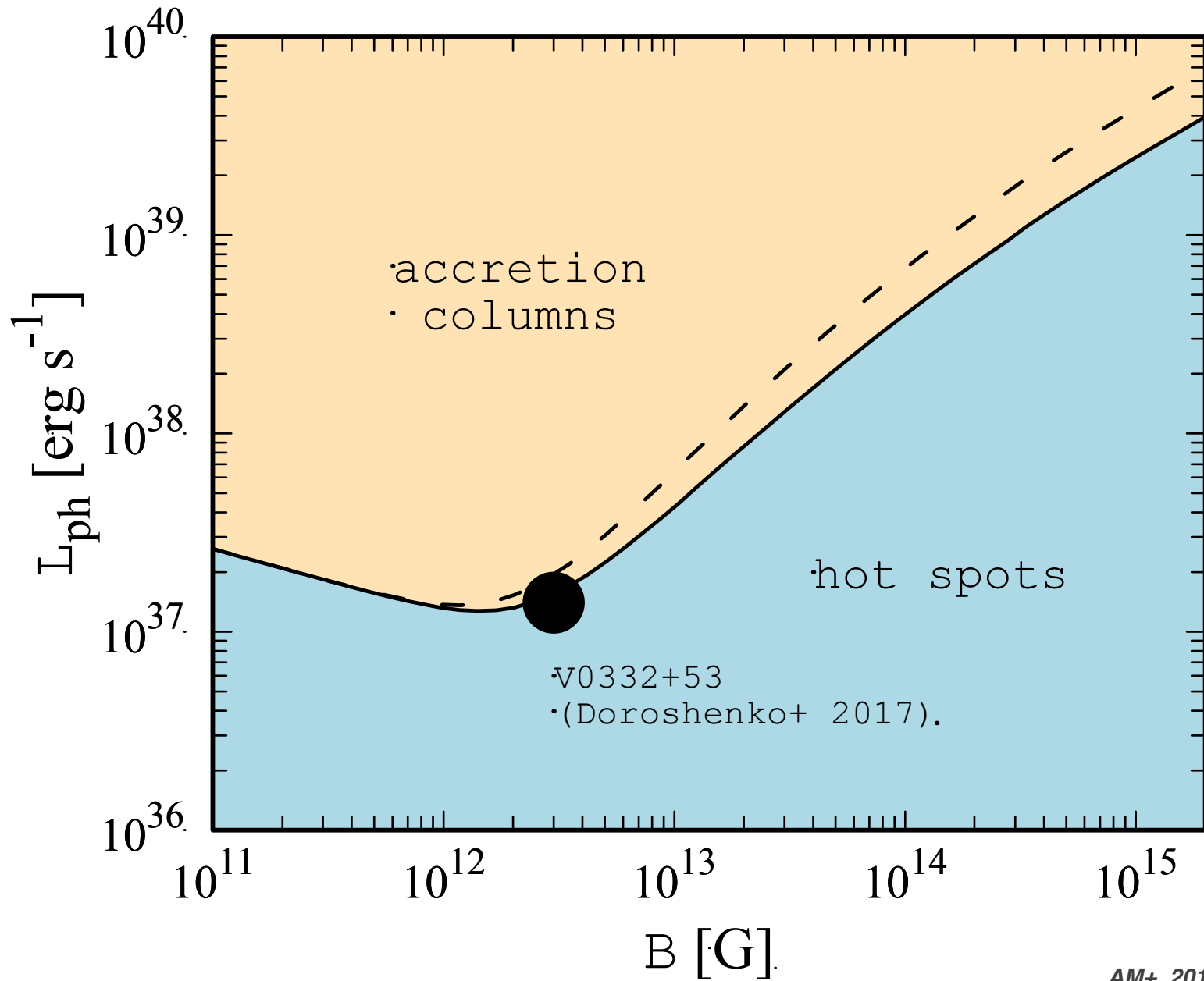
**Low mass
accretion rate**



Critical luminosity



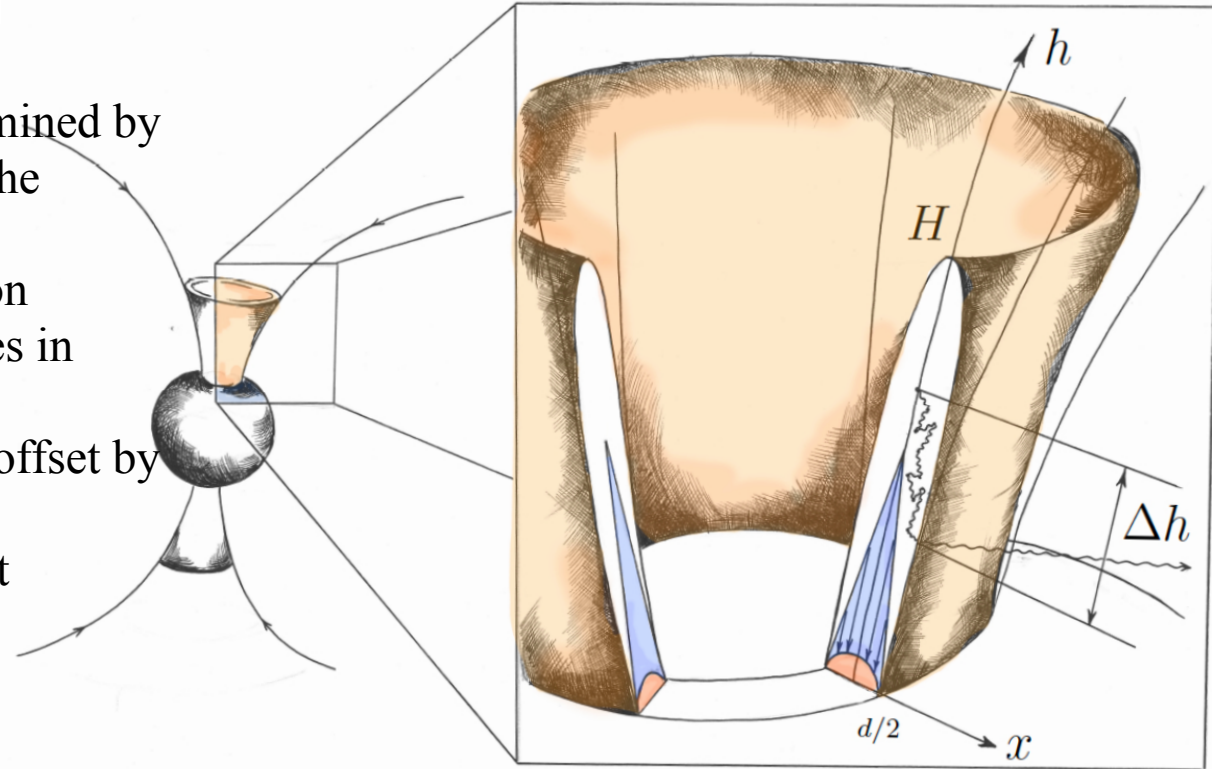
Critical luminosity



Above the critical luminosity: accretion column

A set of assumptions:

- (1) dipole magnetic field;
- (2) geometrical thickness is determined by the thickness of accretion disc at the magnetospheric radius;
- (3) accretion flow stops at radiation dominated shock and slowly settles in inside a sinking region
- (4) the gravitational force will be offset by the radiation pressure gradient
- (5) the gas pressure is unimportant



$$L(H = R) \approx 1.8 \times 10^{39} \left(\frac{l_0/d_0}{50} \right) \left(\frac{\kappa_T}{\kappa_{\perp}} \right) \frac{M}{M_{\odot}} \text{erg s}^{-1}$$

Stable accretion columns cannot be infinitely bright

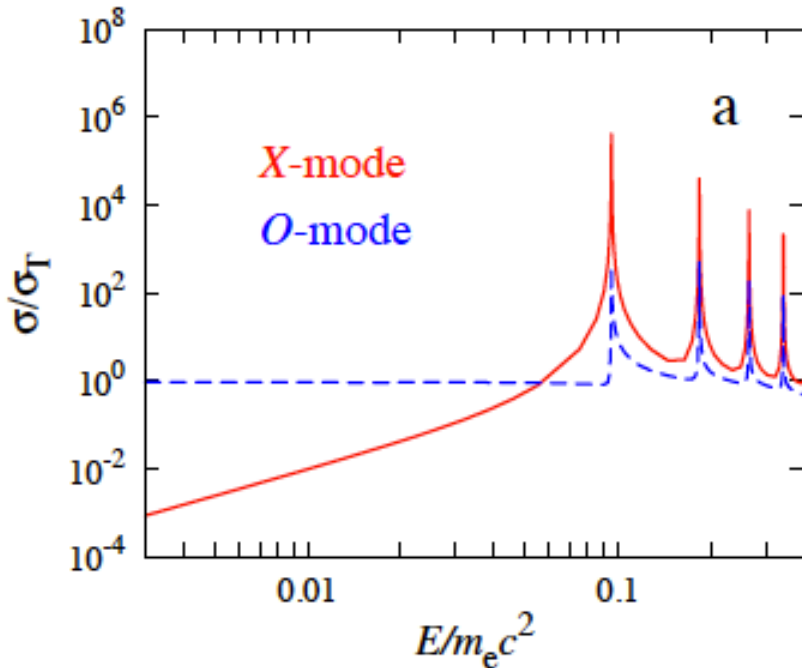
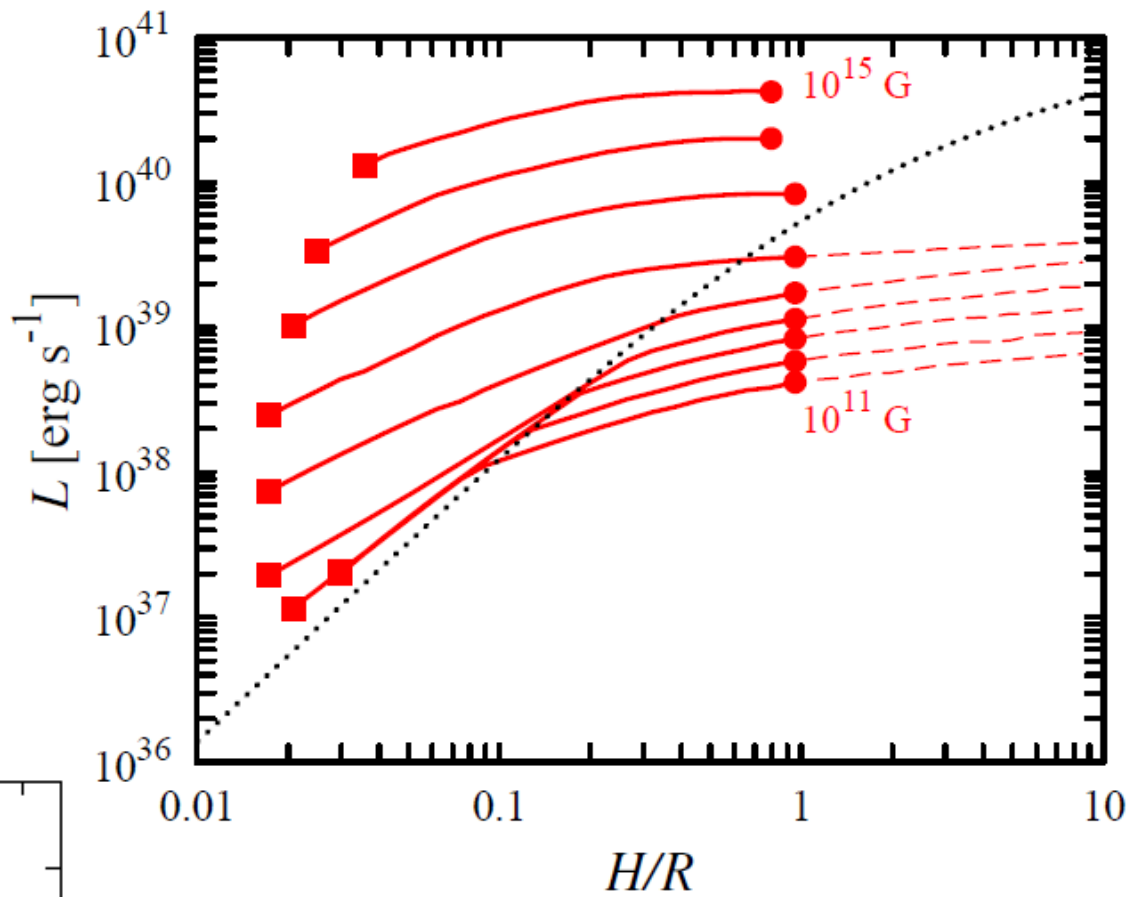
The reason:

The higher the mass accretion rate

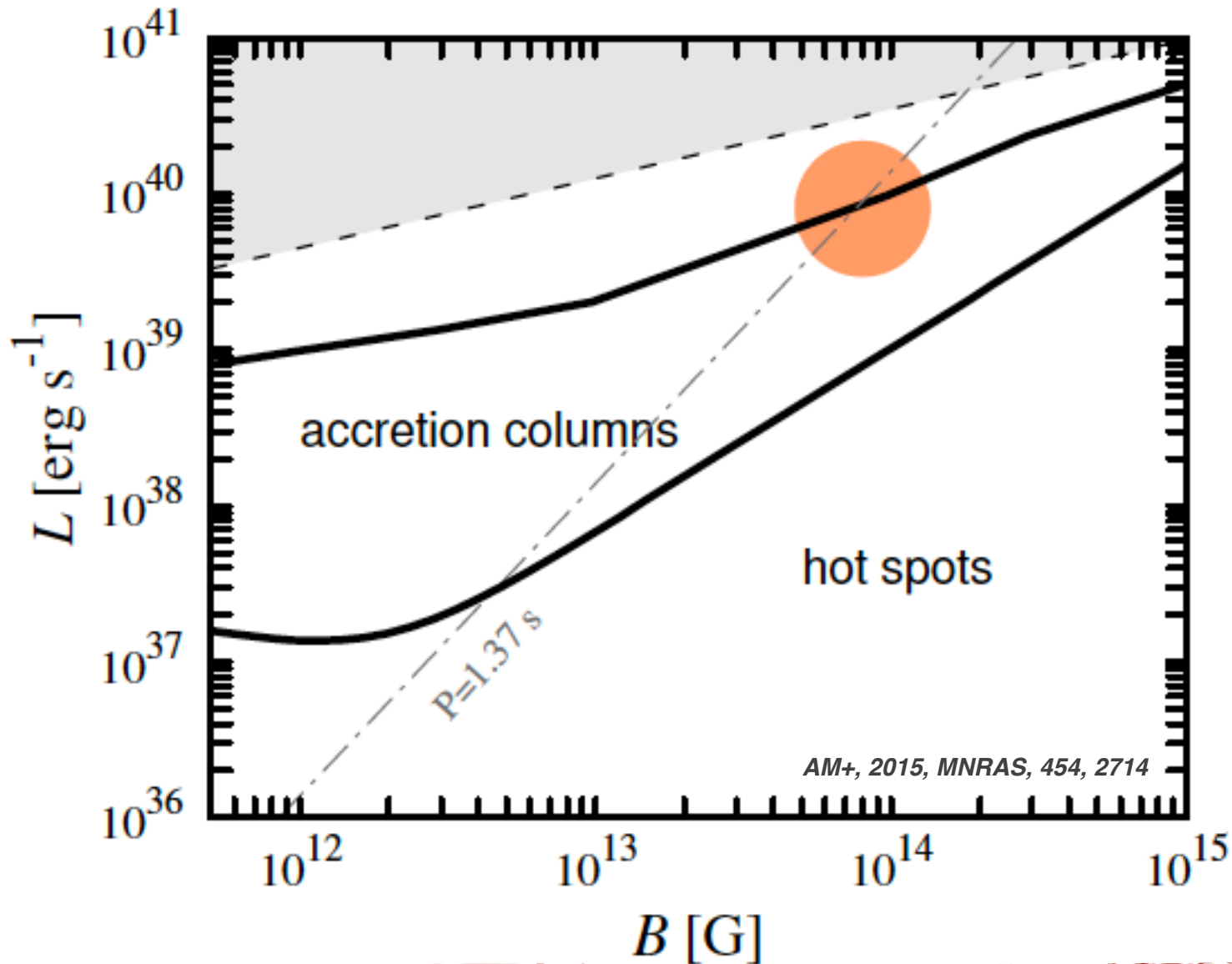
=> the higher the temperature

=> the higher photon energy

=> the higher the cross-sections

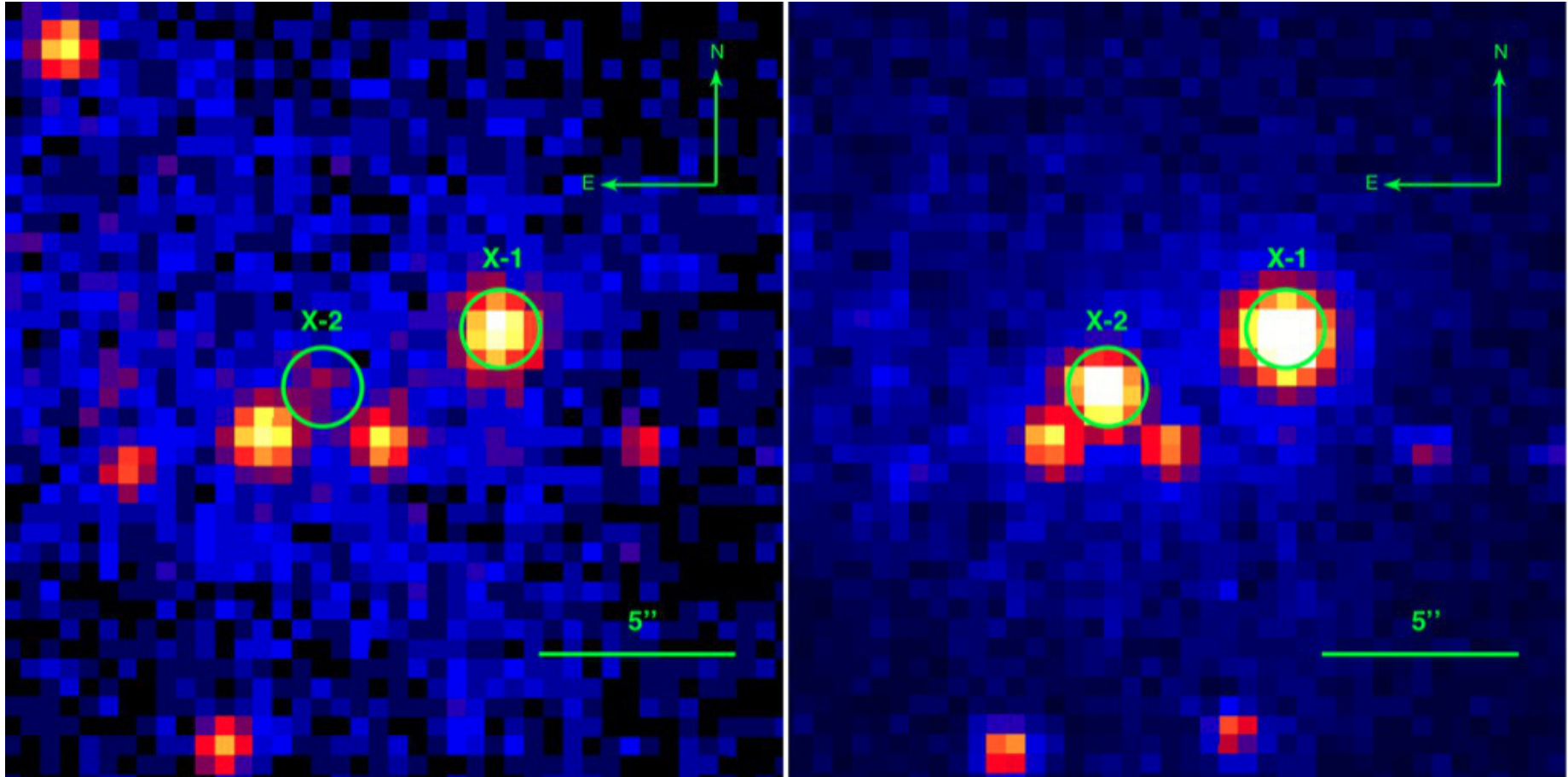


Pulsations from ULX in M82: explanation

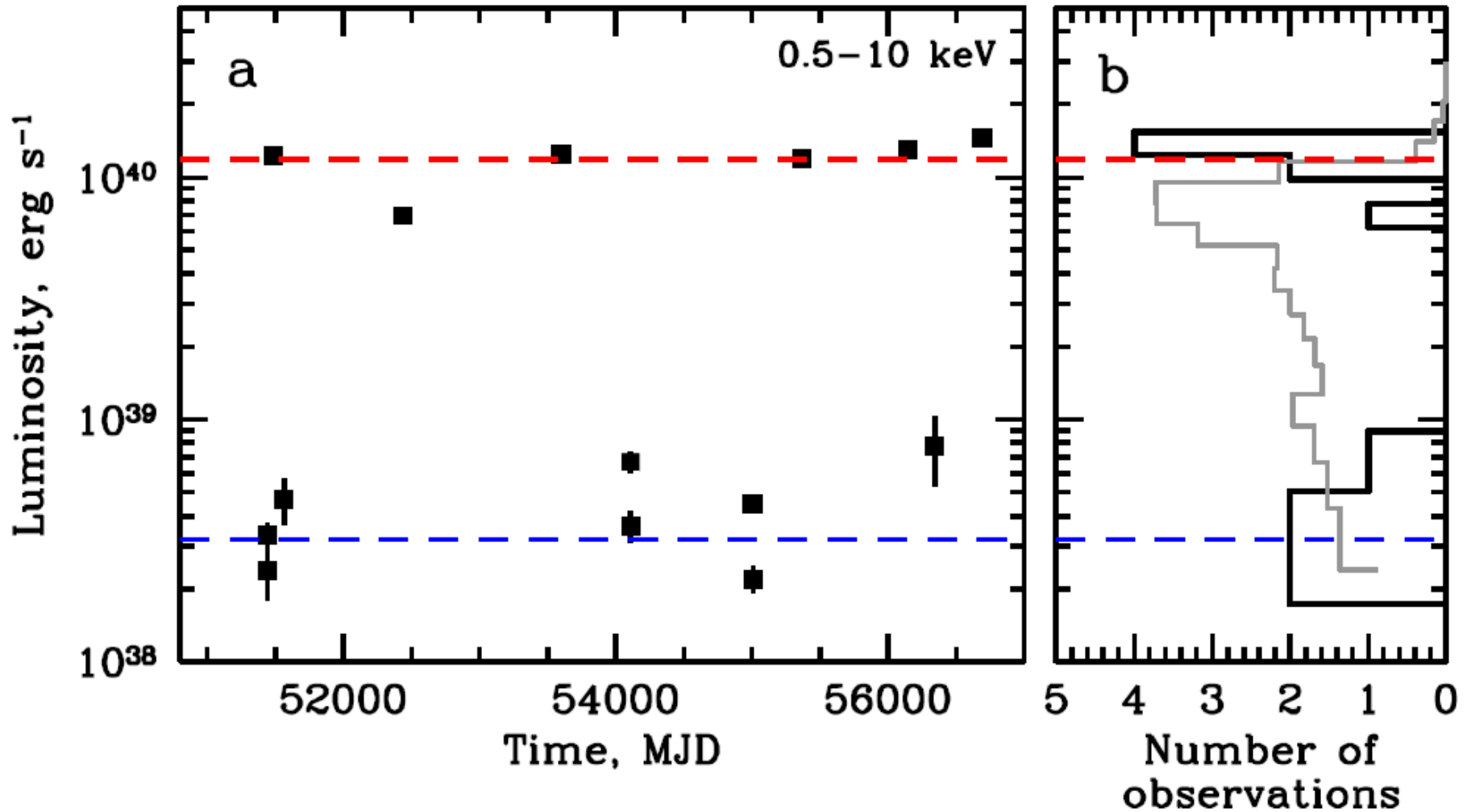


$$L_{\text{lim}}(R) \simeq \frac{GMM_{\text{lim}}}{R} \simeq 4 \times 10^{37} k^{7/2} B_{12}^2 P^{-7/3} M_{1.4}^{-2/3} R_6^5 \text{ erg s}^{-1}$$

M82 as seen by Chandra

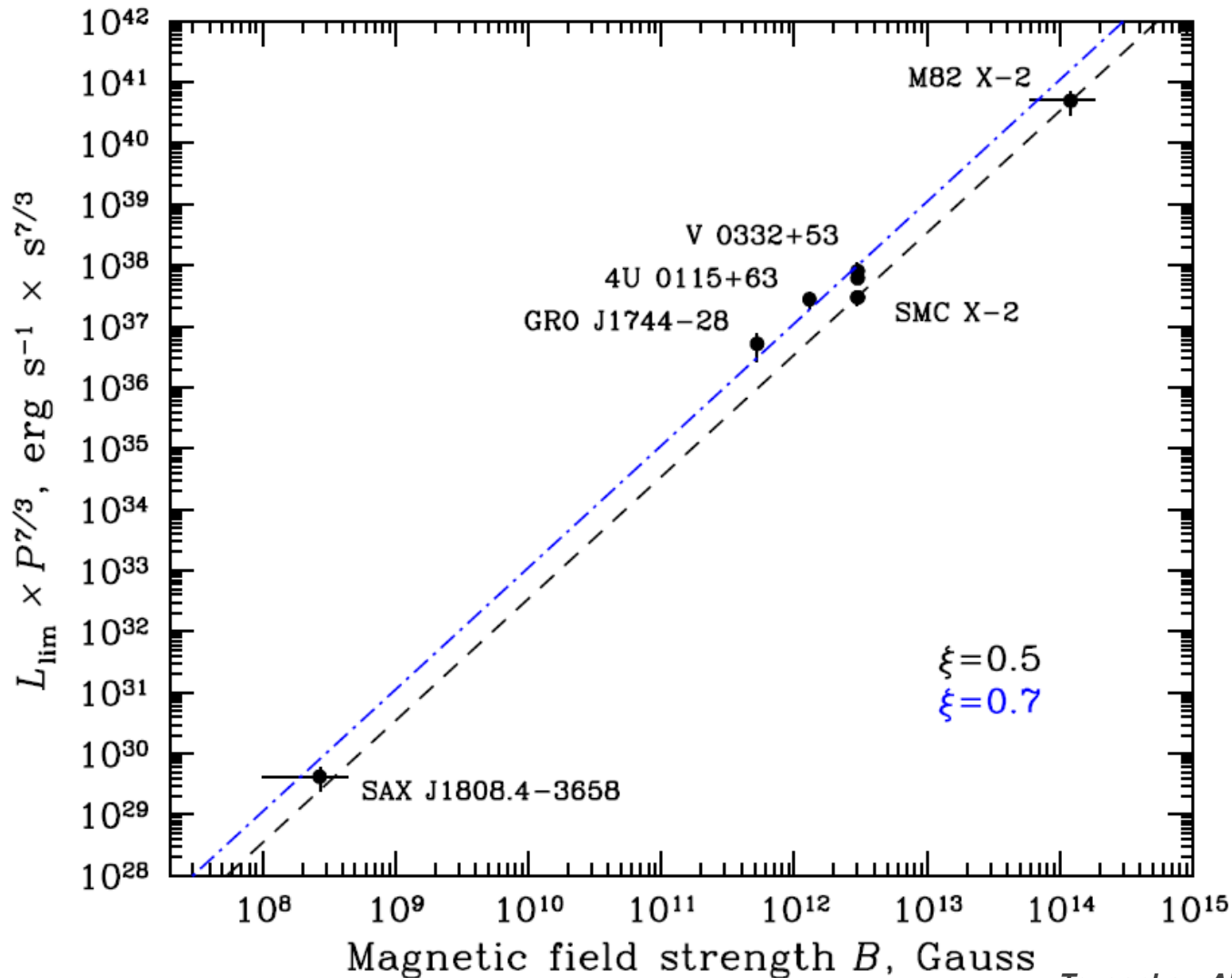


M82 X-2 intensity distribution



Propeller in action

$$L_{\text{lim}}(R) \simeq \frac{GM\dot{M}_{\text{lim}}}{R} \simeq 4 \times 10^{37} k^{7/2} B_{12}^2 P^{-7/3} M_{1.4}^{-2/3} R_6^5 \text{ erg s}^{-1}$$

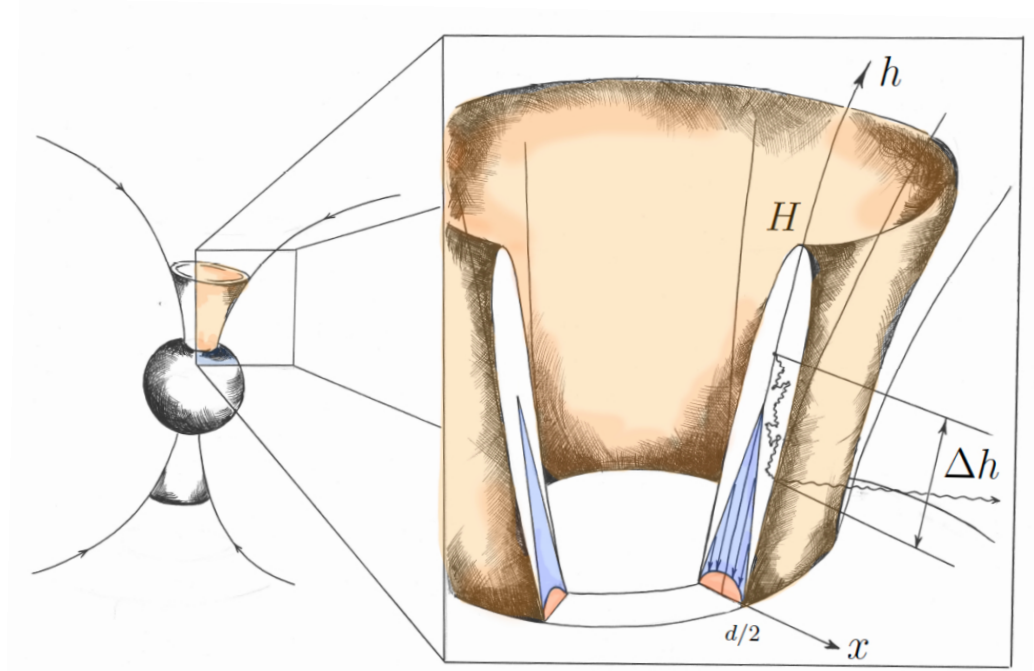


Accretion column: Advection and Neutrino pulsars

Typical time of photon escape:

$$t_{\text{diff}} = \frac{\tau d}{2c} \approx 5 \times 10^{-4} \frac{\dot{m}_{10} d_4^2 \kappa_e}{\beta} \text{ s}$$

$$\frac{\partial}{\partial h} \left[\left(-\frac{\rho GM}{R+h} + \frac{\rho v^2}{2} + \varepsilon_{\text{tot}} + P_{\text{tot}} + 2n_+ m_e c^2 \right) v \right] = Q^-$$

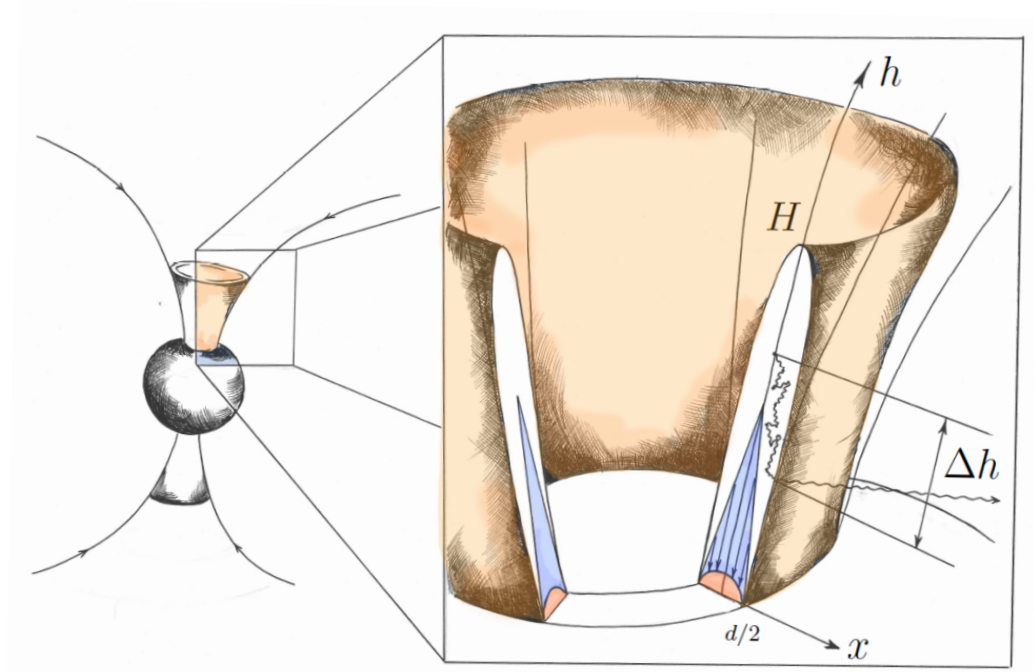


Accretion column: Advection and Neutrino pulsars

Typical time of photon escape:

$$t_{\text{diff}} = \frac{\tau d}{2c} \approx 5 \times 10^{-4} \frac{\dot{m}_{10} d_4^2 \kappa_e}{\beta} \text{ s}$$

$$\frac{\partial}{\partial h} \left[\left(-\frac{\rho GM}{R+h} + \frac{\rho v^2}{2} + \epsilon_{\text{tot}} + P_{\text{tot}} + 2n_+ m_e c^2 \right) v \right] = Q^-$$

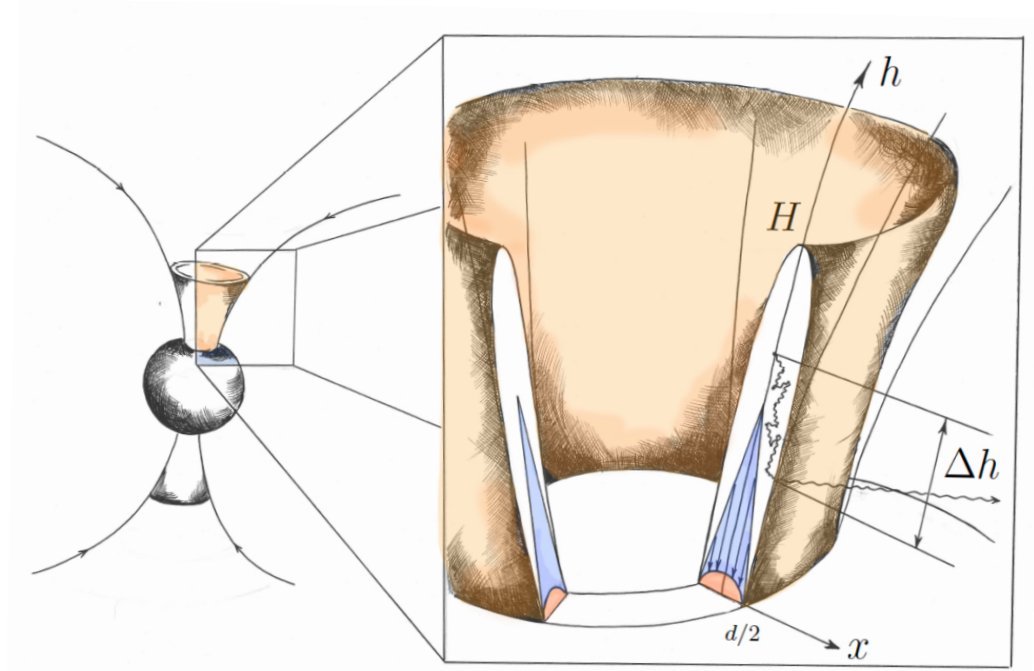


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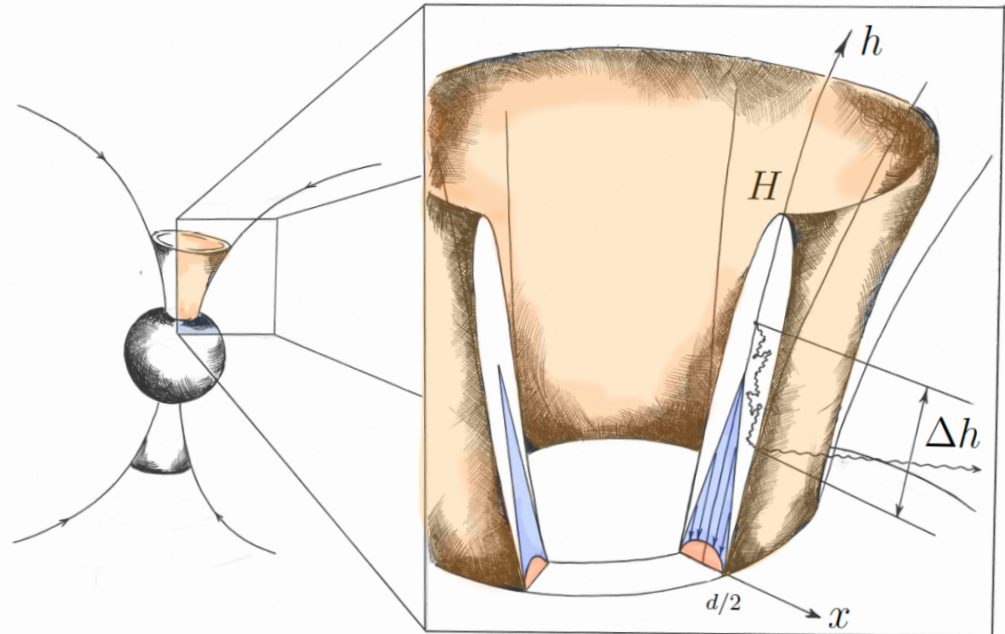
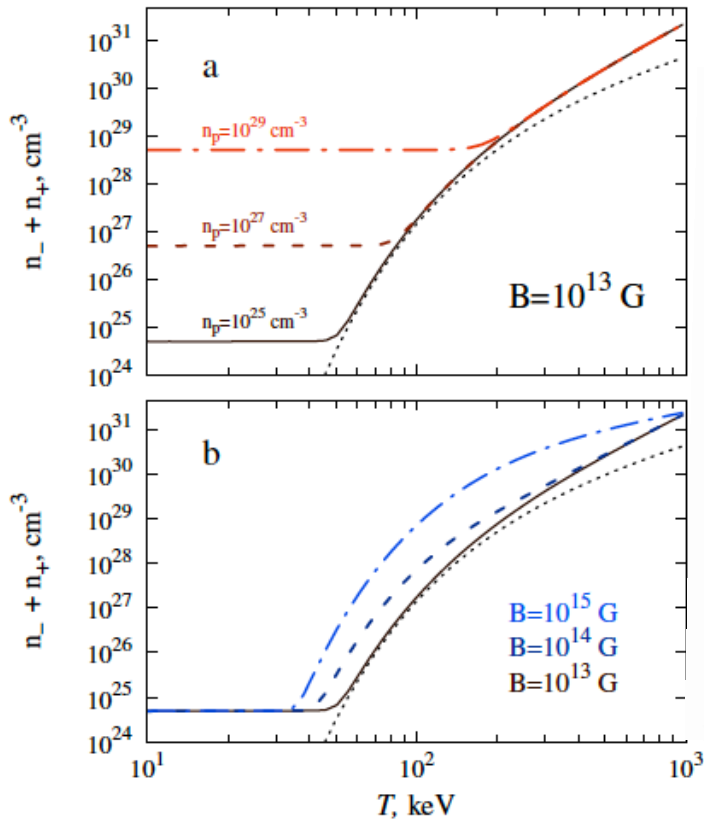


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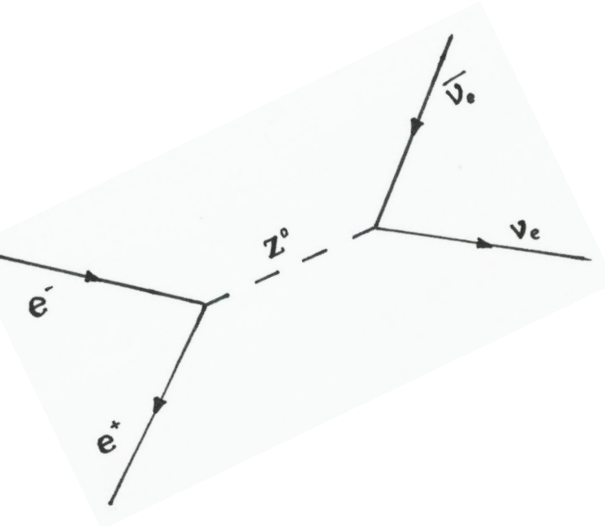
$$\frac{\partial}{\partial h} \left[\left(-\frac{\rho GM}{R+h} + \frac{\rho v^2}{2} + \varepsilon_{\text{tot}} + P_{\text{tot}} + 2n_+ m_e c^2 \right) v \right] = Q^-$$



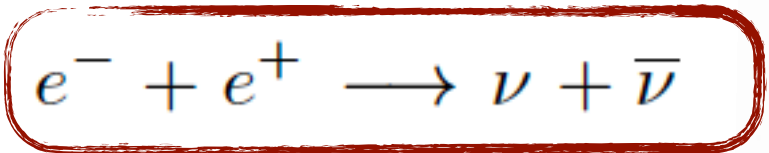
Accretion column: Advection and Neutrino pulsars

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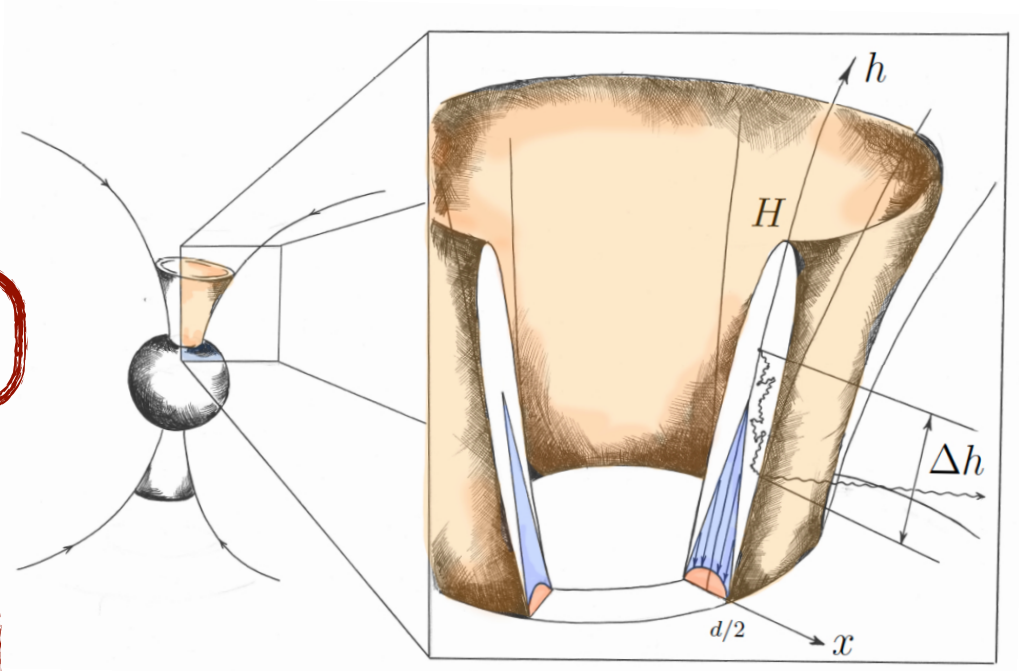


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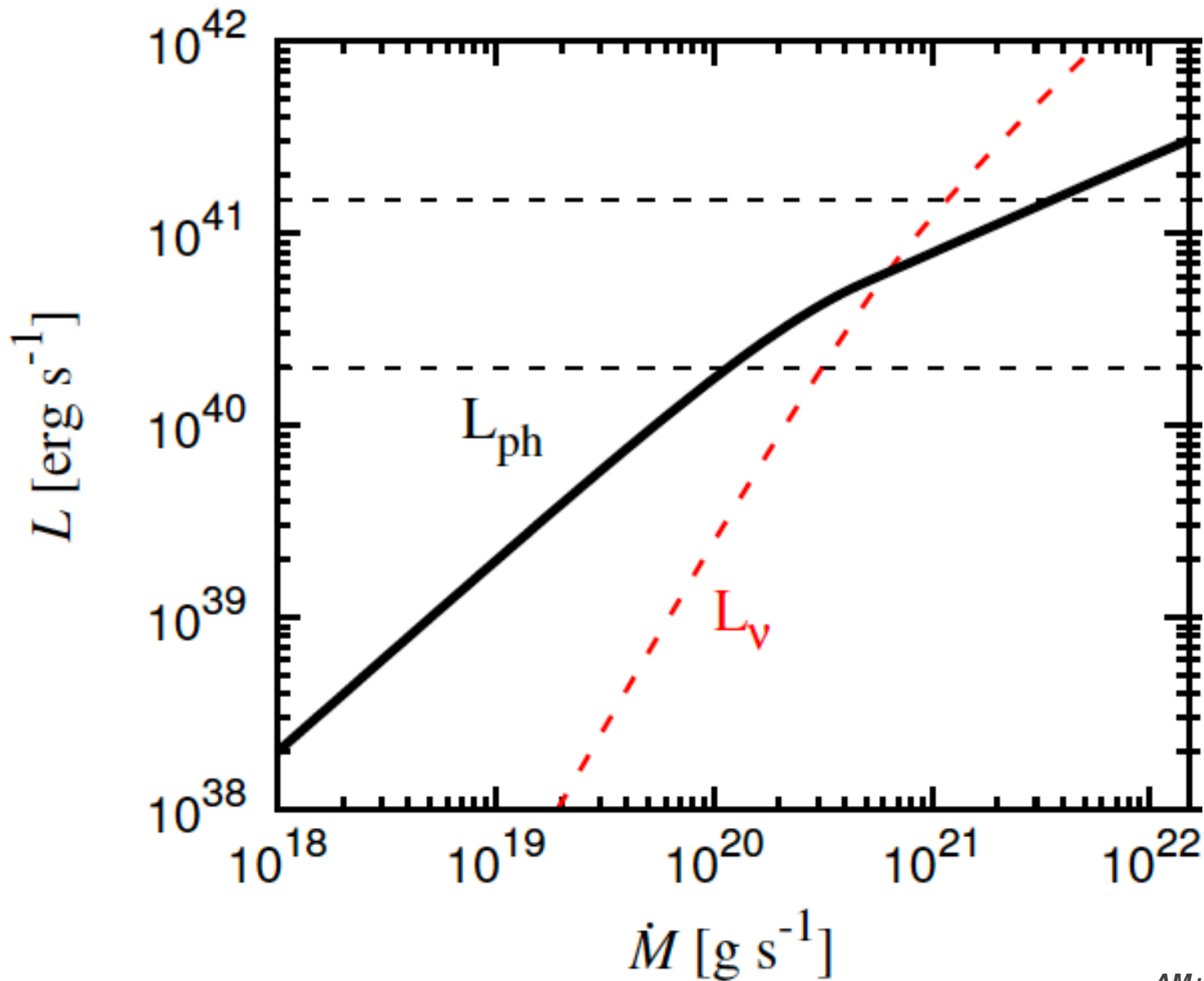


The total accretion luminosity:

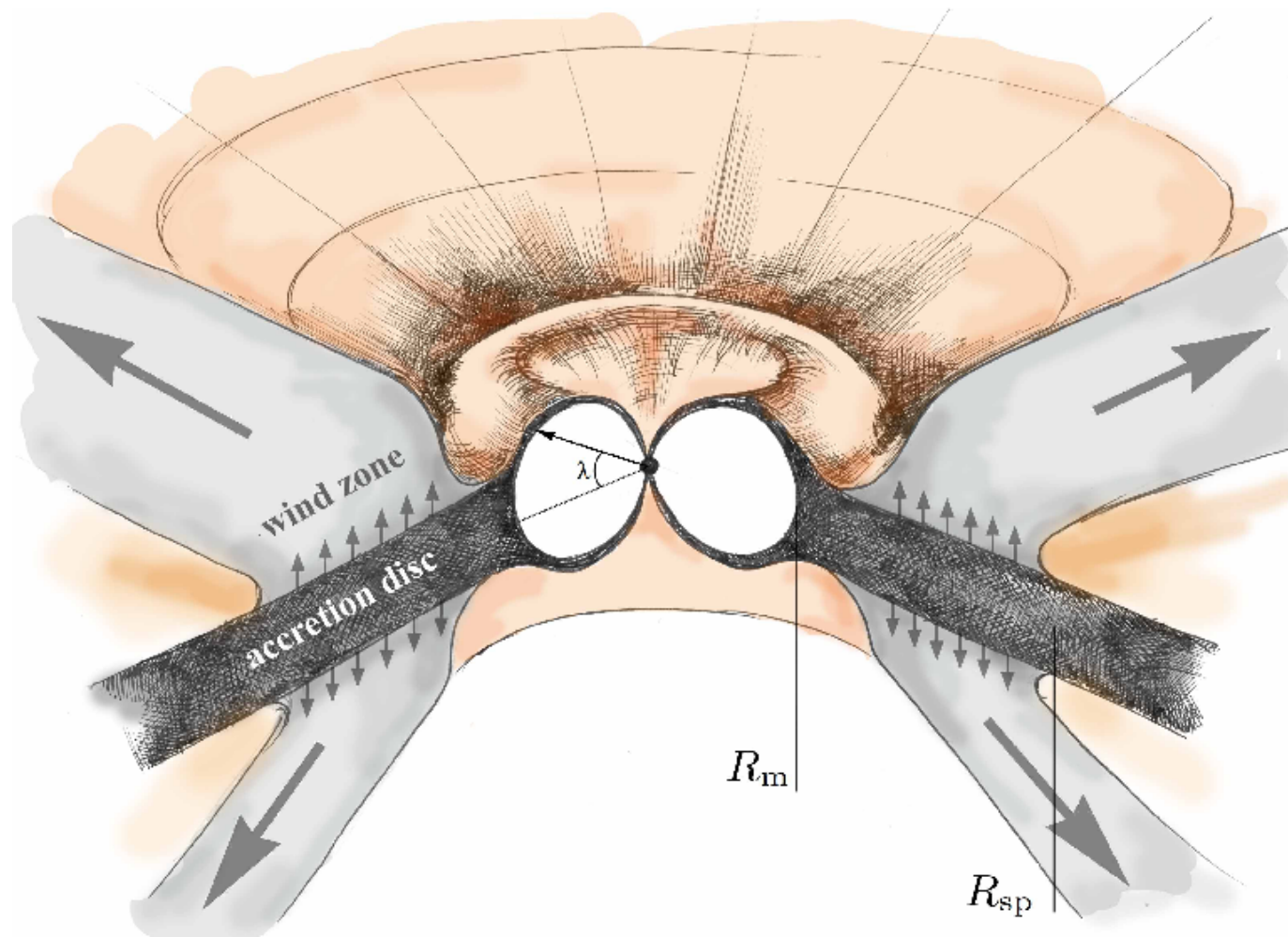
$$L_{\text{tot}} = \frac{GM\dot{M}}{R} = L_{\text{ph}} + L_{\nu}$$



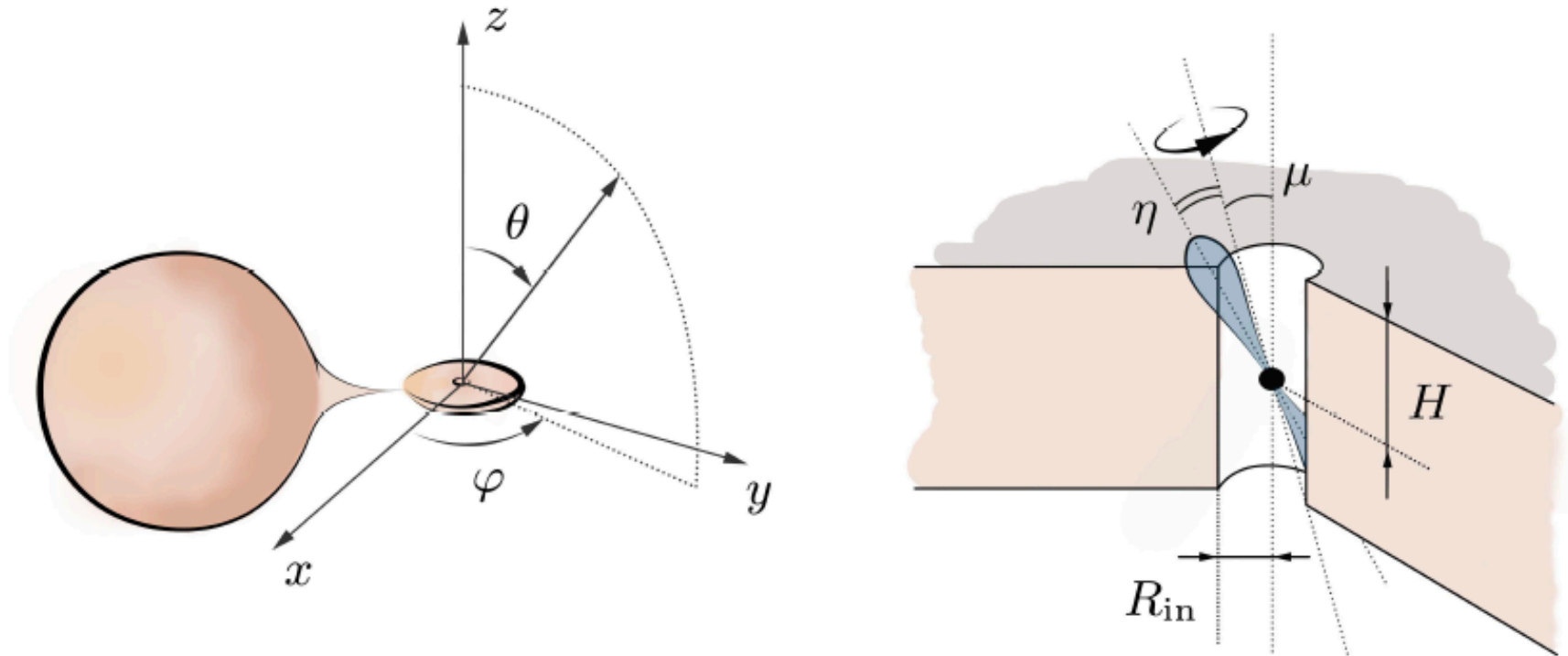
Accretion column: Photon and Neutrino Luminosity



Outflows from accretion discs in ULX pulsars



Geometrical Beaming vs. Pulsed Fraction



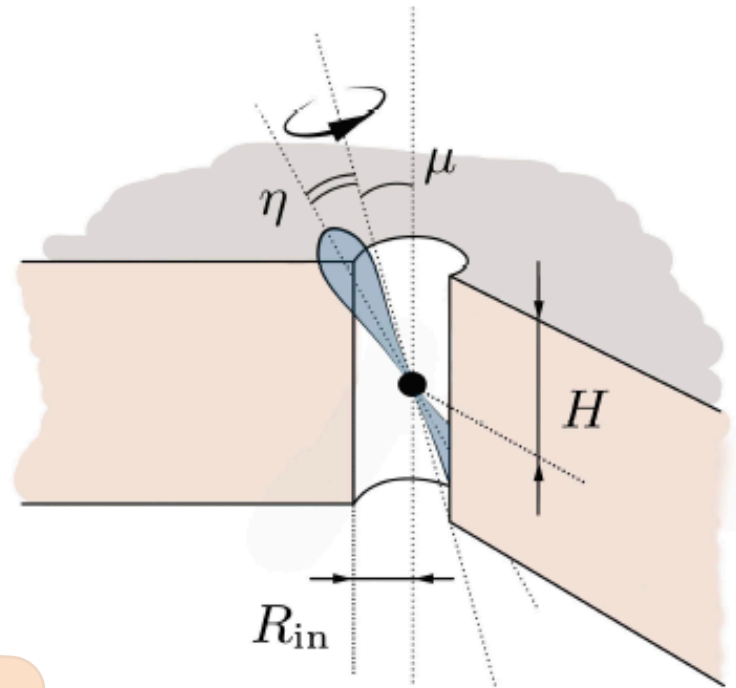
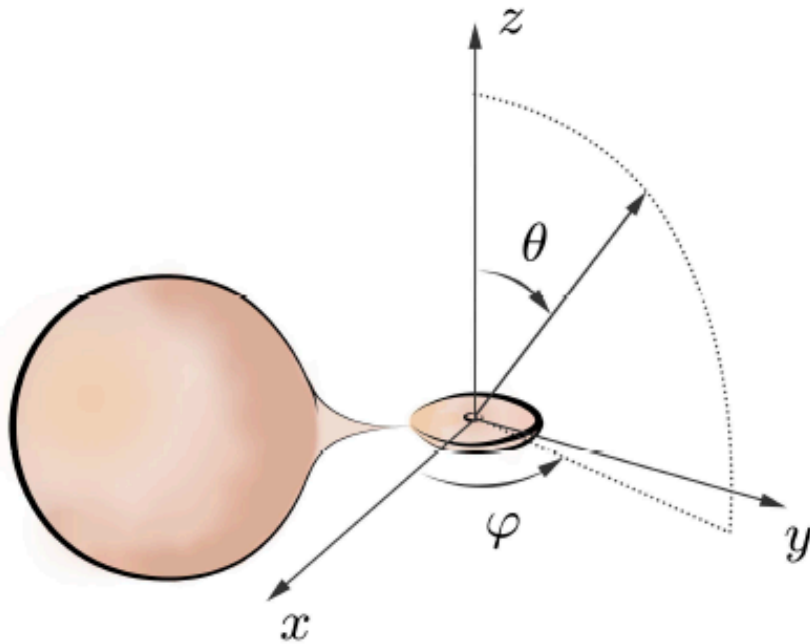
We know **5 pulsating ULXs**.

But, there are only **~15 ULXs** out of **~300** provide the statistics sufficient for detection of pulsations.

(see, e.g., Rodrigues Castillo+, 2020, ApJ, 895)

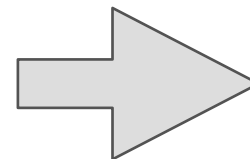
High Pulsed Fraction (~10 percents and more) is a typical feature of ULX pulsars.

Geometrical Beaming vs. Pulsed Fraction



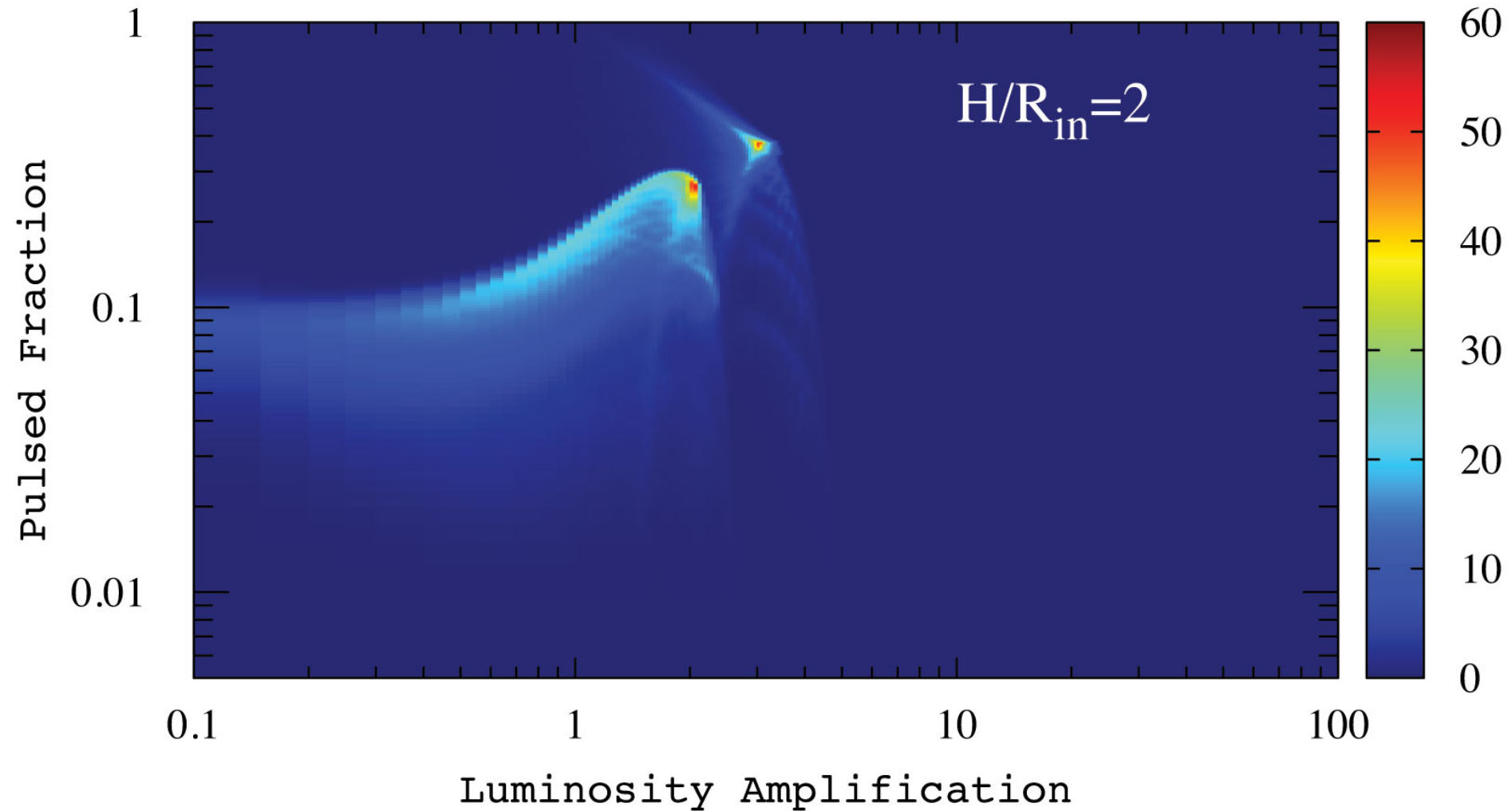
A certain **beam geometry** from a NS
Parameters of rotation in respect to the accretion flow
 H/R parameter

We use Monte Carlo simulations and trace photons
accounting for:
GR effects
&
reflection of photons from the walls of accretion funnel

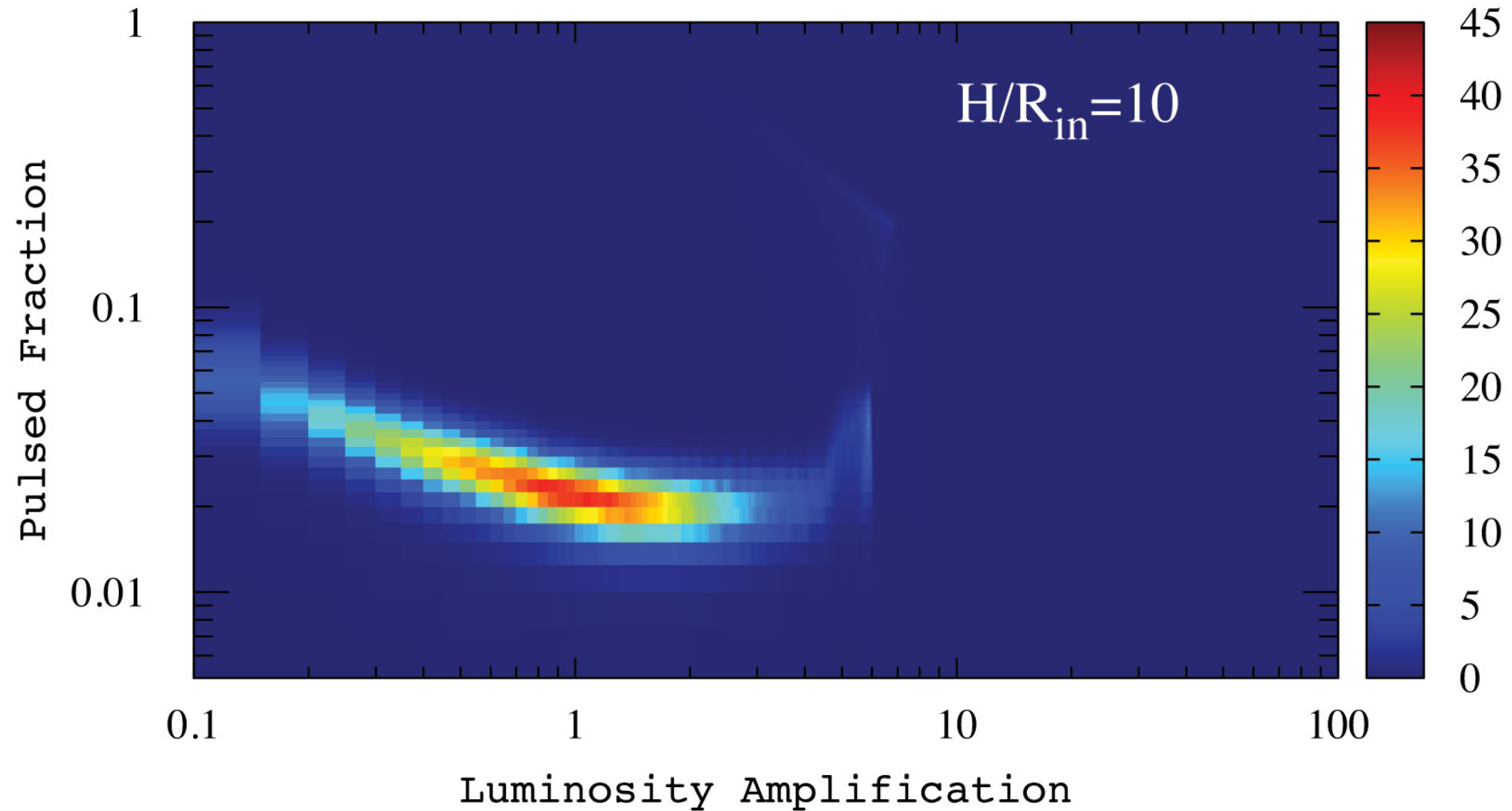


We reproduce
the pulse profiles
for different observers

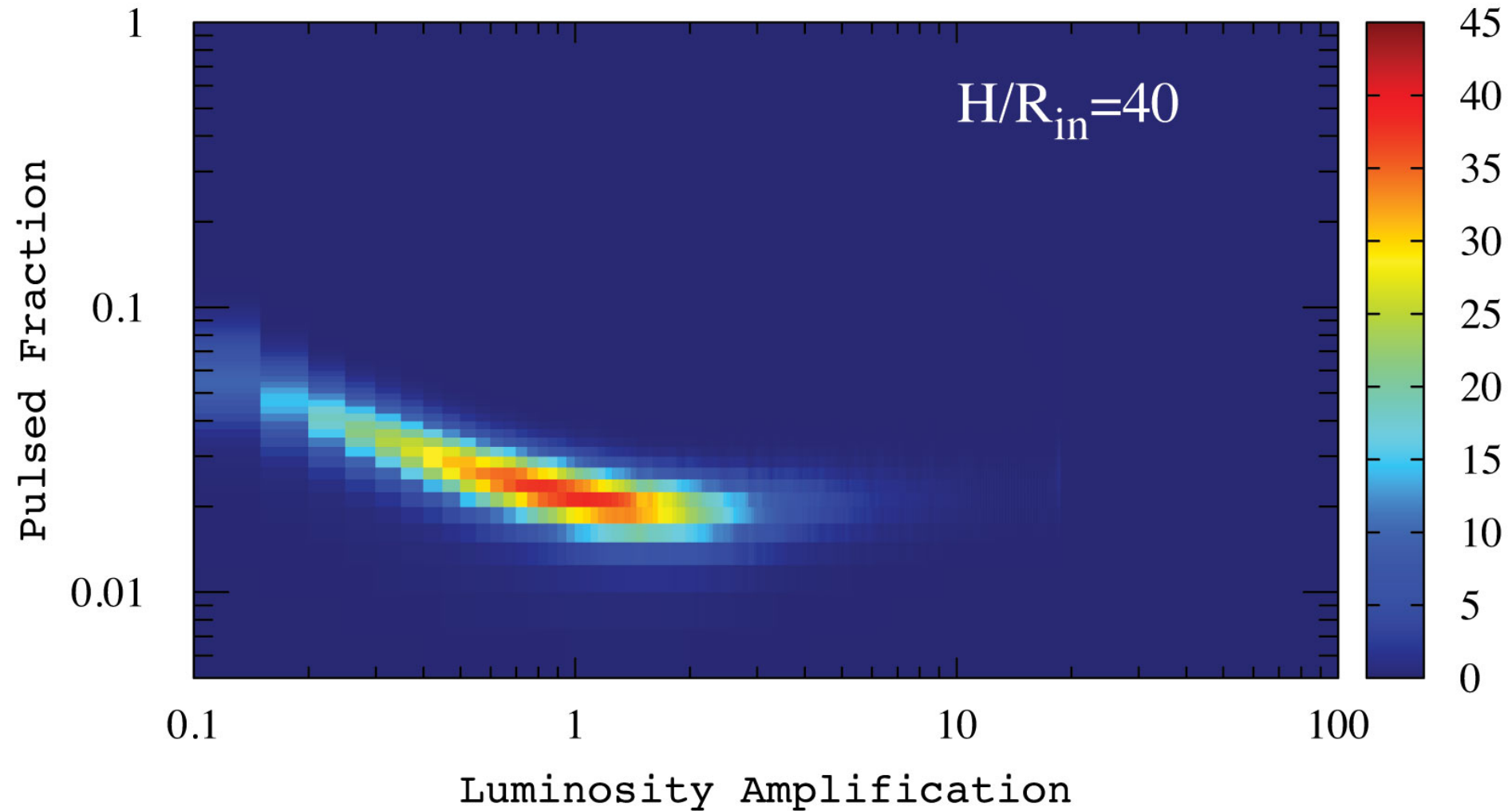
Distribution of ULX pulsars over the PF and Luminosity Amplification Factor



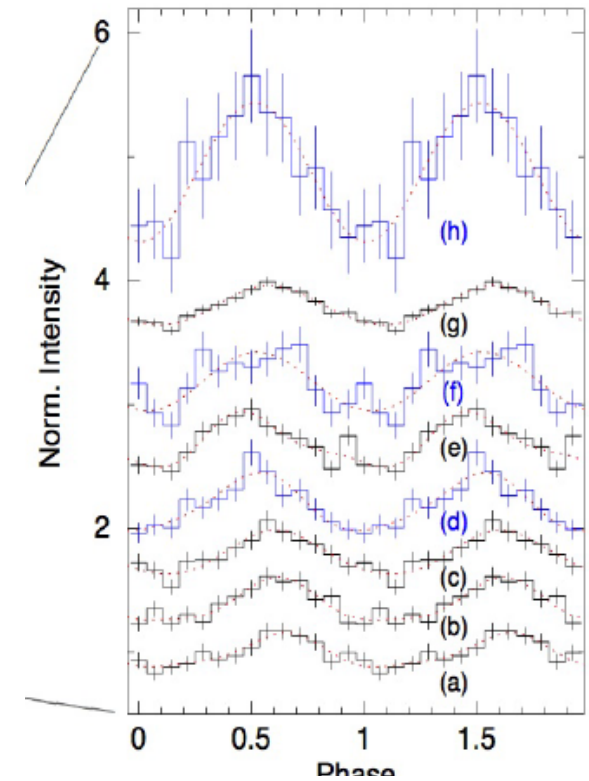
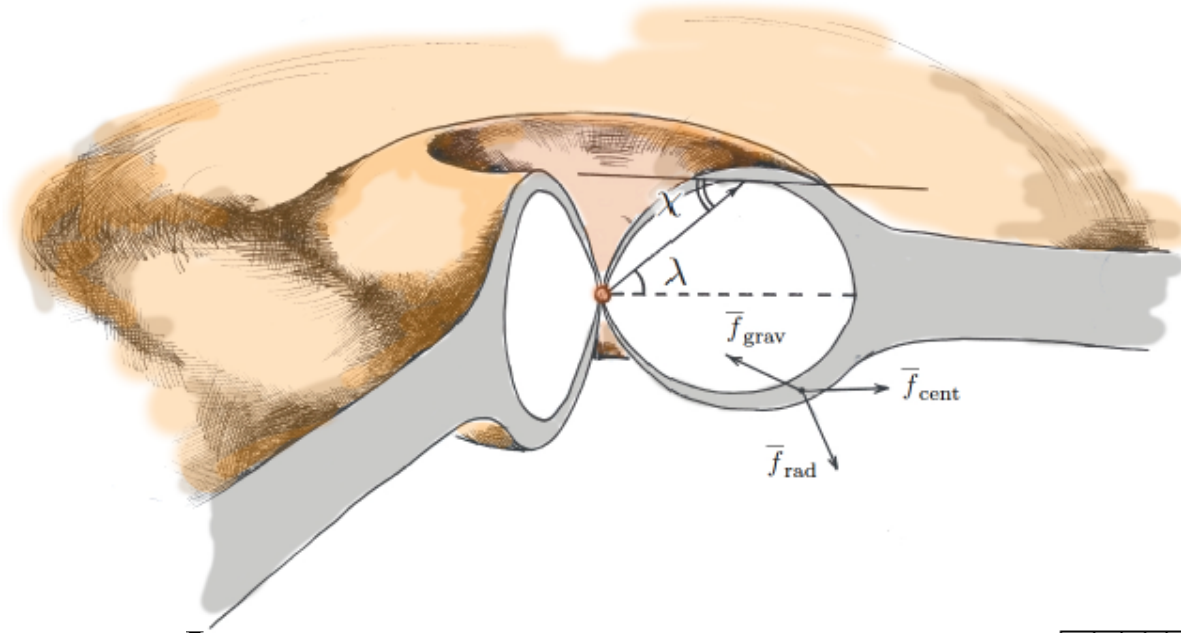
Distribution of ULX pulsars over the PF and Luminosity Amplification Factor



Distribution of ULX pulsars over the PF and Luminosity Amplification Factor



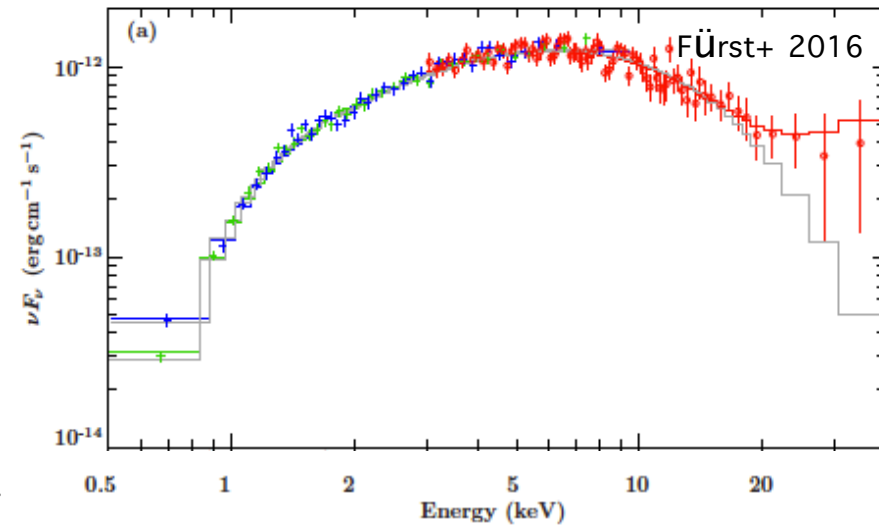
Accretion envelope



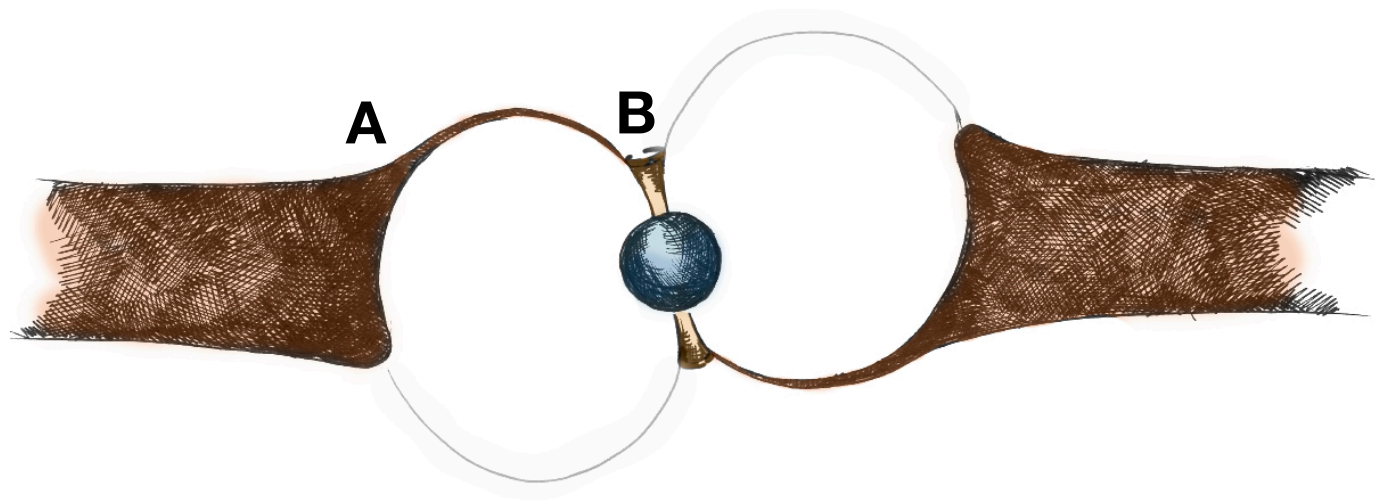
Israel+, Science, 2017

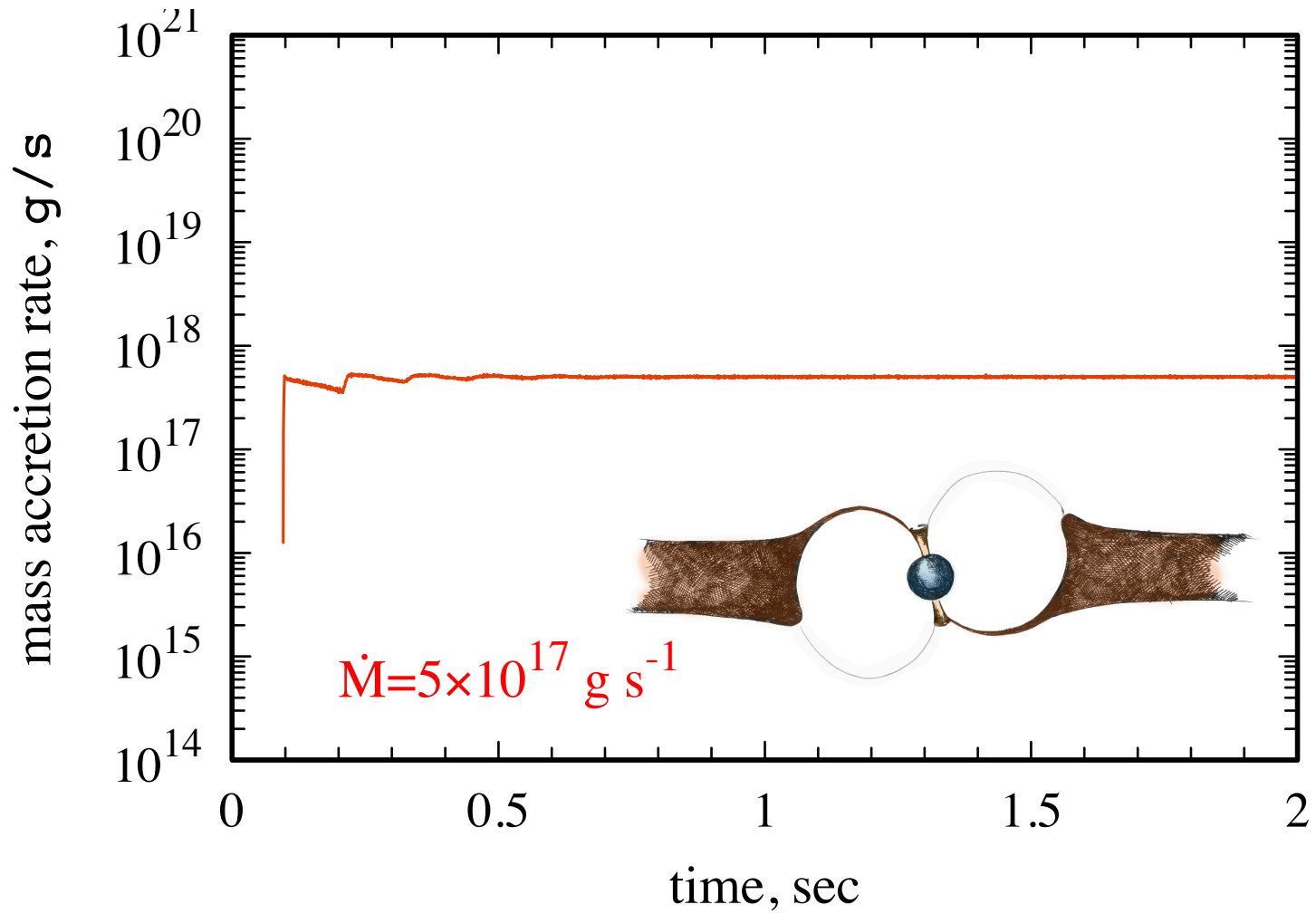
Important consequences:

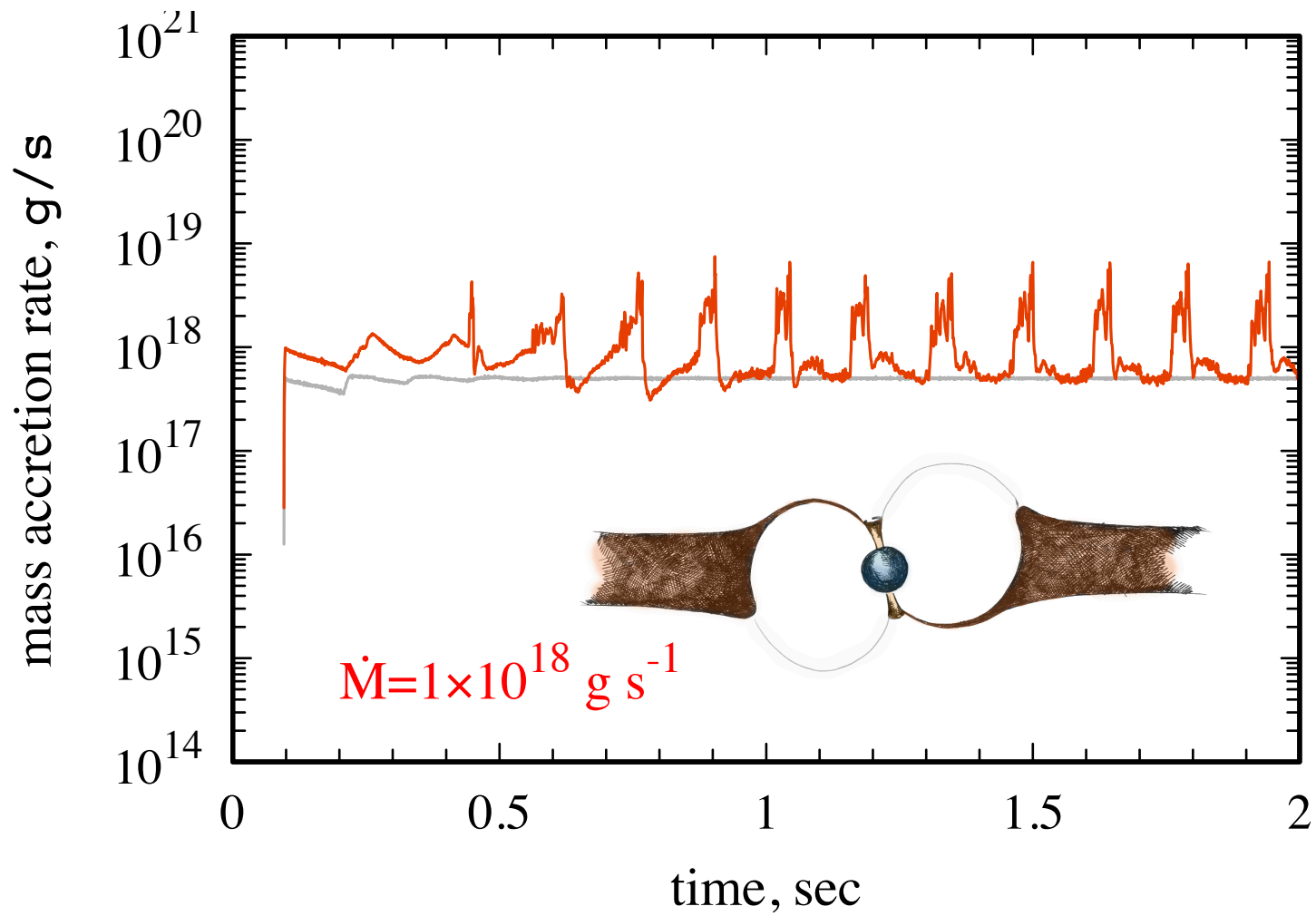
- (1) we do not see directly the central NS in ULX pulsars;
- (2) **smooth pulse profiles** and hardly detected **cyclotron lines**;
- (3) **the energy spectra** of ULXPs are affected by multiple scatterings in the envelope;
- (4) suppressed **power spectra** at high Fourier frequency;
- (5) **super-orbital variability** because of precession of magnetic dipole.

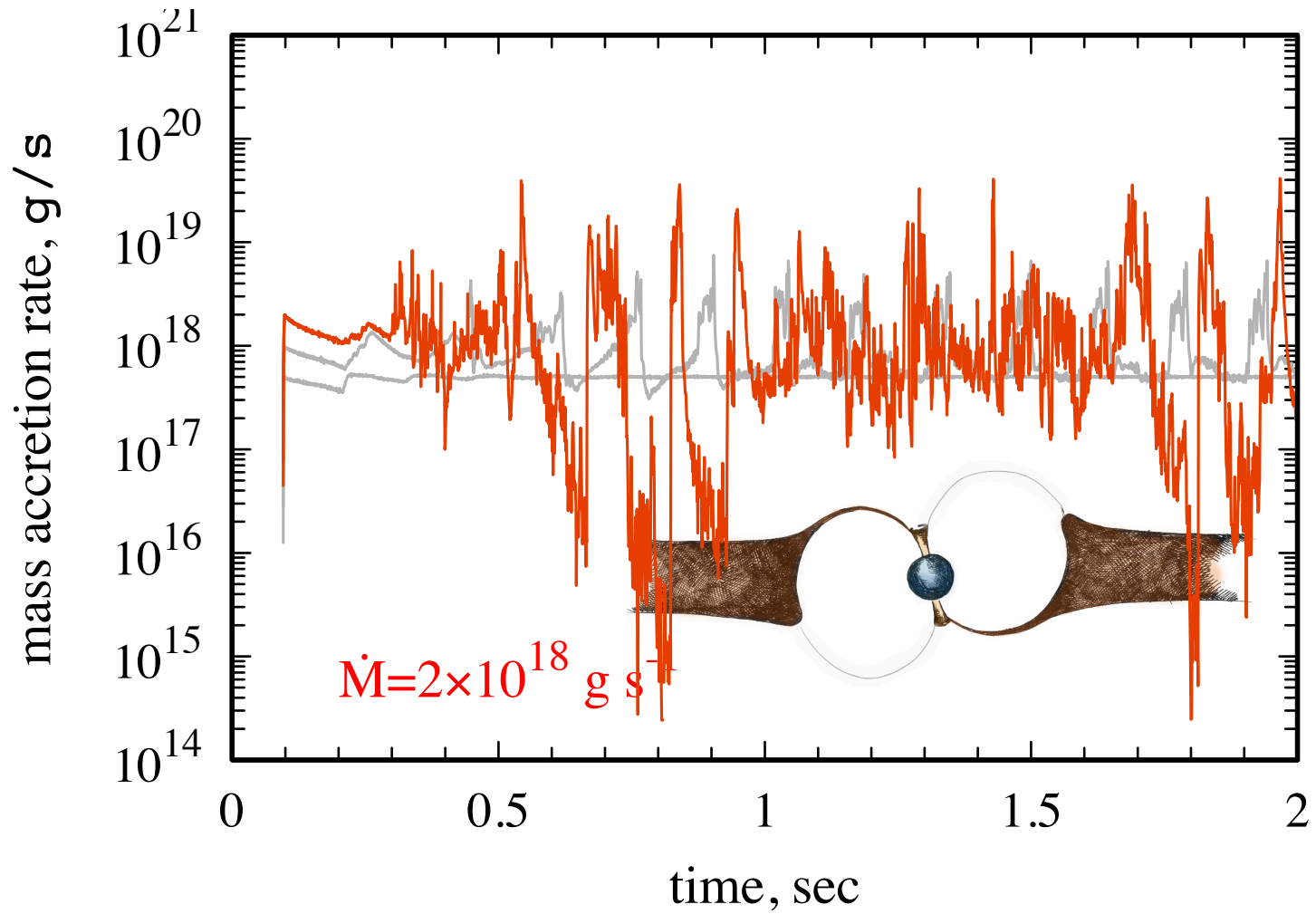


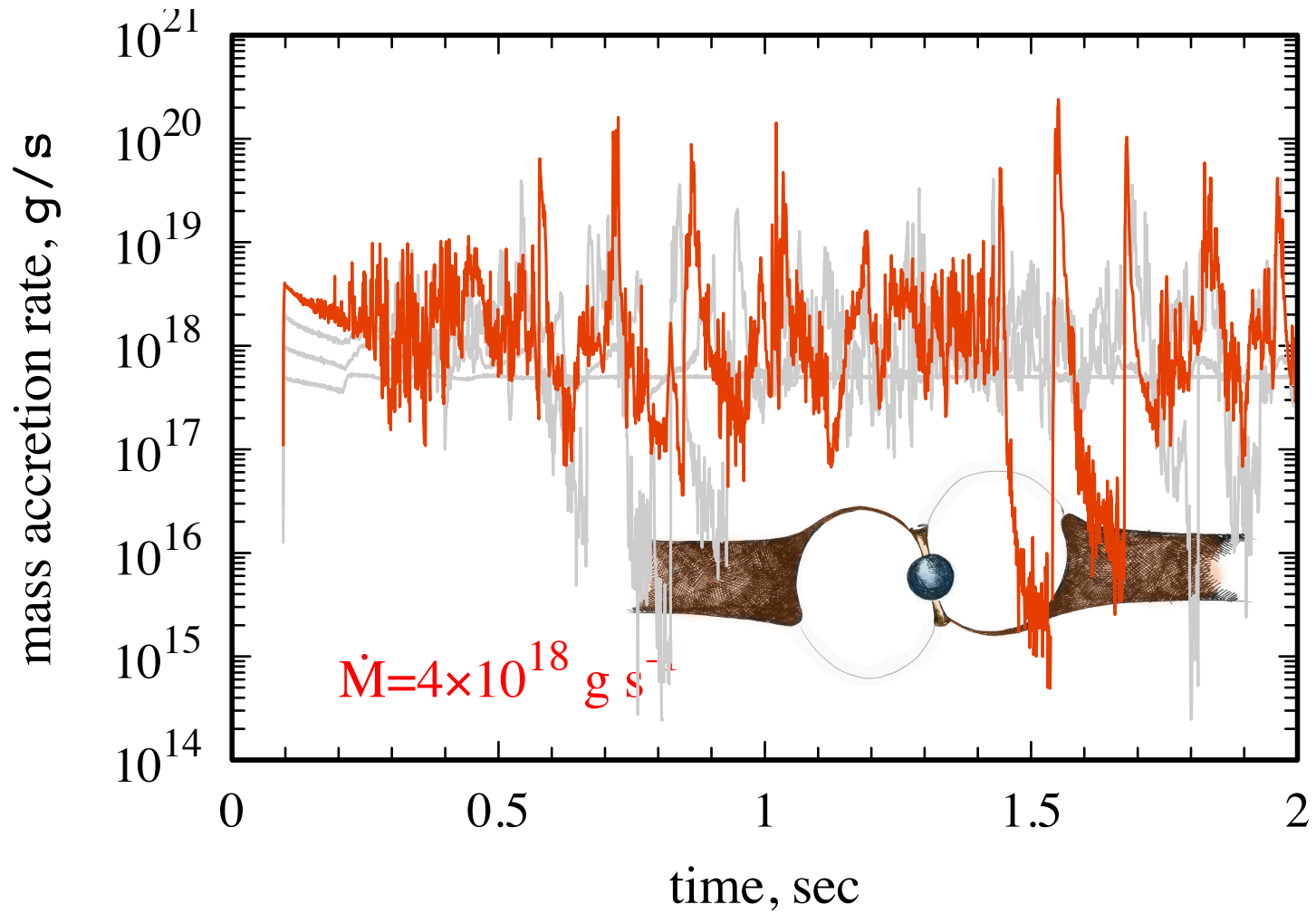
AM+, 2017, MNRAS, 467
AM+, 2019, MNRAS, 484

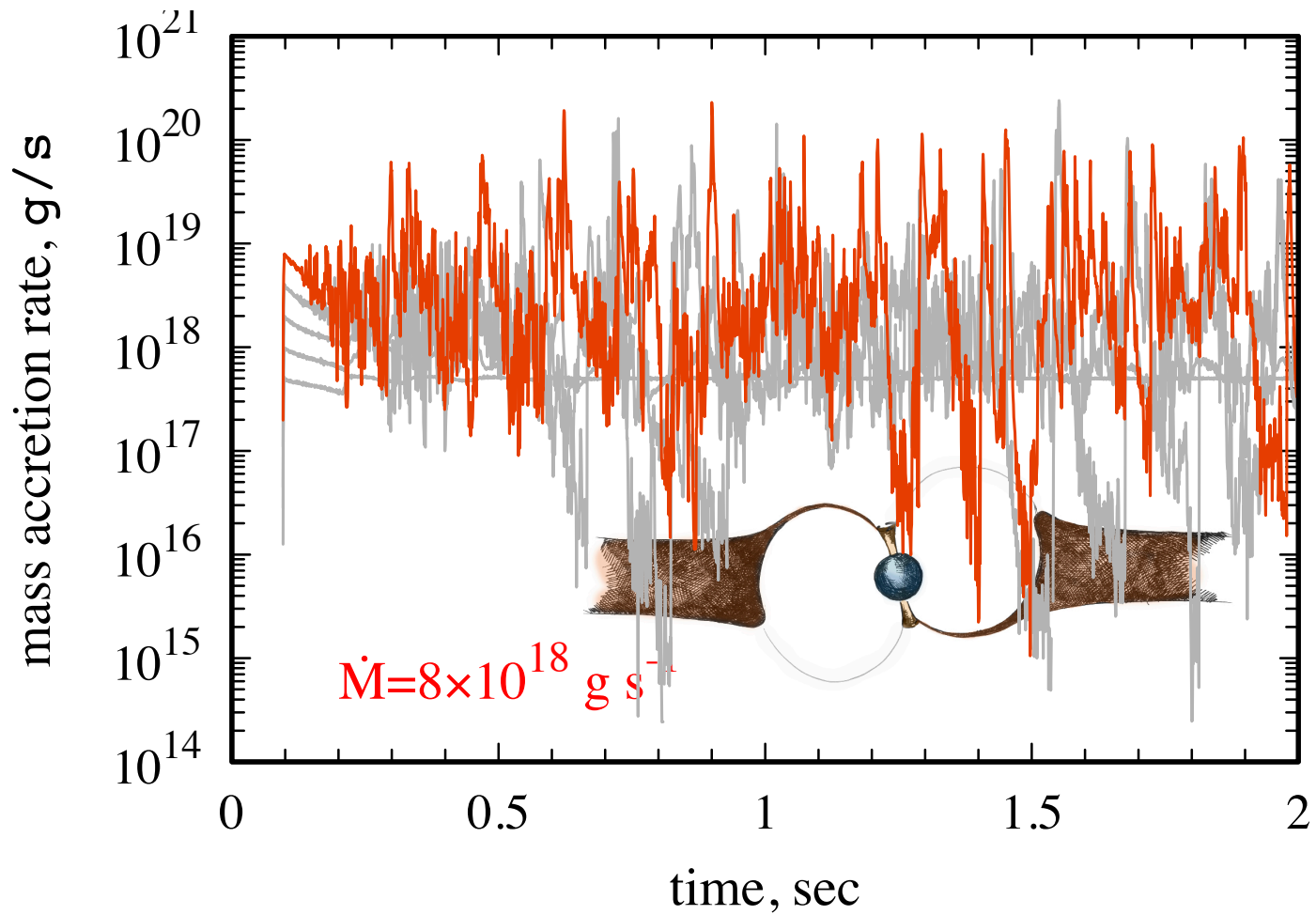






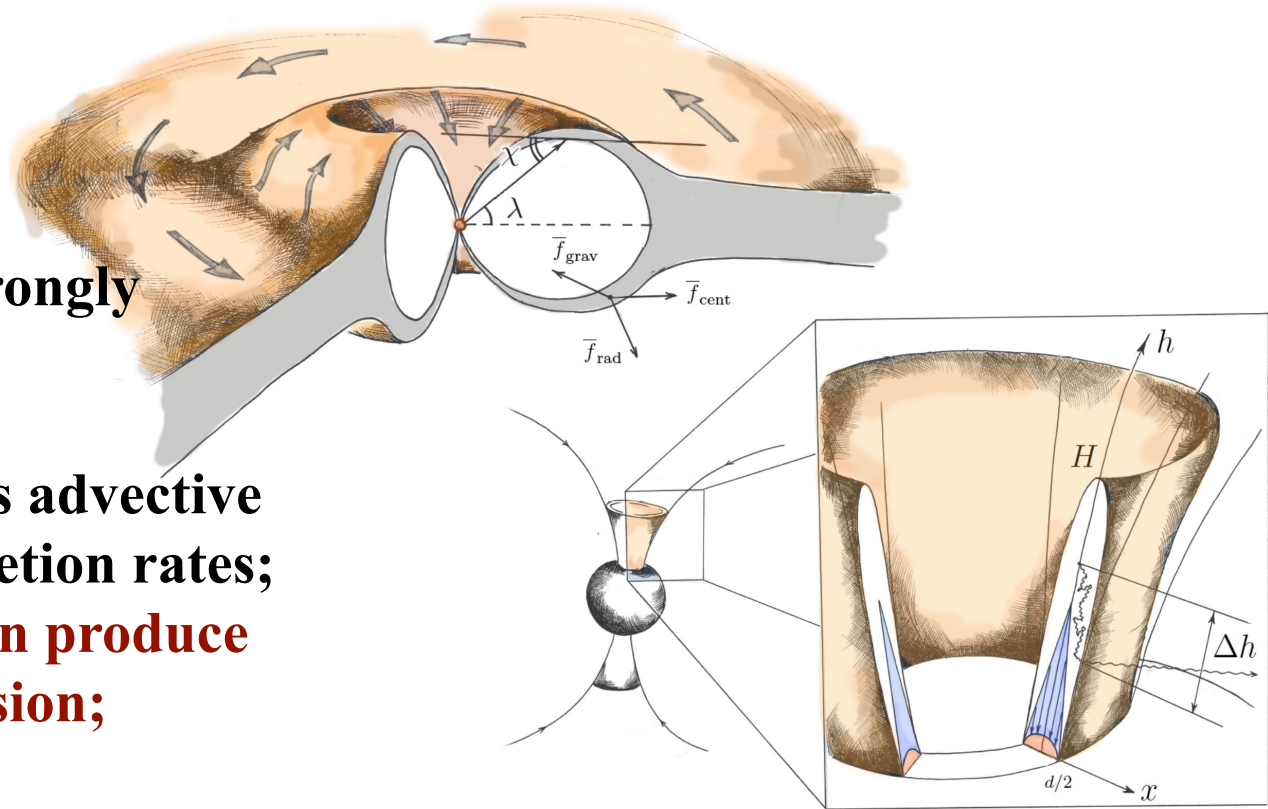






Short Summary

- (1) **Accretion columns are the central engines in ULXs;** their luminosity is strongly affected by geometry of accretion channel;
- (2) The column becomes advective at extreme mass accretion rates; **advective columns can produce strong neutrino emission;**
- (3) Bright ULX pulsars are surrounded by **optically thick envelopes.** The envelopes determine **the observational manifestation** of ULX pulsars;
- (4) **Strong outflow** from the accretion disc in ULX pulsars is possible in the case of relatively weak dipole component of magnetic field



But

many and many details remain unclear and/or debated.

- (1) magnetic field strength
- (2) evolutionary status of ULX pulsars
- (3) fraction of NS among ULXs
- (4) fate of a companion star
- (5) ...

